Heating of braided coronal loops



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Aim: investigate how the pattern of braiding by photospheric footpoint motions affects heating of coronal loops

The model magnetic fields



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Achieved in practice by adding regions of twist to uniform B

-12

Aspect ratio of loops is $\approx 1:10$

Conservative approach: free energy only 3% above potential

Simulation setup

- Take field E3 or S3 and first perform an <u>ideal</u> relaxation
- Then transfer to resistive MHD code: *J*×<u>B</u>≈0, and initialise with a uniform background plasma



20

15

10

5

0

-5

-10

-15

-20

-5

5

0



- Following an initial instability current peaks sharply in both cases
- Peak current falls off quickly for S3
- Magnetic field 'unbraids'. E.g. E3 -











Energy / heating

- To investigate heating, make an appropriate dimensionalisation
- Parameters:
 - **★** B=10G
 - * n=10¹⁵ m⁻³
 - * L=I Mm
 - $* T=2.3 \times 10^{6} K$
 - * Loop dimensions: 6×6×48 Mm

 $* t_0 = 1.45s$

Energy / heating





- Approx twice as much energy released for E3
- More spatially homogeneous heating for E3
- Temperature rise is modest, but so is initial free energy

Structure of final magnetic field

Plots: mean value of $J \cdot B/B \cdot B$ along field lines



Reynolds number comparison for E3

|]| at z=0

η=10-3





Summary

- Resistive relaxation: <u>B</u> field is unbraided (E3) / untwisted (S3). Involves reconnection at multiple <u>J</u> sheets.
- Although "amount" of photospheric driving in the same, relaxation is more efficient for E3:
 - * Current sheets fill the volume more effectively
 - * More energy is released
 - * Homogeneous heating of the loop
- In other words, <u>amount and distribution</u> of energy release dependent on <u>pattern</u> of driving flow.
- This can be measured by computing the "topological entropy" of the photospheric flow
- Energy release constrained by structure of <u>B</u> (periodic orbits)

Thanks for listening

References:

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