

Effective field theory for perturbations in dark energy & modified gravity

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The problem...

- * General Relativity + FRW + standard model particles + observational data:
inconsistent
 - * ... invent dark energy $\sim 70\%$.
- * Or perhaps GR is not the right gravitational theory for cosmological scales...
 - * c.f. Newton & Mercury / Sun system
- * **Models of dark energy:** Λ , quintessence, k-essence, elastic dark energy, ...
- * **Modified gravity:** F(R), Horndeski, galileon, Gauss-Bonnet, Aether, TeVeS, ...
 - * MG... obtain different gravitational potential for the same matter content
- * Lagrangian engineering

Generalized gravitational field equations

- ❖ All modified gravity & dark energy theories have gravitational field equations which can be written as

$$G_{\mu\nu} = 8\pi G T_{\mu\nu} + U_{\mu\nu}$$

- ❖ Stems from an action:

$$\mathcal{L}_{\text{grav}} = \mathcal{L}_{\text{GR}} + \mathcal{L}_{\text{known}} + \mathcal{L}_{\text{dark}}$$

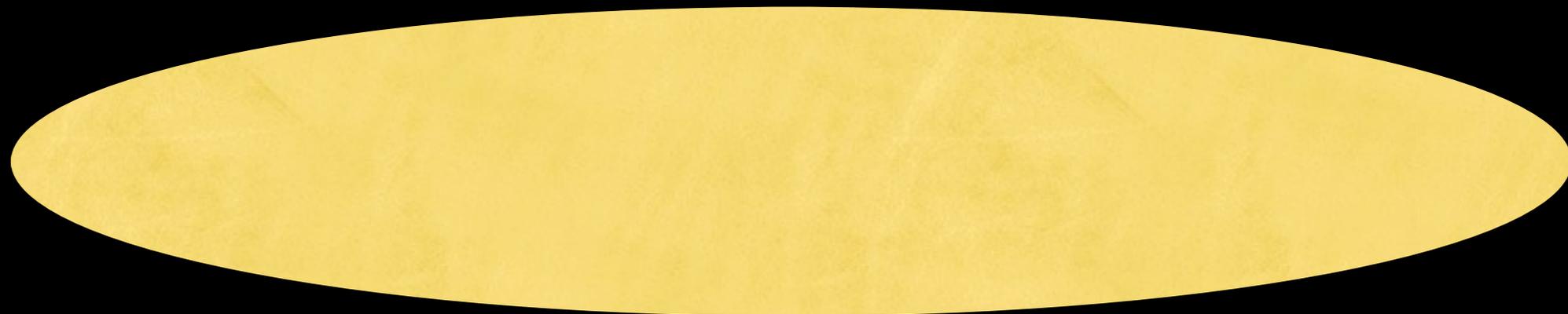
- ❖ At perturbed order: structures...

$$\delta G_{\mu\nu} = 8\pi G \delta T_{\mu\nu} + \delta U_{\mu\nu}$$

- ❖ **Q:** How do we write down the allowed, consistent modifications to the gravity field equations? **A:** *Lagrangian for the dark sector perturbations*

Effective theories: philosophy

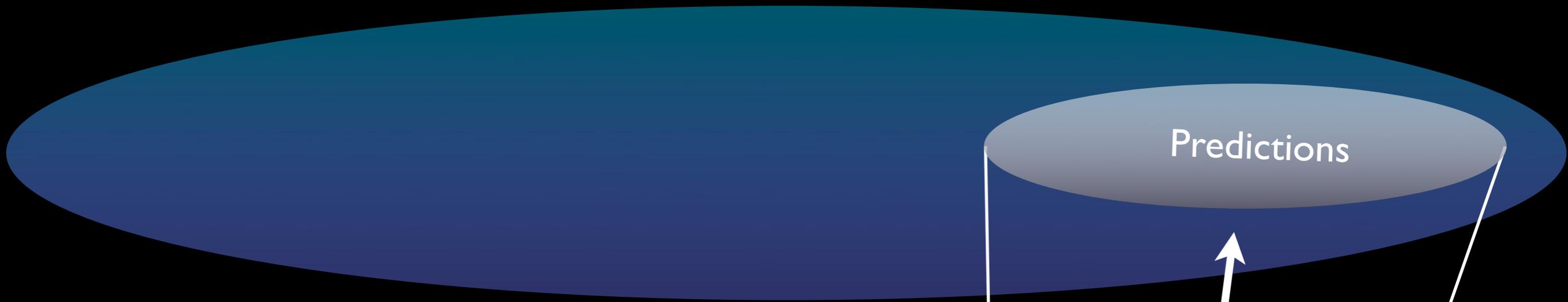
“Observables” space



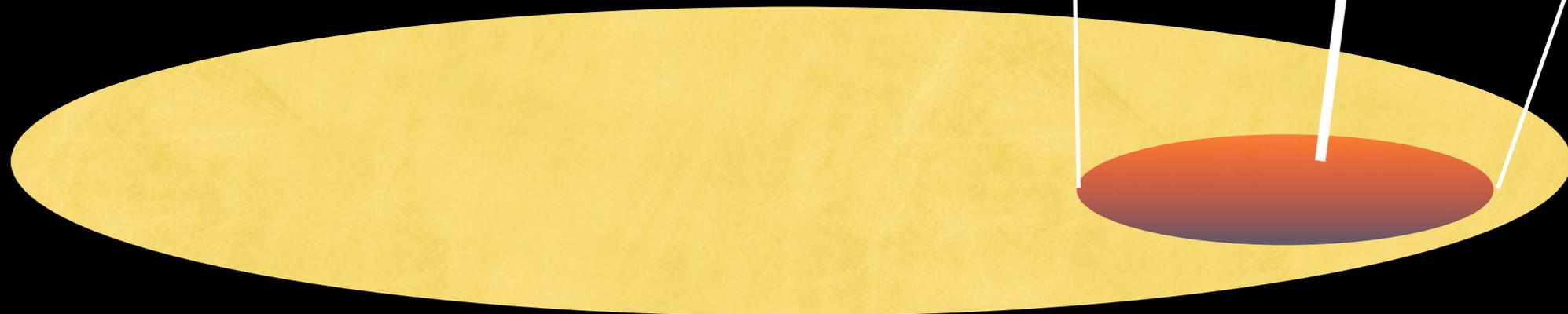
“Theory” space

Effective theories: philosophy

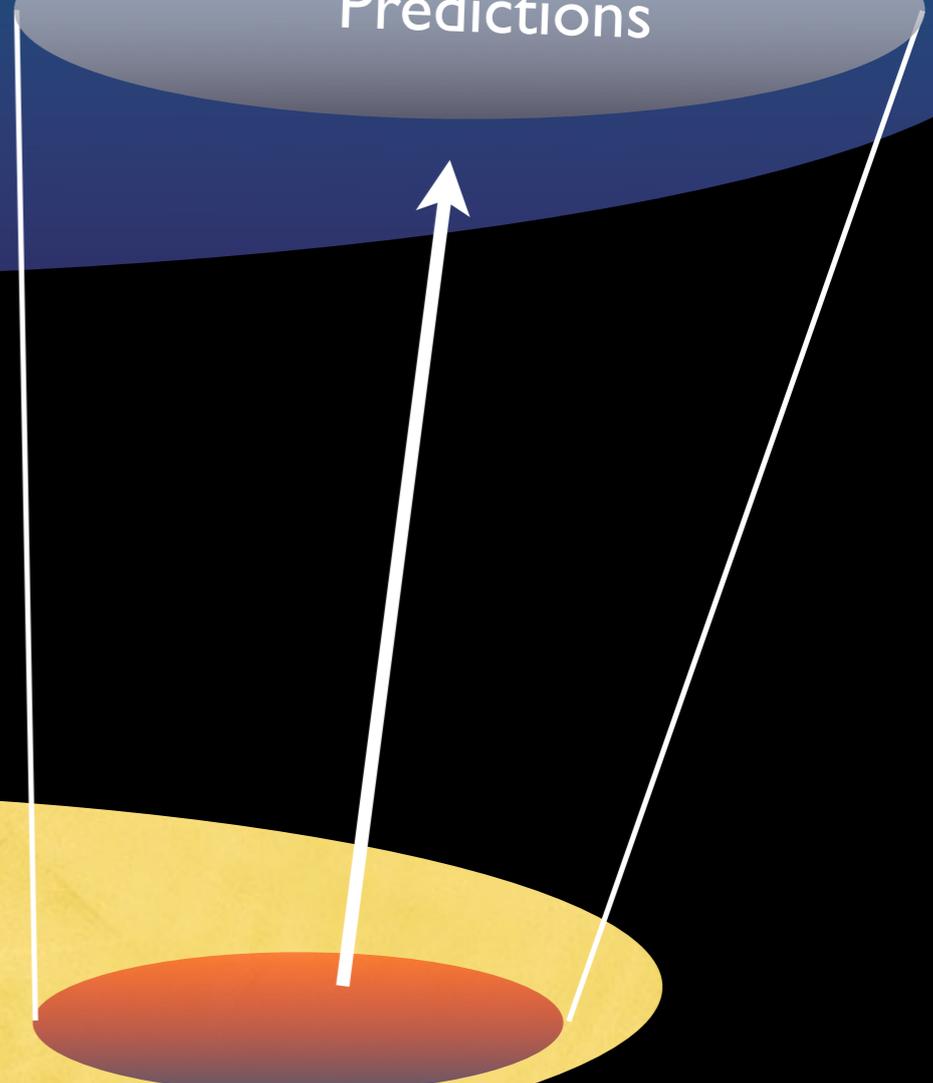
“Observables” space



Predictions

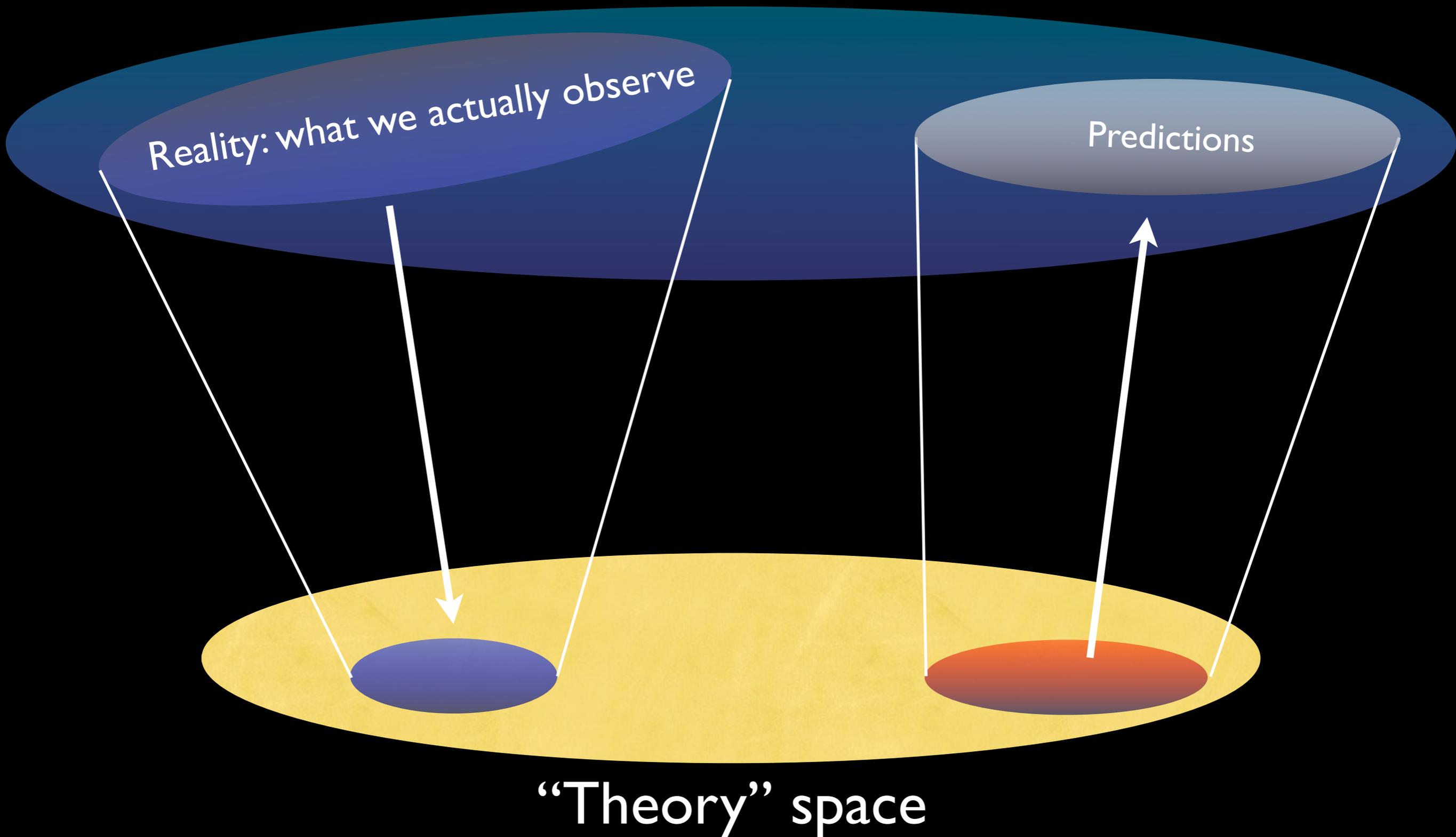


“Theory” space



Effective theories: philosophy

“Observables” space



Effective field theory

- * Some theories may be probing “theory space” which is already ruled out by observations... but you don’t know that until you write down the theory & do the calculations
- * Want to write down “*everything*” with all possible “free parameters” and find out what values they can take to be *consistent* with observational data
 - * *Particle physicists* do this... field content + symmetry: work out how to measure the free parameters (interaction terms, masses, etc...)
- * Doing this with a Lagrangian enables *clear physical interpretation* of constraints on parameters

Lagrangian for perturbations

The Lagrangian for perturbations is a *quadratic functional* in perturbations of fields...

$$\mathcal{L}_{\{2\}} = \sum_{A=1}^N \sum_{B=1}^N \mathbf{G}_{AB}^{\{0\}} \delta X^{(A)} \delta X^{(B)}$$

The Lagrangian for perturbations is equivalent to the *second measure-weighted variation* of the action...

$$\delta^2 S = \int d^4x \sqrt{-g} \underbrace{\left[\frac{1}{\sqrt{-g}} \delta^2 (\sqrt{-g} \mathcal{L}) \right]}_{\diamond^2 \mathcal{L}} = \int d^4x \sqrt{-g} \mathcal{L}_{\{2\}}$$

This allows us to

- (1) “dream up” a Lagrangian for perturbations,
- (2) explicitly calculate for some known theory... compare with established theories

Map from general parameterization to established theories,
but we are not limited by them!

Nothing extra

Field content: just the metric

Lagrangian for perturbations

$$\mathcal{L}_{\{2\}} = \frac{1}{8} \mathcal{W}^{\mu\nu\alpha\beta} \delta g_{\mu\nu} \delta g_{\alpha\beta}$$

Gravitational effects...

$$\delta U^{\mu\nu} = -\frac{1}{2} \left[\mathcal{W}^{\mu\nu\alpha\beta} + U^{\mu\nu} g^{\alpha\beta} \right] \delta g_{\alpha\beta}$$

Elastic dark energy, Feirz-Pauli & massive gravity in GR,...

Scalar fields

$$\mathcal{L} = \mathcal{L}(g_{\mu\nu}, \phi, \nabla_{\mu}\phi)$$

$$\mathcal{L}_{\{2\}} = \mathcal{L}_{\{2\}}(\delta g_{\mu\nu}, \delta\phi, \nabla_{\mu}\delta\phi)$$

Lagrangian for perturbations

$$\begin{aligned} \mathcal{L}_{\{2\}} = & \mathcal{A}\delta\phi\delta\phi + \mathcal{B}^{\mu}\delta\phi\nabla_{\mu}\delta\phi + \frac{1}{2}\mathcal{C}^{\mu\nu}\nabla_{\mu}\delta\phi\nabla_{\nu}\delta\phi \\ & + \frac{1}{4}\left[\mathcal{V}^{\mu\nu}\delta\phi\delta g_{\mu\nu} + \mathcal{Y}^{\alpha\mu\nu}\nabla_{\alpha}\delta\phi\delta g_{\mu\nu} + \frac{1}{2}\mathcal{W}^{\mu\nu\alpha\beta}\delta g_{\alpha\beta}\delta g_{\mu\nu}\right] \end{aligned}$$

Gravitational effects...

$$\delta U^{\mu\nu} = -\frac{1}{2}\left[\mathcal{V}^{\mu\nu}\delta\phi + \mathcal{Y}^{\alpha\mu\nu}\nabla_{\alpha}\delta\phi + \mathcal{W}^{\mu\nu\alpha\beta}\delta g_{\alpha\beta} + U^{\mu\nu}g^{\alpha\beta}\delta g_{\alpha\beta}\right]$$

Quintessence, k-essence, Lorentz violating theories,...

Total number of free functions

Scalar field case before imposing anything

Tensors in $L_{\{2\}}$	$\left\{ \mathcal{A}, \mathcal{B}^\mu, \mathcal{C}^{\mu\nu}, \mathcal{V}^{\mu\nu}, \mathcal{Y}^{\alpha\mu\nu}, \mathcal{W}^{\mu\nu\alpha\beta} \right\}$
# components	1 4 10 10 40 24

$$\mathcal{L} = \mathcal{L}(g_{\mu\nu}, \phi, \nabla_\mu \phi)$$

$$\begin{aligned} \mathcal{X} &\equiv -\frac{1}{2} g^{\mu\nu} \nabla_\mu \phi \nabla_\nu \phi \\ &= \frac{1}{2a^2} \dot{\phi}^2 \end{aligned}$$

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Total: 90

Impose *isotropy*: 90 \rightarrow 14

Impose *linking*: 14 \rightarrow 11

Impose $\mathcal{L} = \mathcal{L}(\phi, \mathcal{X})$: 11 \rightarrow 3

Imposing symmetry on the background

e.g. isotropy of spatial sections...

everything becomes compatible with FRW

$$(3+1) \text{ decomposition } g_{\mu\nu} = \gamma_{\mu\nu} - u_{\mu}u_{\nu} \quad \gamma_{\mu\nu}u^{\nu} = 0 \quad u_{\mu}u^{\mu} = -1$$

Coefficients in Lagrangian become:

$$\mathcal{W}_{\mu\nu\alpha\beta} = A_{\mathcal{W}}u_{\mu}u_{\nu}u_{\alpha}u_{\beta} + B_{\mathcal{W}}\left(\gamma_{\mu\nu}u_{\alpha}u_{\beta} + \gamma_{\alpha\beta}u_{\mu}u_{\nu}\right) \\ + 2C_{\mathcal{W}}\left(\gamma_{\mu(\alpha}u_{\beta)}u_{\nu} + \gamma_{\nu(\alpha}u_{\beta)}u_{\mu}\right) + \mathcal{E}_{\mu\nu\alpha\beta},$$

$$\mathcal{E}_{\mu\nu\alpha\beta} = D_{\mathcal{W}}\gamma_{\mu\nu}\gamma_{\alpha\beta} + 2E_{\mathcal{W}}\gamma_{\mu(\alpha}\gamma_{\beta)\nu}.$$

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$$A = A_{\mathcal{A}},$$

$$B^{\mu} = A_{\mathcal{B}}u^{\mu},$$

$$C_{\mu\nu} = A_{\mathcal{C}}u_{\mu}u_{\nu} + B_{\mathcal{C}}\gamma_{\mu\nu},$$

$$V_{\mu\nu} = A_{\mathcal{V}}u_{\mu}u_{\nu} + B_{\mathcal{V}}\gamma_{\mu\nu},$$

$$\mathcal{Y}_{\alpha\mu\nu} = A_{\mathcal{Y}}u_{\alpha}u_{\mu}u_{\nu} + B_{\mathcal{Y}}u_{\alpha}\gamma_{\mu\nu} + 2C_{\mathcal{Y}}\gamma_{\alpha(\mu}u_{\nu)}.$$

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$$B^{\mu} = A_{\mathcal{B}}u^{\mu},$$

$$C_{\mu\nu} = A_{\mathcal{C}}u_{\mu}u_{\nu} + B_{\mathcal{C}}\gamma_{\mu\nu},$$

$$V_{\mu\nu} = A_{\mathcal{V}}u_{\mu}u_{\nu} + B_{\mathcal{V}}\gamma_{\mu\nu},$$

Impose some "theory" structure:

$$\mathcal{L} = \mathcal{L}(\phi, \mathcal{X}) \quad \text{reduce } 14 \rightarrow 3$$

$$\mathcal{Y}_{\alpha\mu\nu} = A_{\mathcal{Y}}u_{\alpha}u_{\mu}u_{\nu} + B_{\mathcal{Y}}u_{\alpha}\gamma_{\mu\nu} + 2C_{\mathcal{Y}}\gamma_{\alpha(\mu}u_{\nu)}.$$

Function	(a) EDE	(b) $\mathcal{L} = \mathcal{L}(\phi, \mathcal{X})$	(c) $\mathcal{L} = F(\mathcal{X})$	(d) $\mathcal{L} = \mathcal{X} - V(\phi)$
$A_{\mathcal{Y}}$	0	$-2(\mathcal{L}_{,\mathcal{X}\phi}\dot{\phi}^2 - \mathcal{L}_{,\phi})$	0	$-2V'$
$B_{\mathcal{Y}}$	0	$-2\mathcal{L}_{,\phi}$	0	$2V'$
$A_{\mathcal{Y}}$	0	$-2(\mathcal{L}_{,\mathcal{X}\mathcal{X}}\dot{\phi}^2 + \mathcal{L}_{,\mathcal{X}}\dot{\phi})$	$-2(F''\dot{\phi}^2 + F'\dot{\phi})$	$-2\dot{\phi}$
$B_{\mathcal{Y}}$	0	$-2\mathcal{L}_{,\mathcal{X}}\dot{\phi}$	$-2F'\dot{\phi}$	$-2\dot{\phi}$
$C_{\mathcal{Y}}$	0	$2\mathcal{L}_{,\mathcal{X}}\dot{\phi}$	$2F'\dot{\phi}$	$2\dot{\phi}$
$A_{\mathcal{W}}$	$-\rho$	$-(\mathcal{L}_{,\mathcal{X}\mathcal{X}}\dot{\phi}^4 + 2\rho + P)$	$-(F''\dot{\phi}^4 + 2\rho + P)$	$-(2\rho + P)$
$B_{\mathcal{W}}$	P	$-\rho$	$-\rho$	$-\rho$
$C_{\mathcal{W}}$	$-P$	ρ	ρ	ρ
$D_{\mathcal{W}}$	$\beta - P - \frac{2}{3}\mu$	$-P$	$-P$	$-P$
$E_{\mathcal{W}}$	$\mu + P$	P	P	P

Parameterizing entropy

$$\delta P = w\delta\rho + P\Gamma$$

$$w\Gamma = (\alpha - w) \left[\delta - 3\mathcal{H}\beta(1+w)\theta \right]$$

Standard: Weller & Lewis: $\alpha = c_s^2$ $\beta = 1$

Nothing extra $\beta = 0$

Quintessence $\alpha = 1$ $\beta = 1$

Pure k-essence $\beta = 0$

Physically, what does $\alpha \neq 1$ mean?

Is α always a sound speed?

α is neither group nor phase velocity of waves

$$\alpha = \left(1 + 2\mathcal{X} \frac{\mathcal{L}_{,\mathcal{X}\mathcal{X}}}{\mathcal{L}_{,\mathcal{X}}} \right)^{-1},$$

$$\beta = \frac{2a\mathcal{L}_{,\phi}}{3\mathcal{H}\mathcal{L}_{,\mathcal{X}}\sqrt{2\mathcal{X}}} \left[1 + \mathcal{X} \left(\frac{\mathcal{L}_{,\mathcal{X}\mathcal{X}}}{\mathcal{L}_{,\mathcal{X}}} - \frac{\mathcal{L}_{,\mathcal{X}\phi}}{\mathcal{L}_{,\phi}} \right) \right] \frac{\alpha}{\alpha - w}.$$

Effective metric that scalar field perturbations “feel” ...

$$C^{\mu\nu} = \mathcal{L}_{,\mathcal{X}}g^{\mu\nu} + \mathcal{L}_{,\mathcal{X}}\alpha^{-1}(\alpha - 1)u^\mu u^\nu$$

Scope of using the Lagrangian for perturbations

❖ Formalism

- ❖ The Lagrangian for perturbations
- ❖ The perturbed dark energy-momentum tensor
- ❖ Lagrangian & Eulerian perturbations (Stuckleberg completion / deformation vector)

❖ Examples

- ❖ Nothing extra
- ❖ Scalar & vector fields (*aether*, *TeV*S, ...)
- ❖ “high order” derivative theories (*galileon*, $F(R)$, *Gauss-Bonnet*, ...)

❖ Applications

- ❖ Cosmological perturbations
- ❖ Entropy & anisotropic stresses
- ❖ Massive gravity
- ❖ Modified gravity

Summary

- ❖ Construct coherent consistent modifications to the gravitational field equations at perturbed order
- ❖ All freedom inside “background” tensors
- ❖ Encompass theories never before considered: $\mathcal{L}_{\{2\}}$ needs scalars in background *and* perturbed field variables: more freedom!
- ❖ Encompass massive gravity & “high derivative” theories: e.g. galileon, Horndeski, Brans-Dicke.
- ❖ In a model independent way compute cosmological observables (CMB, lensing, $P(k)$, ...): rule in / out before requiring the actual theory!
- ❖ Impose symmetry on background (e.g. isotropy... compatible with FRW)... allows split of BG tensors

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