

Extracting foreground-obscured μ -type distortion anisotropies with future CMB satellites

Mathieu Remazeilles



The University of Manchester

Remazeilles & Chluba
[MNRAS 478, 807 \(2018\)](#)

15th Marcel Grossmann Meeting
Rome, 1-7 July 2018

μ -type CMB spectral distortions

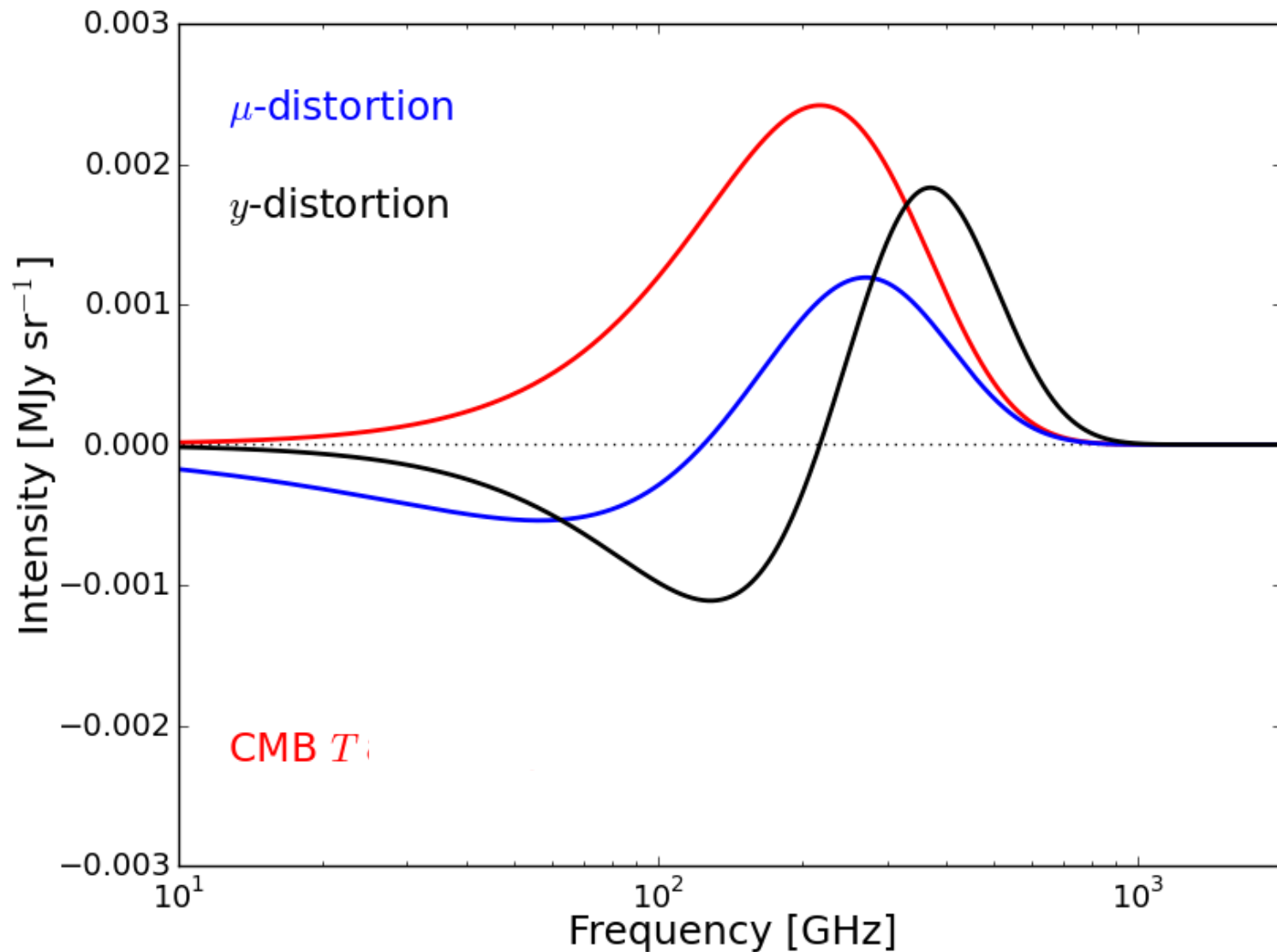
Sunyaev &
Zeldovich 1970

- At pre-recombination era ($10^4 < z < 2 \times 10^6$), energy injections into primordial plasma prevent *brehmsstrahlung* and *double Compton scattering* to create photons to maintain Planck's equilibrium, leading to Bose-Einstein (BE) equilibrium:

$$n_{BE} = n_{Pl} + \underbrace{\mu \frac{e^x}{(e^x - 1)^2} \left(\frac{x}{2.19} - 1 \right)}_{\text{spectral signature of } \mu\text{-distortion}} \quad x \equiv h\nu/kT_{CMB}$$

- Caused by exciting physical processes at redshifts $z > 10^4$:
 - ✓ Dissipation of primordial small-scale acoustic modes – Silk 1968
 - ✓ Annihilation / decay of relic particles – Hu & Silk 1993
 - ✓ Evaporation of primordial black holes – Carr et al 2010
 - ✓ Inflation with non-Bunch-Davies vacuum – Ganc & Komatsu 2012
- LCDM: $|\mu| = 2.3 \times 10^{-8}$ – Chluba 2016
- COBE/FIRAS: $|\mu| < 9 \times 10^{-5}$ – Fixsen et al 1996

Spectral signature of distortions



Distinct spectral signatures!

→ **Multi-frequency observations allow to disentangle those signals**

Anisotropic μ -type distortions

Aside from average CMB monopole distortions:

- Anisotropies of μ -type distortions (spectral-spatial anisotropies)

$$C_{\ell}^{\mu\mu} = 144 C_{\ell}^{TT,SW} f_{NL}^2 \langle \mu \rangle^2$$

- Correlation between CMB temperature and μ -type distortion anisotropies:

$$C_{\ell}^{\mu \times T} = 12 C_{\ell}^{TT,SW} \rho(\ell) f_{NL} \langle \mu \rangle$$

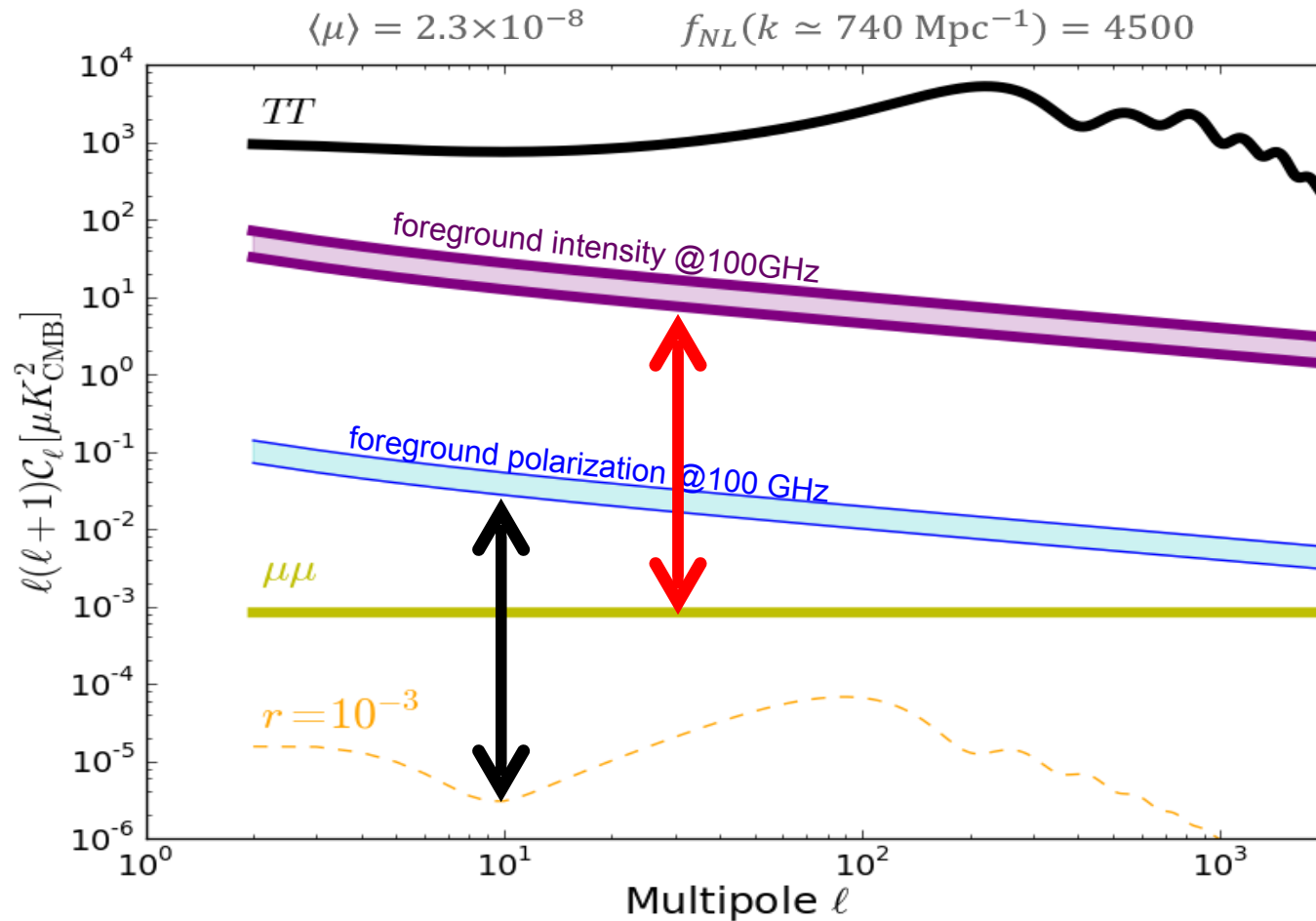
- Sources:

- ✓ *Inflation with non-Bunch-davies initial state* – [Ganc & Komatsu 2012](#)
- ✓ *Damping of primordial small-scale acoustic modes in the presence of enhanced primordial non-Gaussianity in the ultra-squeezed limit*
[Pajer & Zaldarriaga 2012](#)

N.B.: Scale-dependent non-Gaussianity $f_{NL}(k) = f_{NL}(k_0) \left(\frac{k}{k_0}\right)^{n_{NL}-1}$ with $n_{NL} \simeq 1.6$

allows for:
$$\begin{cases} f_{NL}(k \simeq 740 \text{ Mpc}^{-1}) \simeq 4500 & (\mu\text{-distortion probe}) \\ f_{NL}(k_0 = 0.05 \text{ Mpc}^{-1}) \simeq 5 & (\text{CMB probe}) \end{cases}$$

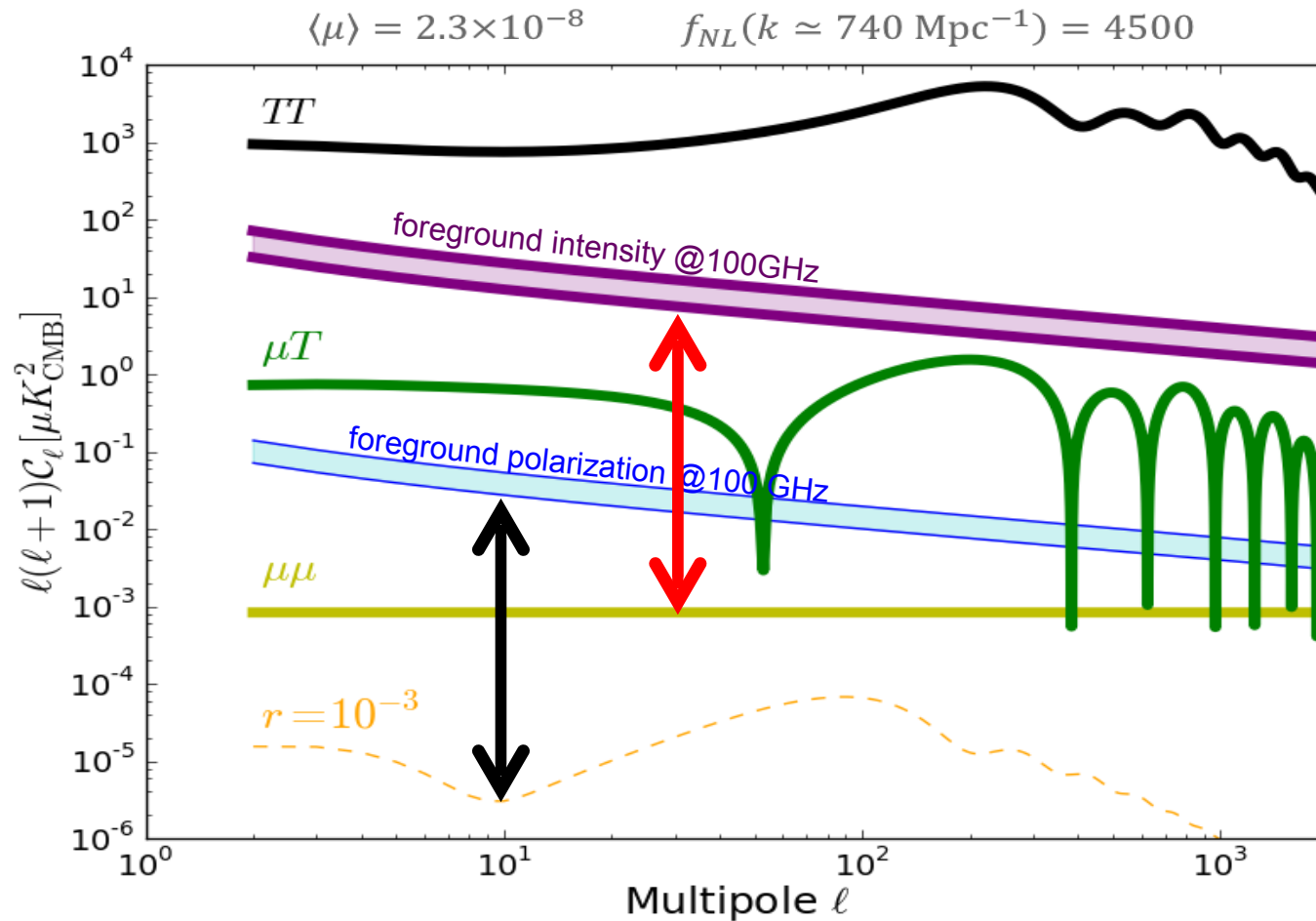
μ -type distortion anisotropies



**Same dynamic range (signal-to-foregrounds ratio)
than for B-modes at $r = 10^{-3}$**

→ A science case for future CMB satellites!

μ -type distortion anisotropies



μ -T cross-power spectrum:

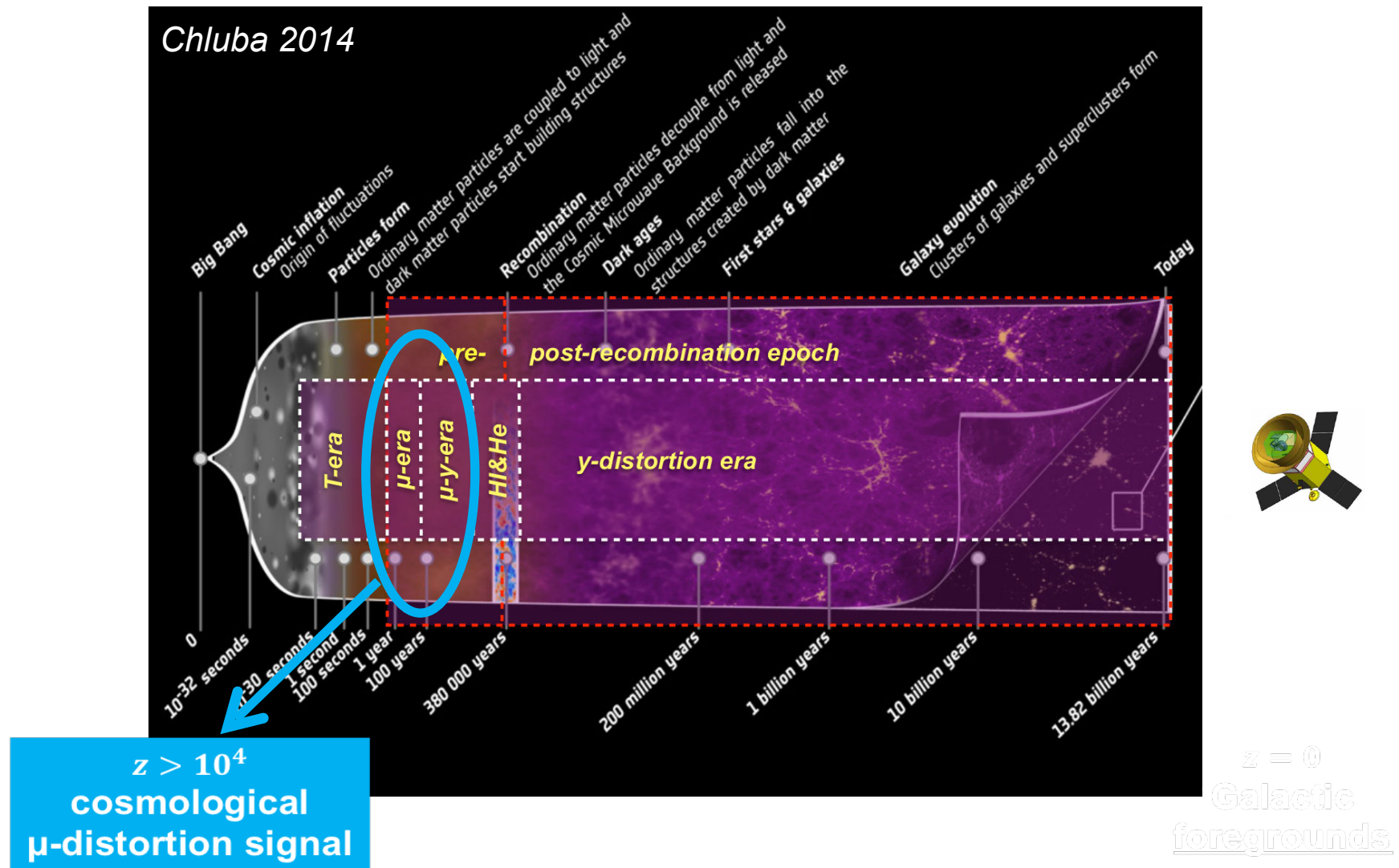
Enhanced μ -type distortion signal through correlation with CMB temperature anisotropies!

→ **Accessible signal for future CMB satellites!**

Questions

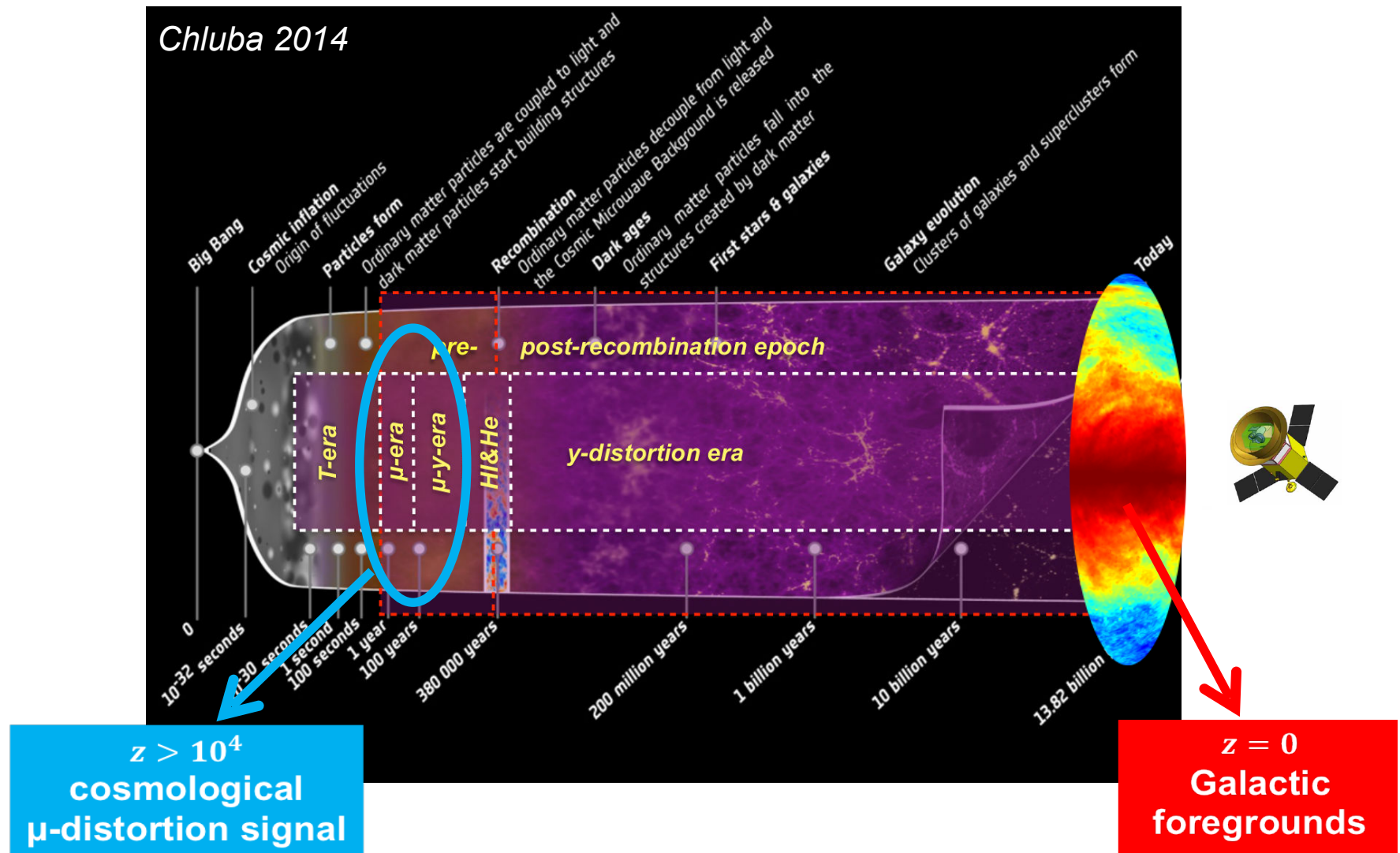
- *Can we detect the μ - T correlated signal with future CMB satellites?*
- *What constraints on $f_{NL}(k \simeq 740 \text{ Mpc}^{-1})$ can be achieved in the presence of foregrounds?*

The problem of foregrounds



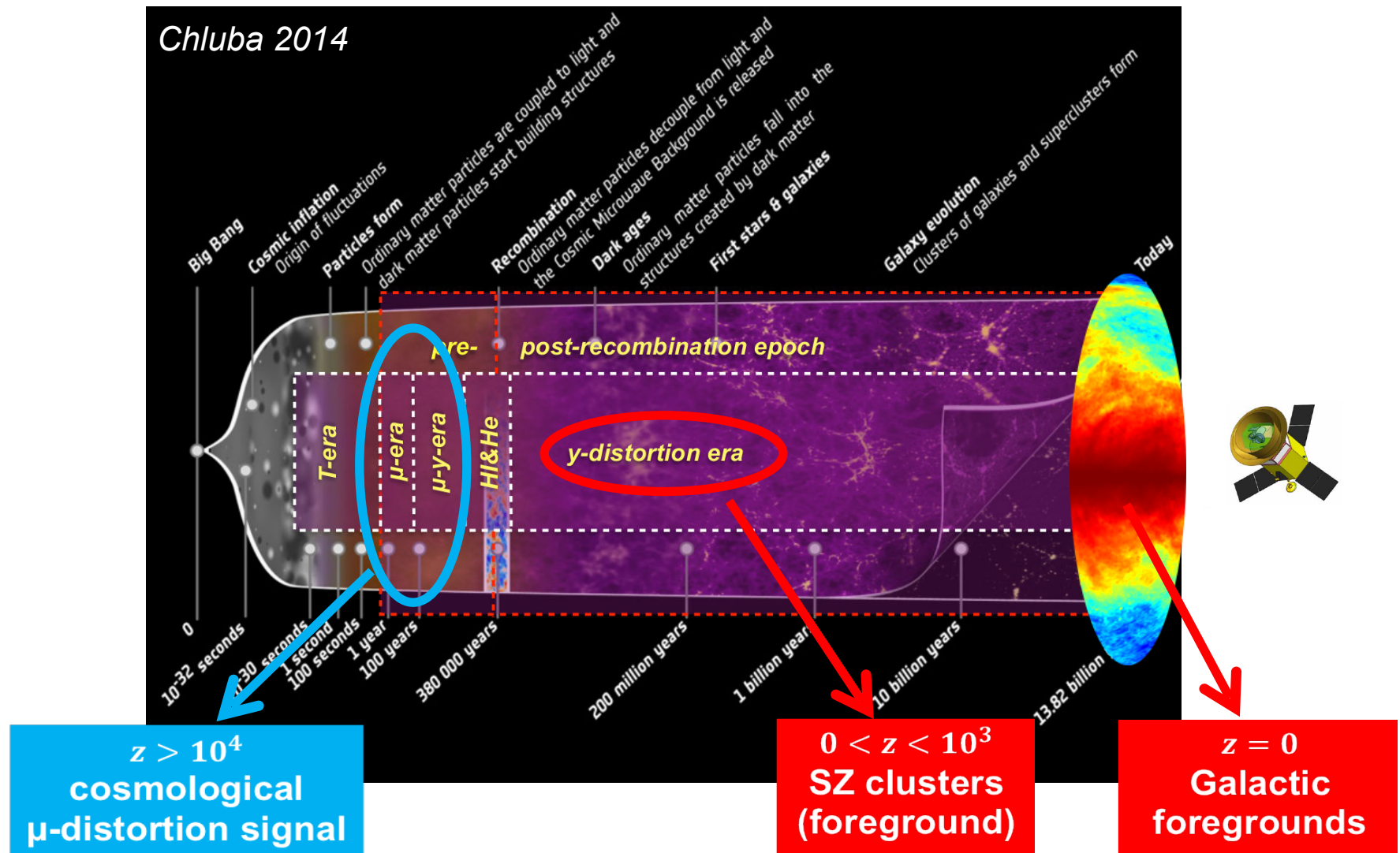
μ -type spectral distortions open a new window to probe physics occurring behind the last-scattering surface, where the universe is invisible!

The problem of foregrounds



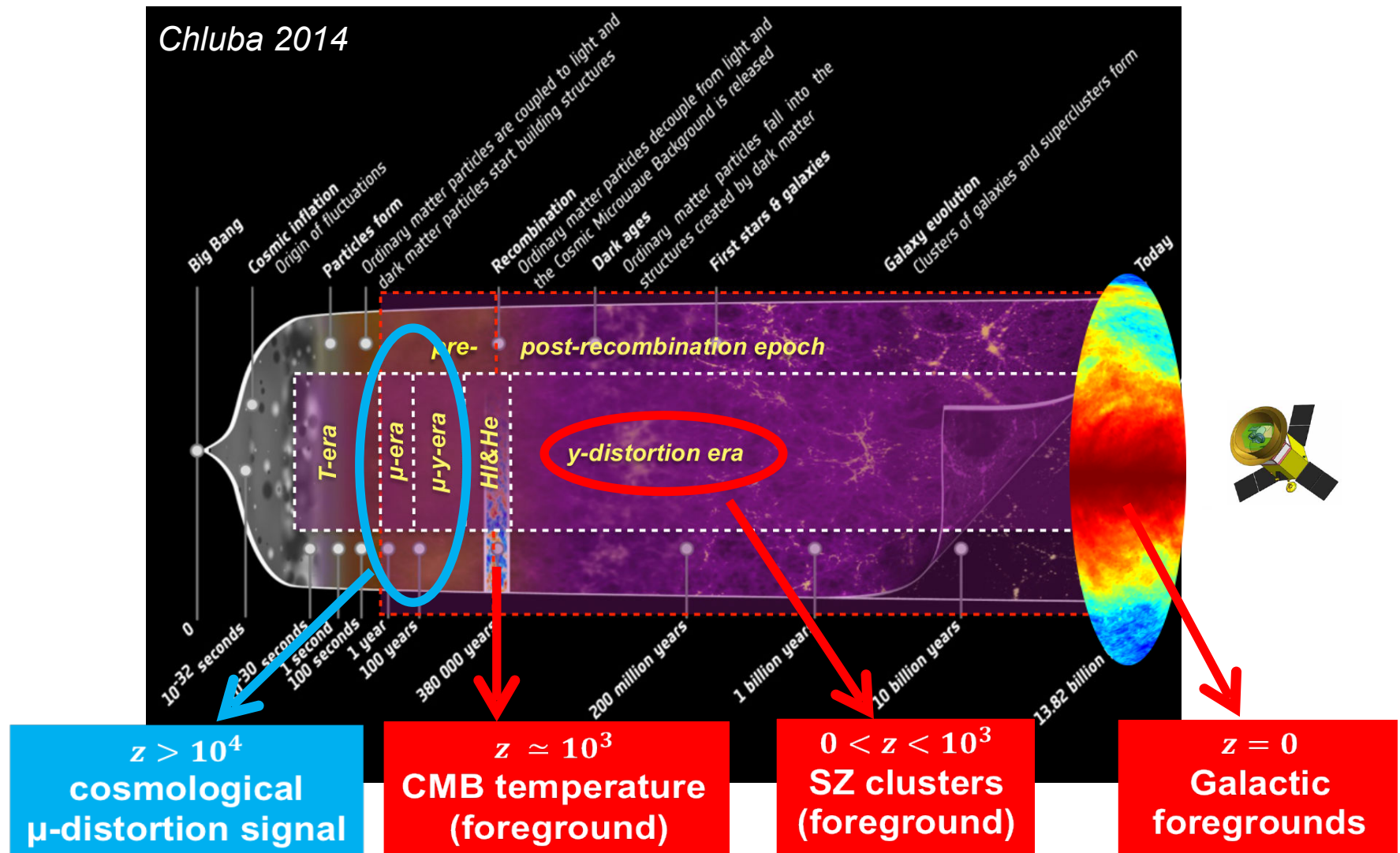
μ -type spectral distortions open a new window to probe physics occurring behind the last-scattering surface, where the universe is invisible!

The problem of foregrounds



μ -type spectral distortions open a new window to probe physics occurring behind the last-scattering surface, where the universe is invisible!

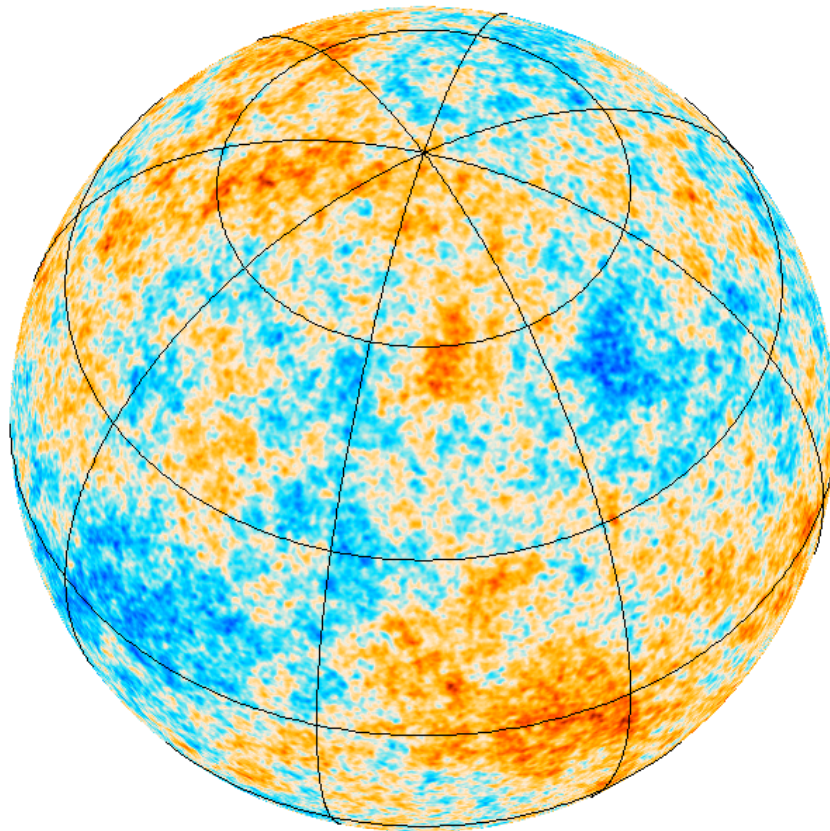
The problem of foregrounds



μ -type spectral distortions open a new window to probe physics occurring behind the last-scattering surface, where the universe is invisible!

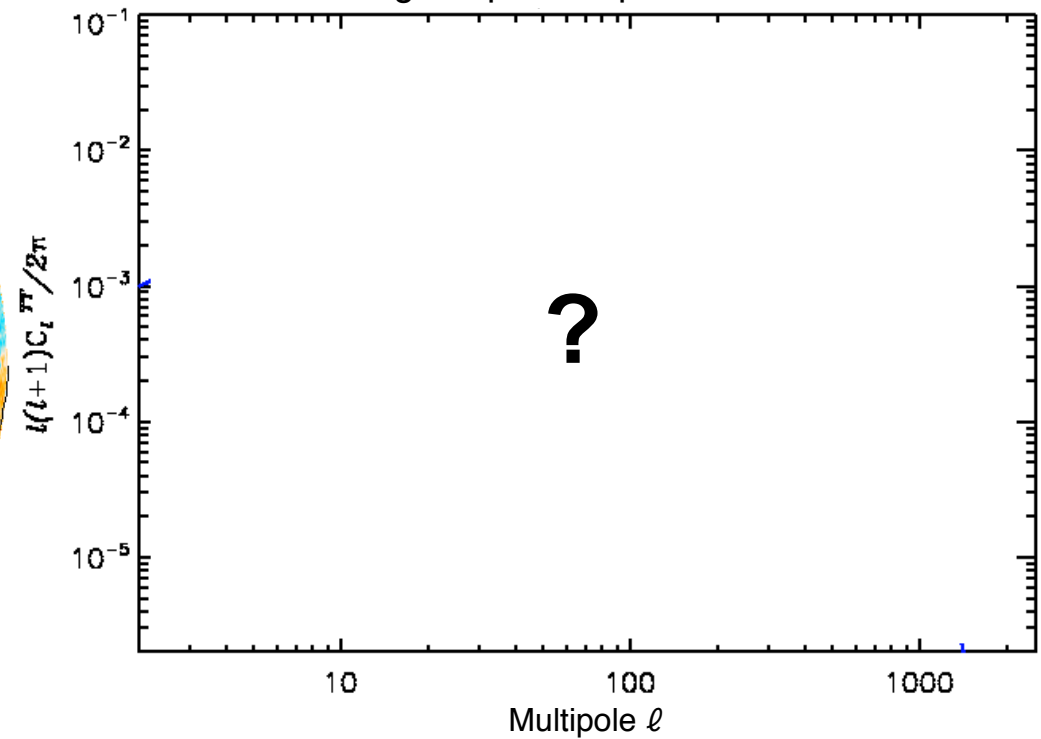
Component separation : the problem

μ -distortion anisotropies



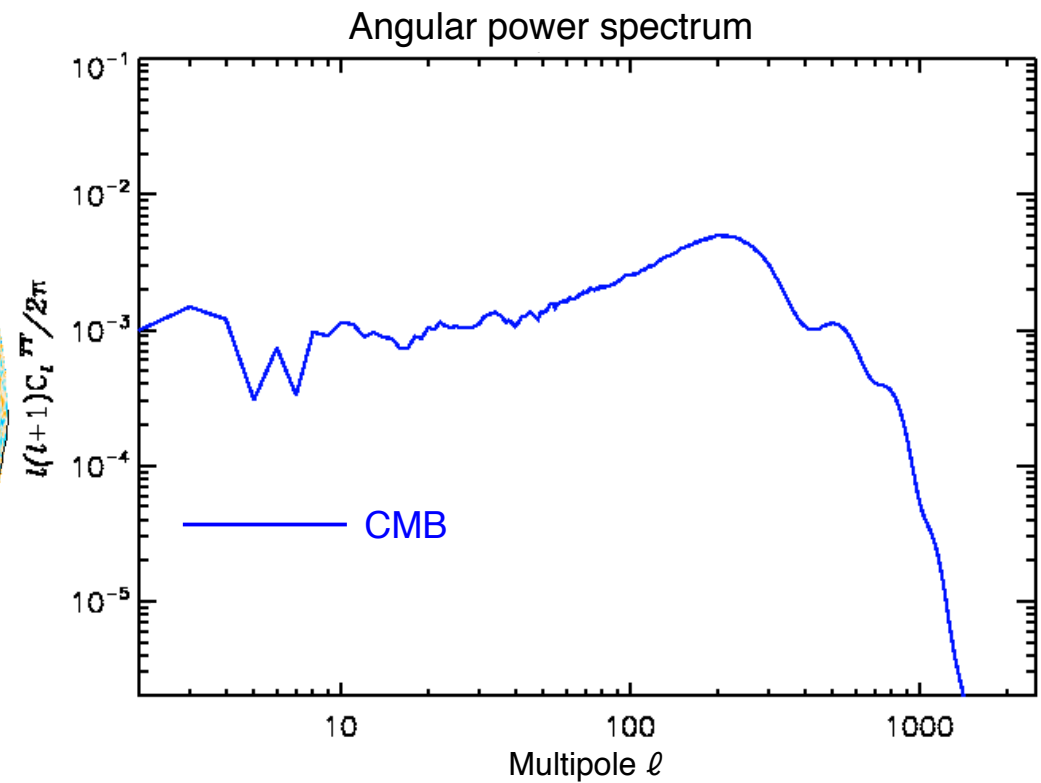
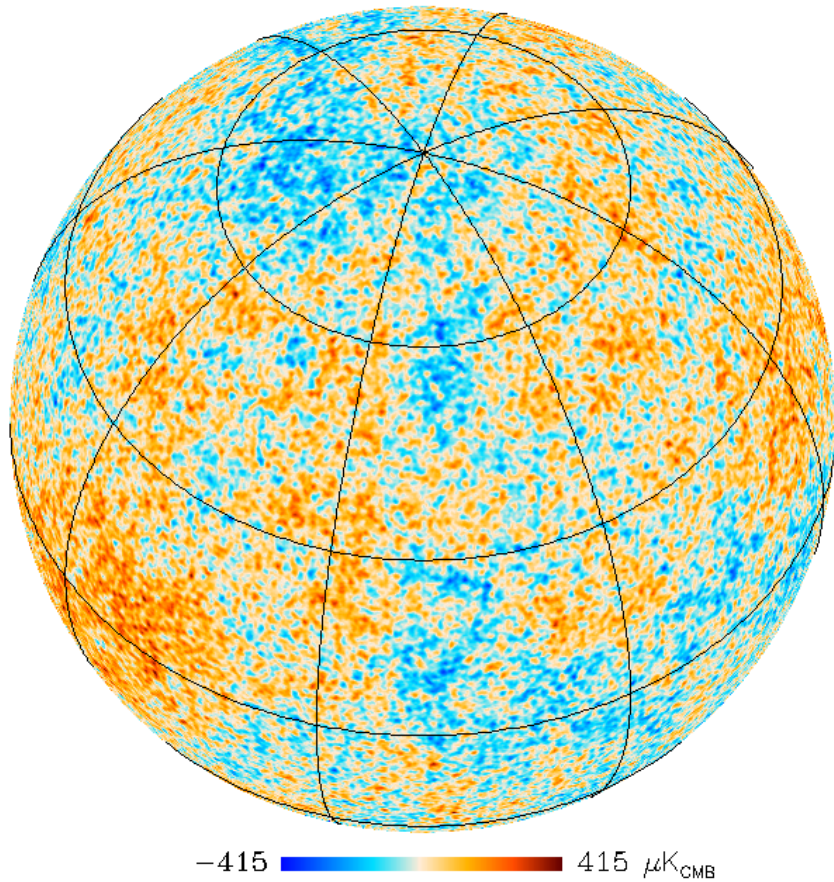
-0.015  0.015 μK_{CMB}

Angular power spectrum



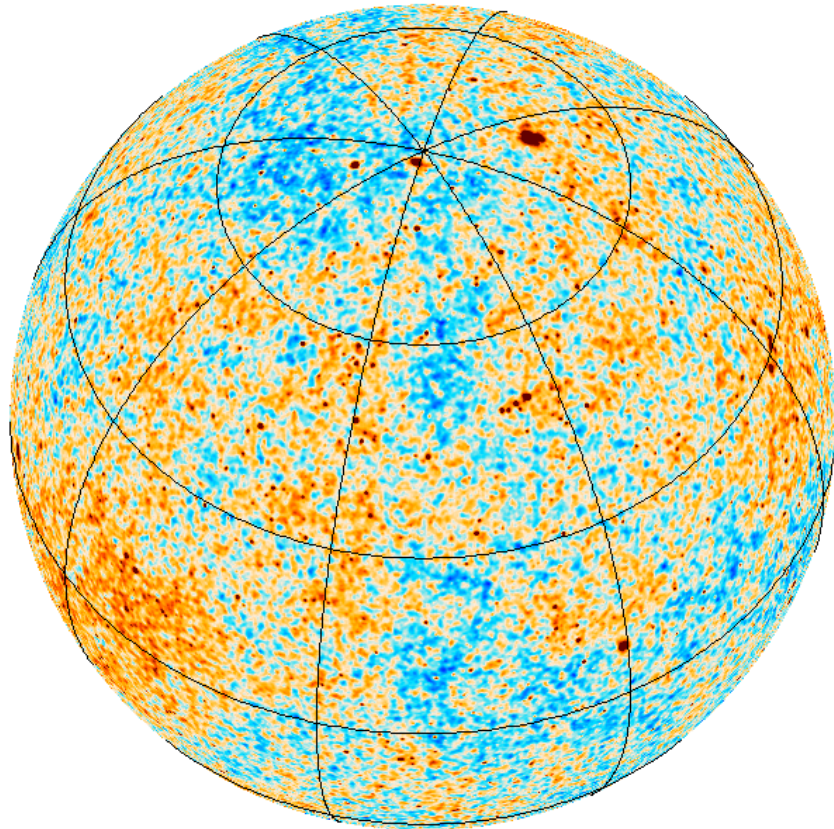
Component separation : the problem

μ -distortion + CMB temperature anisotropies

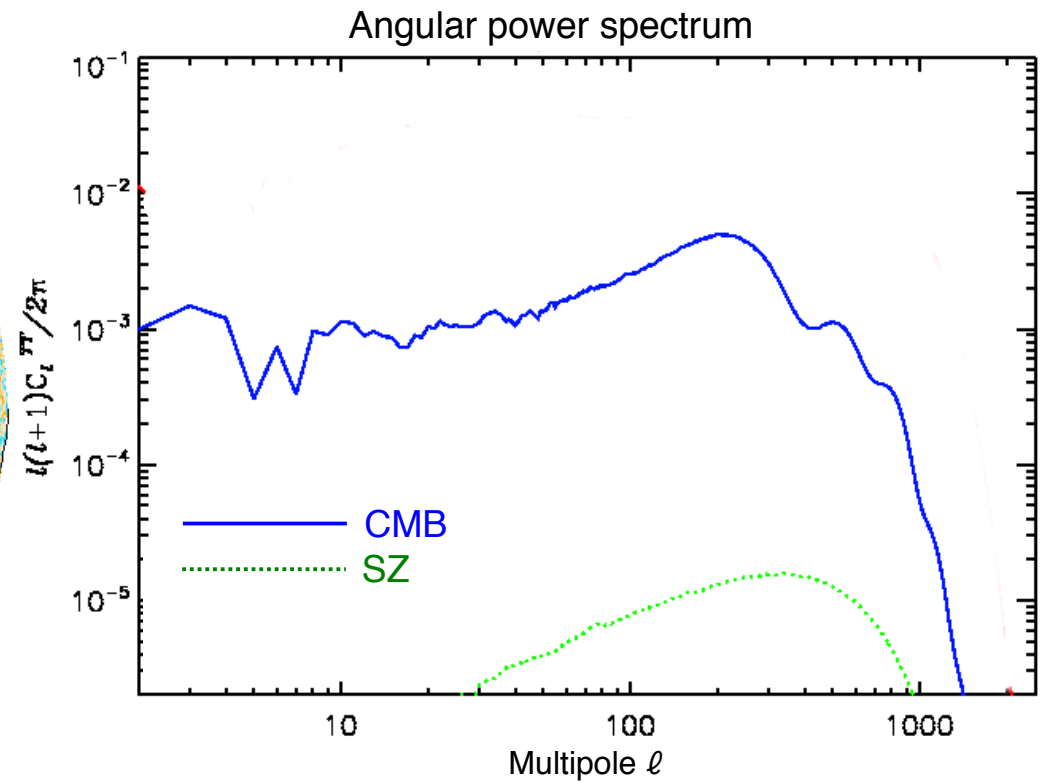


Component separation : the problem

μ -distortion + CMB + SZ

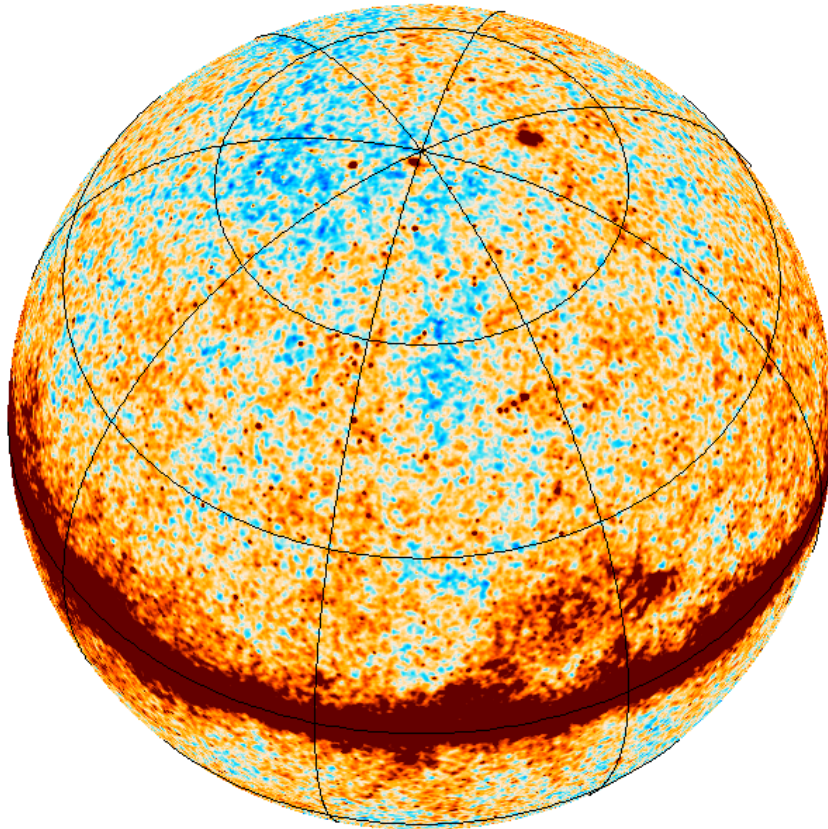


-415 415 μK_{CMB}

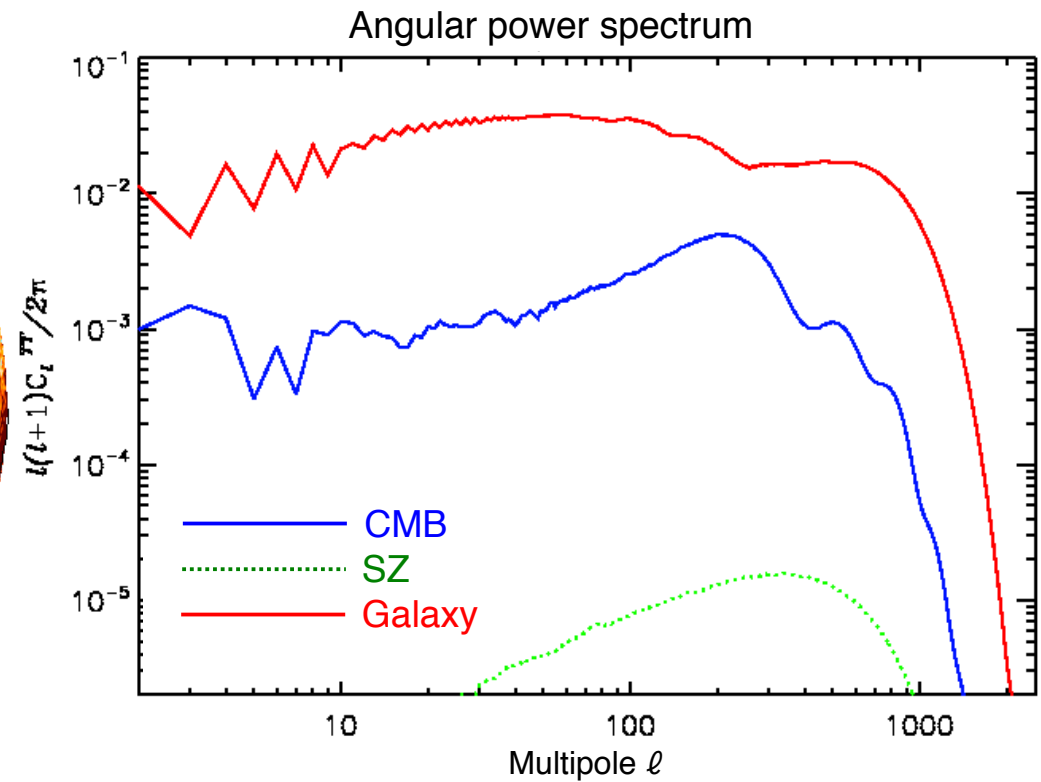


Component separation : the problem

μ -distortion + CMB + SZ + Galactic

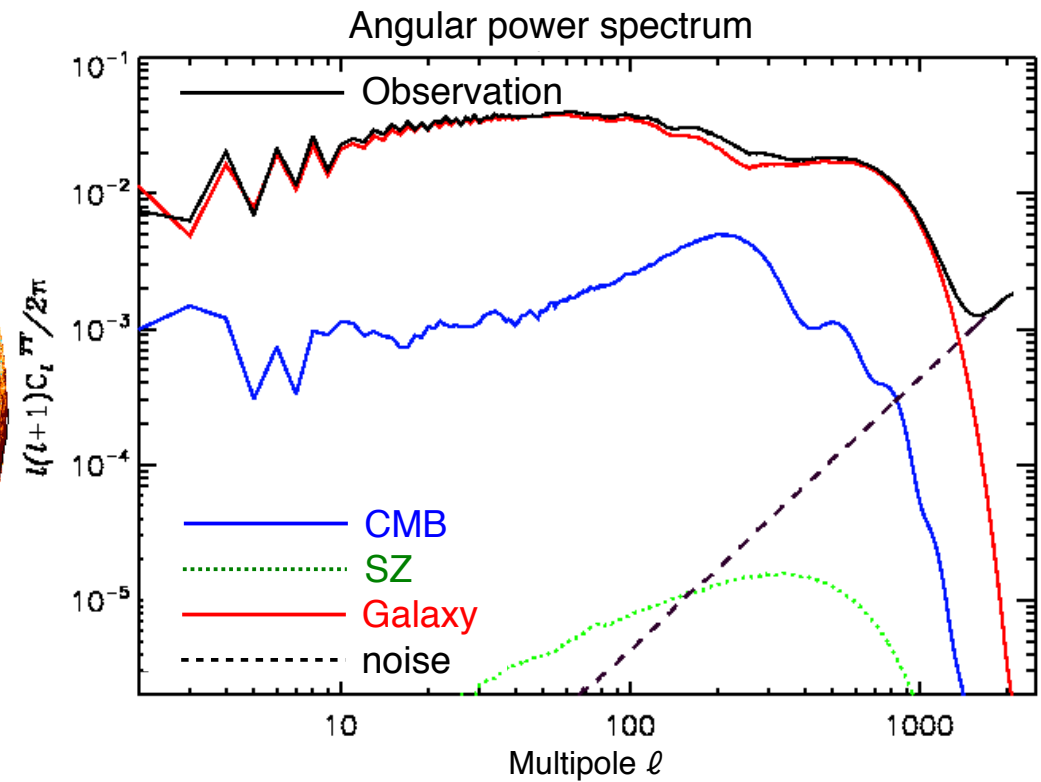
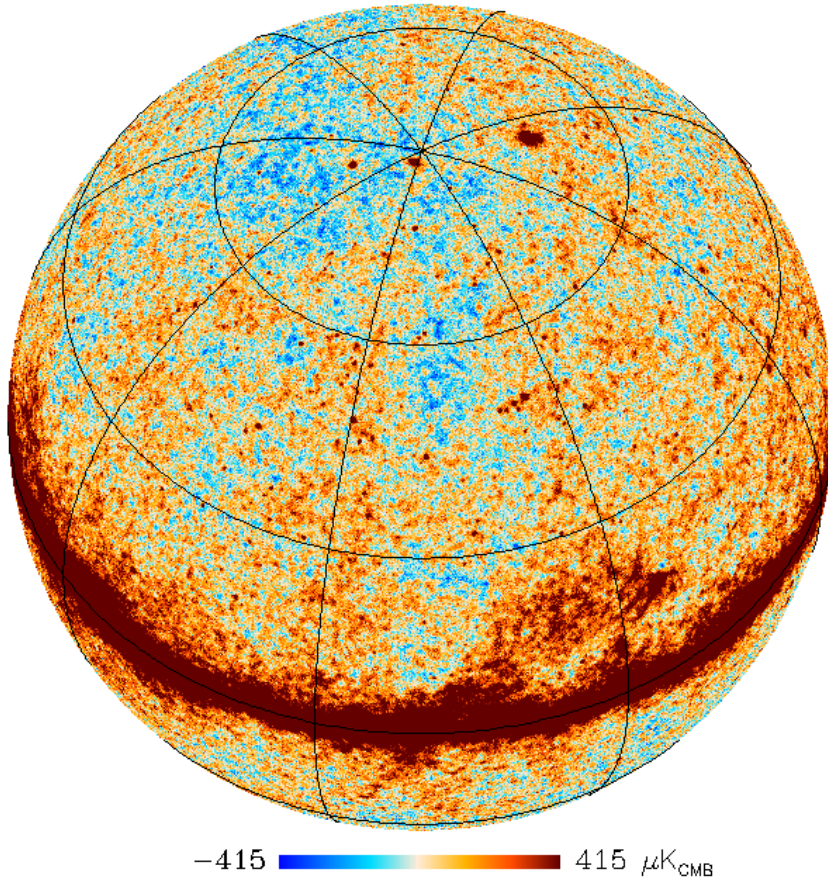


-415  415 μK_{CMB}



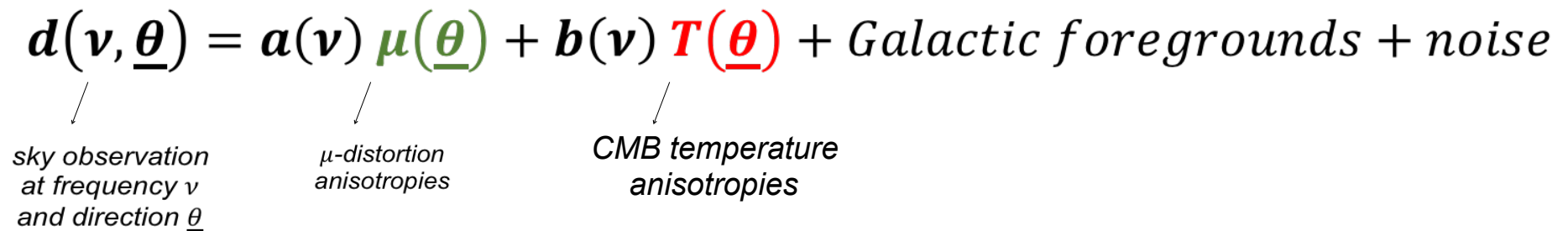
Component separation : the problem

μ -distortion + CMB + SZ + Galactic + noise



The problem of the CMB temp. foreground

$$d(\nu, \underline{\theta}) = a(\nu) \mu(\underline{\theta}) + b(\nu) T(\underline{\theta}) + \text{Galactic foregrounds} + \text{noise}$$


sky observation at frequency ν and direction $\underline{\theta}$ μ -distortion anisotropies CMB temperature anisotropies

- CMB T anisotropies: a major foreground to μ -type distortion anisotropies
- CMB T anisotropies: a foreground that is correlated to the μ -anisotropy signal!

→ *This might be an issue for component separation*

Component separation: standard ILC

$$d(\nu, \underline{\theta}) = a(\nu) \mu(\underline{\theta}) + b(\nu) T(\underline{\theta}) + \text{Galactic foregrounds} + \text{noise}$$

- Weighted linear combination of frequency maps:

$$\hat{\mu}(\underline{\theta}) = \sum_{\nu} w(\nu) d(\nu, \underline{\theta}) \quad \text{such that} \quad \begin{cases} \mathbf{w}^t \langle d d^t \rangle \mathbf{w} \text{ min. var.} \\ \sum_{\nu} w(\nu) a(\nu) = 1 \end{cases}$$

- Analytic solution: $\mathbf{w}^t = \frac{a^t C^{-1}}{a^t C^{-1} a} \quad (C \equiv \langle d d^t \rangle)$
[Benett et al, 2003](#)
[Tegmark et al, 2003](#)
[Eriksen et al, 2004](#)
[Delabrouille et al, 2009](#)
- Problem:

→ residual CMB T anisotropies in the ILC reconstructed μ -map

$$\hat{\mu}(\underline{\theta}) = \mu(\underline{\theta}) + (\mathbf{w}^t \mathbf{b}) T(\underline{\theta}) + \dots$$

→ residual TT correlations in the measured μ - T cross-power spectrum!

$$\langle \hat{\mu}(\underline{\theta}) T(\underline{\theta}') \rangle = \underbrace{\langle \mu(\underline{\theta}) T(\underline{\theta}') \rangle}_{\text{signal}} + \underbrace{\varepsilon \langle T(\underline{\theta}) T(\underline{\theta}') \rangle}_{\substack{TT \text{ residuals} \\ \text{bury the signal!}}} + \dots$$

Component separation: Constrained-ILC

$$d(\nu, \underline{\theta}) = a(\nu) \mu(\underline{\theta}) + b(\nu) T(\underline{\theta}) + \text{Galactic foregrounds} + \text{noise}$$

- Weighted linear combination of frequency maps:

$$\hat{\mu}(\underline{\theta}) = \sum_{\nu} w(\nu) d(\nu, \underline{\theta}) \quad \text{such that} \quad \begin{cases} w^t \langle d d^t \rangle w \text{ min. var.} \\ \sum_{\nu} w(\nu) a(\nu) = 1 \\ \sum_{\nu} w(\nu) b(\nu) = 0 \rightarrow \text{orthogonality to CMB } T \text{ spectrum} \end{cases}$$

- Analytic solution: $w^t = \frac{(b^t C^{-1} b) a^t C^{-1} - (a^t C^{-1} b) b^t C^{-1}}{(a^t C^{-1} a)(b^t C^{-1} b) - (a^t C^{-1} b)^2}$ Remazeilles et al, 2011
Remazeilles & Chluba, 2018

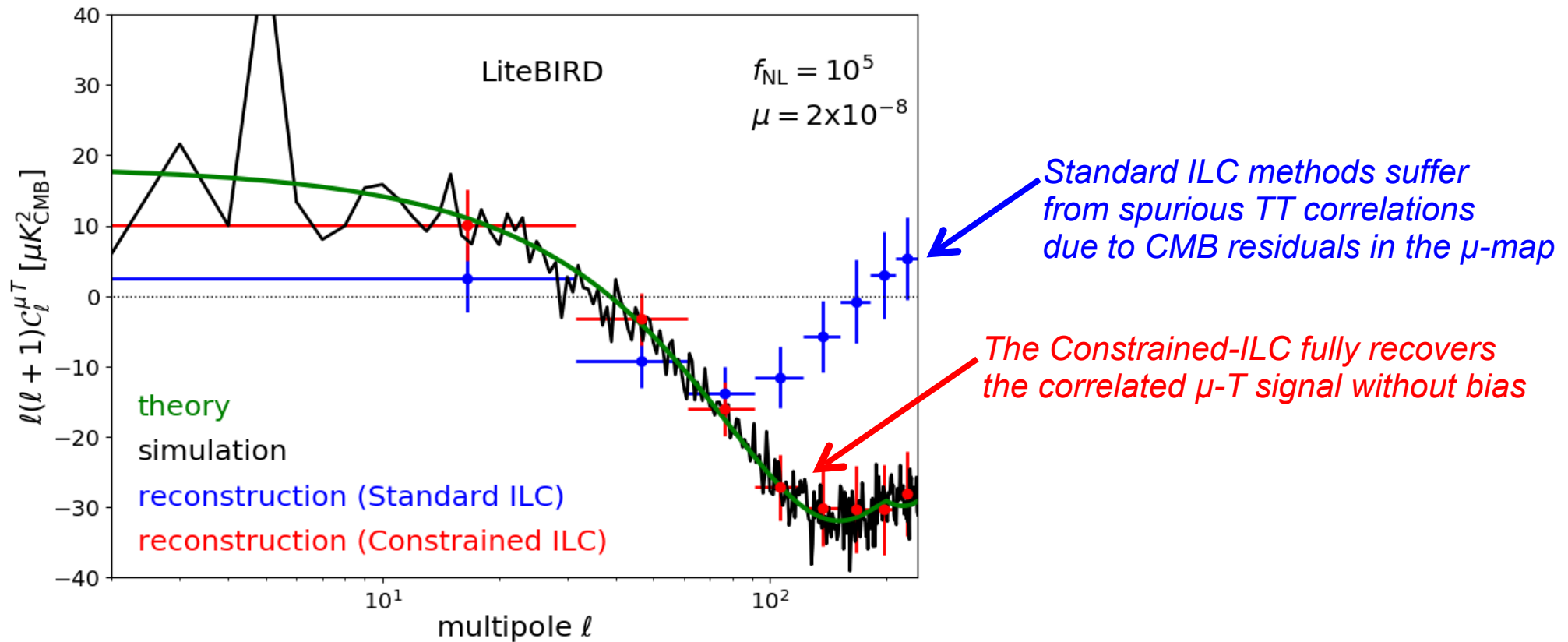
→ No more residual CMB T anisotropies in the reconstructed μ -map

$$\hat{\mu}(\underline{\theta}) = \mu(\underline{\theta}) + \underbrace{(w^t b)}_{=0} T(\underline{\theta}) + \dots$$

→ No more residual TT correlations in the μ - T cross-power spectrum!

$$\langle \hat{\mu}(\underline{\theta}) T(\underline{\theta}') \rangle = \langle \mu(\underline{\theta}) T(\underline{\theta}') \rangle + \varepsilon \langle \cancel{T(\underline{\theta}) T(\underline{\theta}')} \rangle + \dots$$

Standard ILC vs Constrained ILC

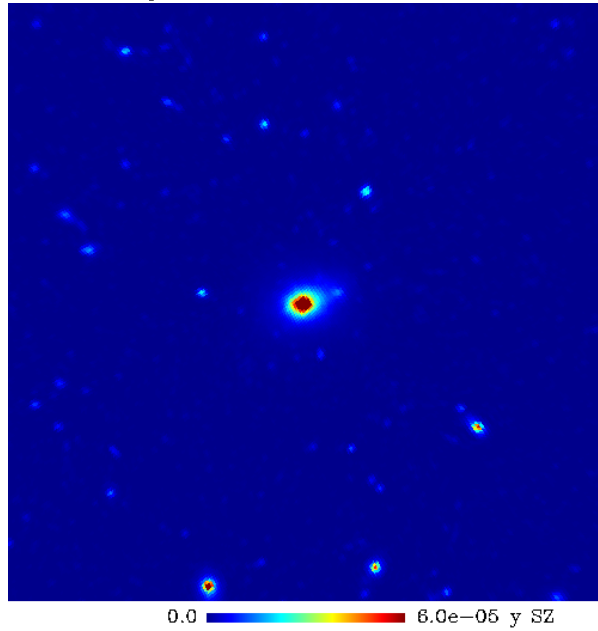


In light of these considerations, the constraints on $C_\ell^{\mu T}$ from Planck data by Khatri & Sunyaev (2015) should be taken cautiously

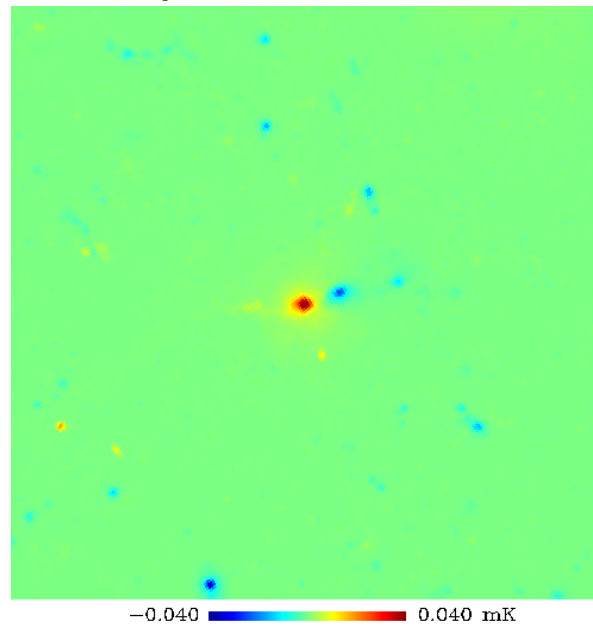
*The Constrained-ILC idea
can also be used to kill residual y-distortions
in CMB temperature map*

Standard ILC

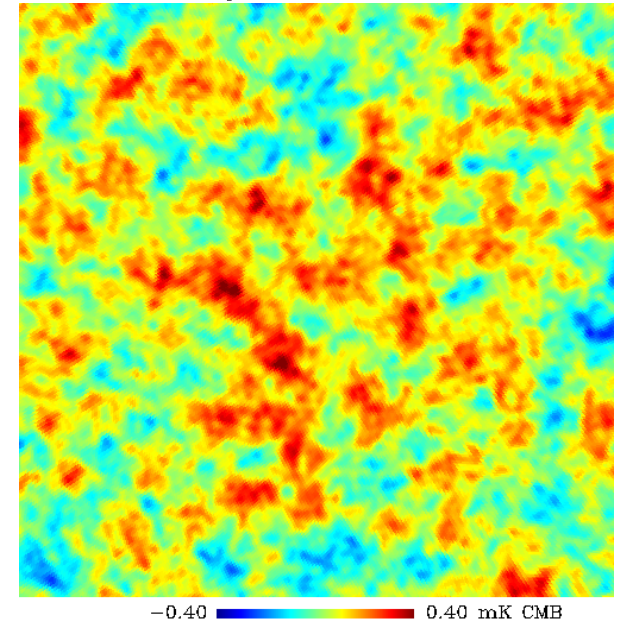
input thermal SZ



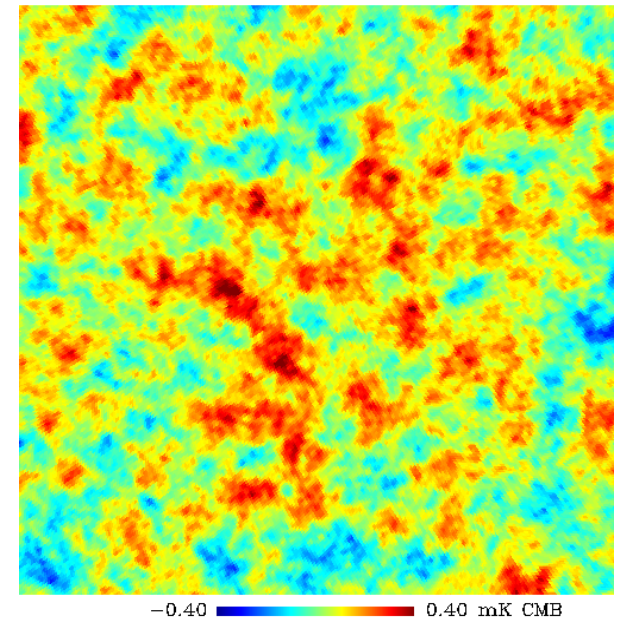
input kinetic SZ



input CMB



ILC

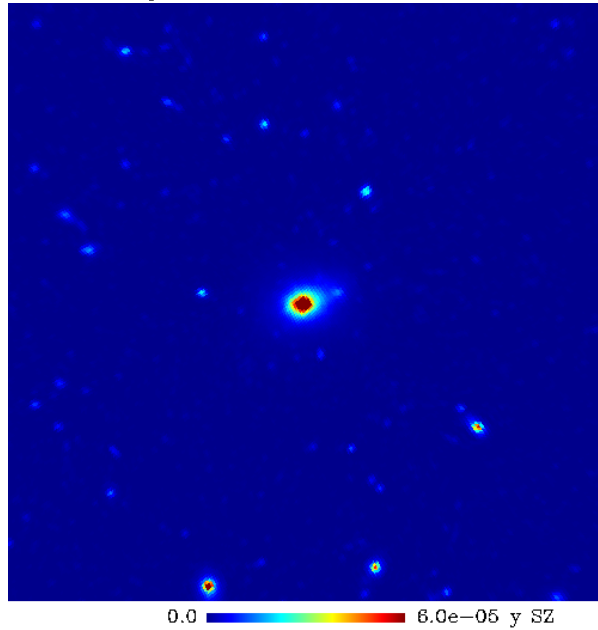


$$\mathbf{w}^t = \frac{\mathbf{a}^t \mathbf{C}^{-1}}{\mathbf{a}^t \mathbf{C}^{-1} \mathbf{a}}$$

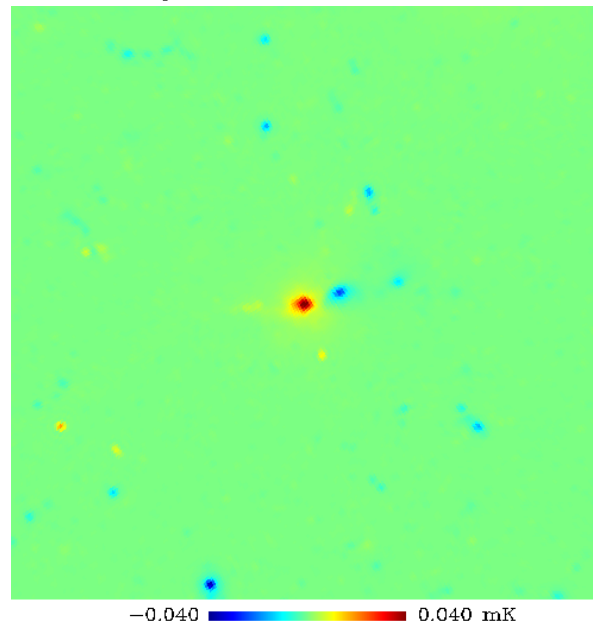
Bennett et al (2003), Tegmark et al (2003)
Eriksen et al (2004), Delabrouille et al (2009)

Standard ILC

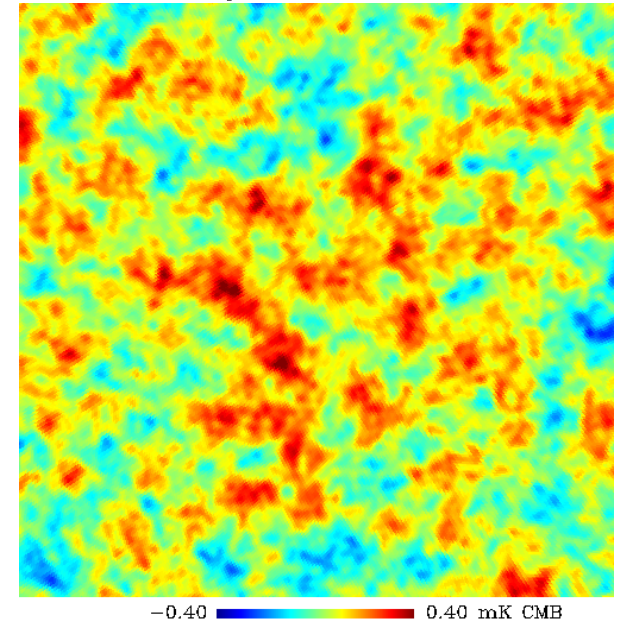
input thermal SZ



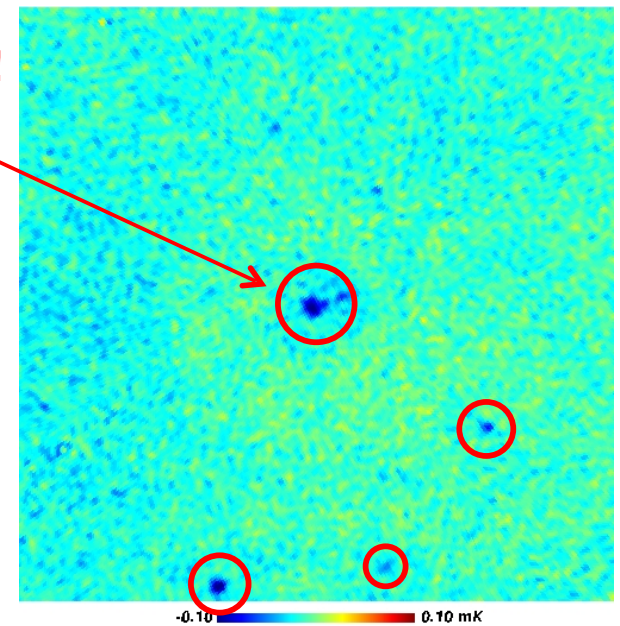
input kinetic SZ



input CMB



error: ILC - CMB



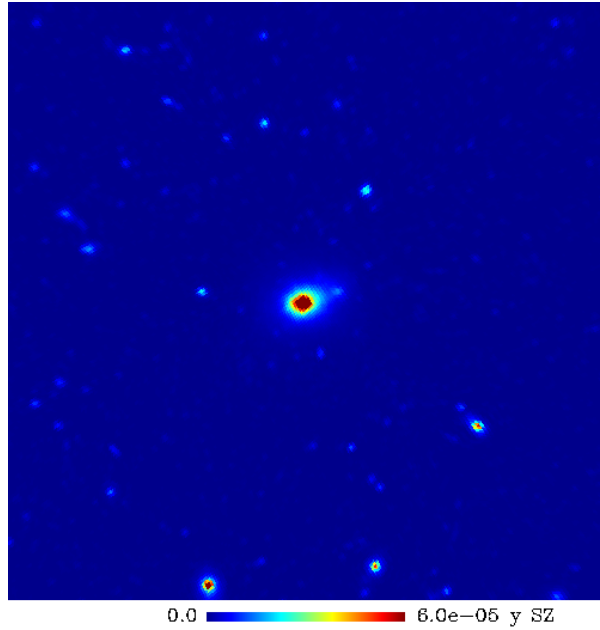
Thermal SZ / y-distortion residuals!
(clusters in the CMB)

$$\mathbf{w}^t = \frac{\mathbf{a}^t \mathbf{C}^{-1}}{\mathbf{a}^t \mathbf{C}^{-1} \mathbf{a}}$$

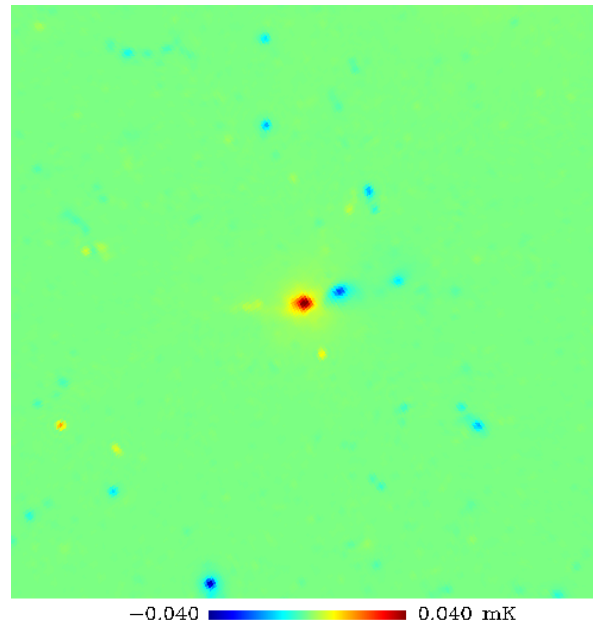
Bennett et al (2003), Tegmark et al (2003)
Eriksen et al (2004), Delabrouille et al (2009)

Constrained-ILC

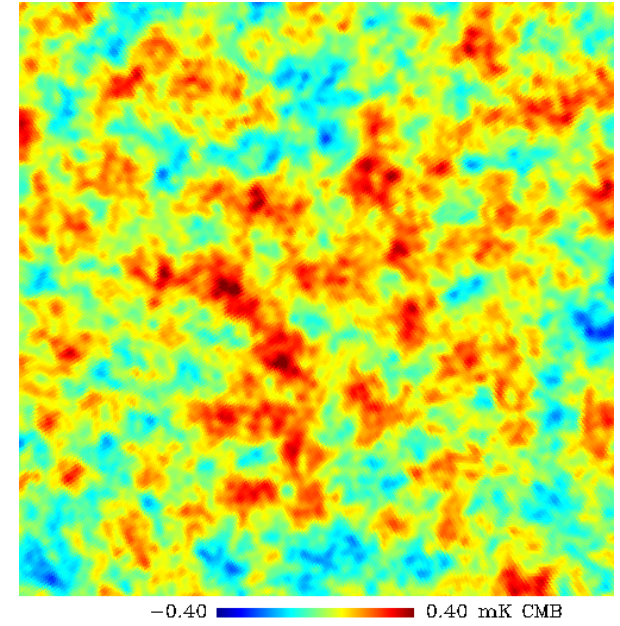
input thermal SZ



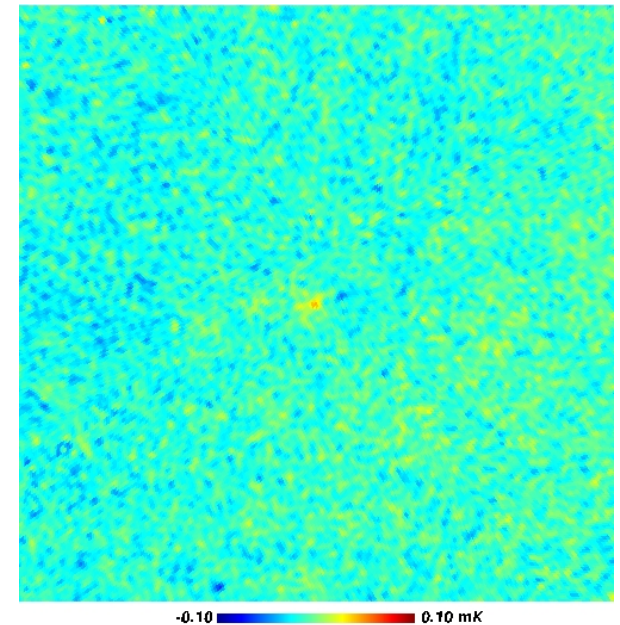
input kinetic SZ



input CMB



error: Constrained ILC - CMB



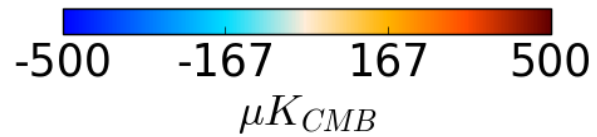
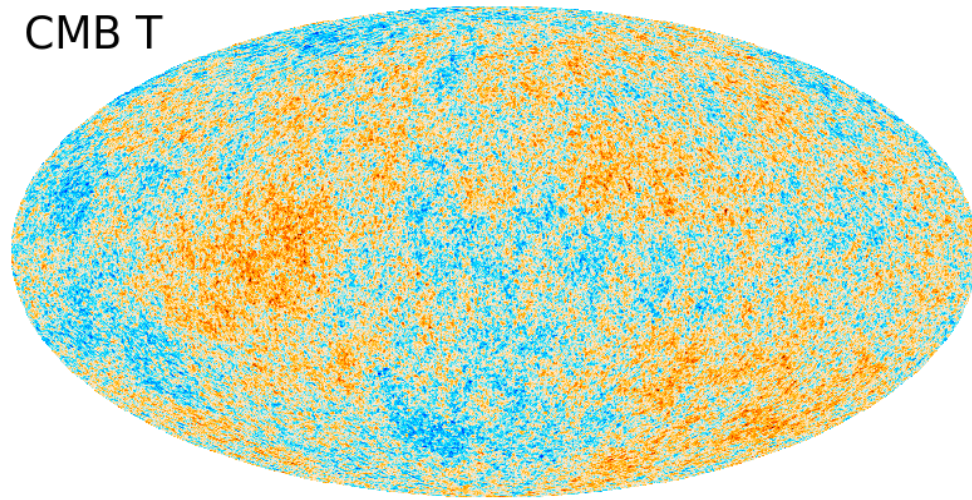
$$\mathbf{w}^t = \frac{(\mathbf{b}^t \mathbf{C}^{-1} \mathbf{b}) \mathbf{a}^t \mathbf{C}^{-1} - (\mathbf{a}^t \mathbf{C}^{-1} \mathbf{b}) \mathbf{b}^t \mathbf{C}^{-1}}{(\mathbf{a}^t \mathbf{C}^{-1} \mathbf{a})(\mathbf{b}^t \mathbf{C}^{-1} \mathbf{b}) - (\mathbf{a}^t \mathbf{C}^{-1} \mathbf{b})^2}$$

Remazeilles, Delabrouille, Cardoso, MNRAS (2011)

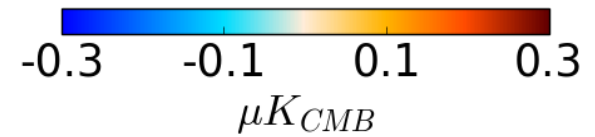
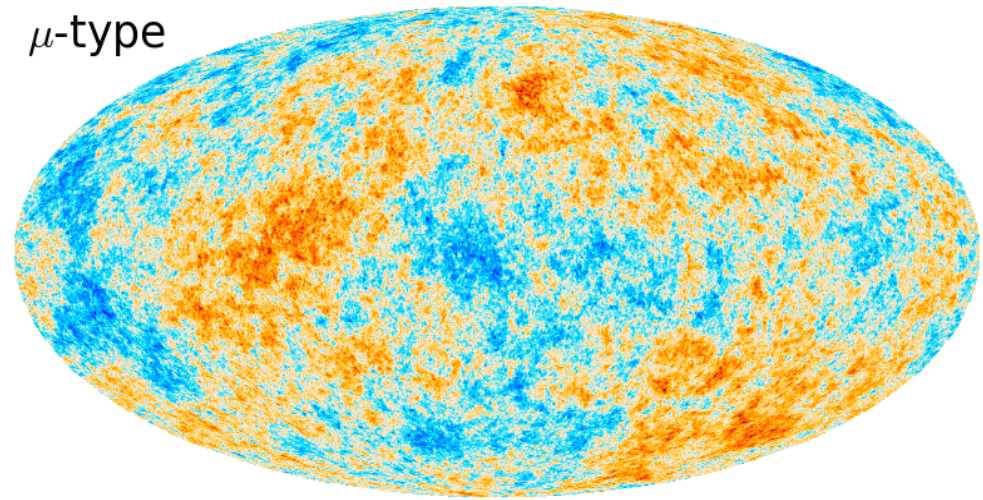
*Forecasts on μ -distortion anisotropies
from
sky simulations of future CMB satellites*

Simulation of correlated μ and T fields

CMB T



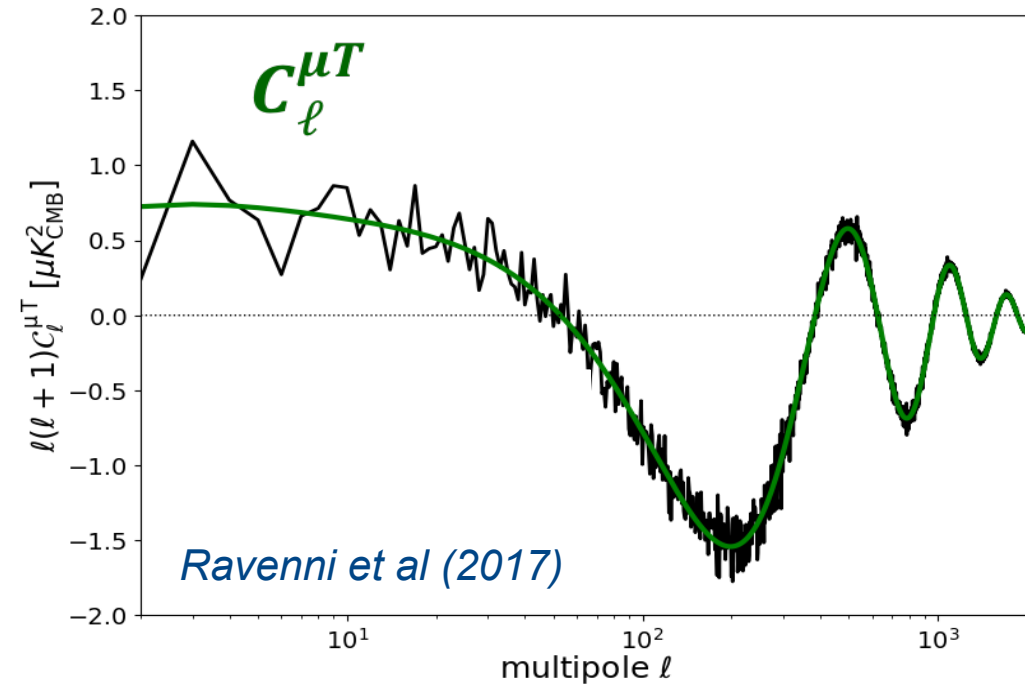
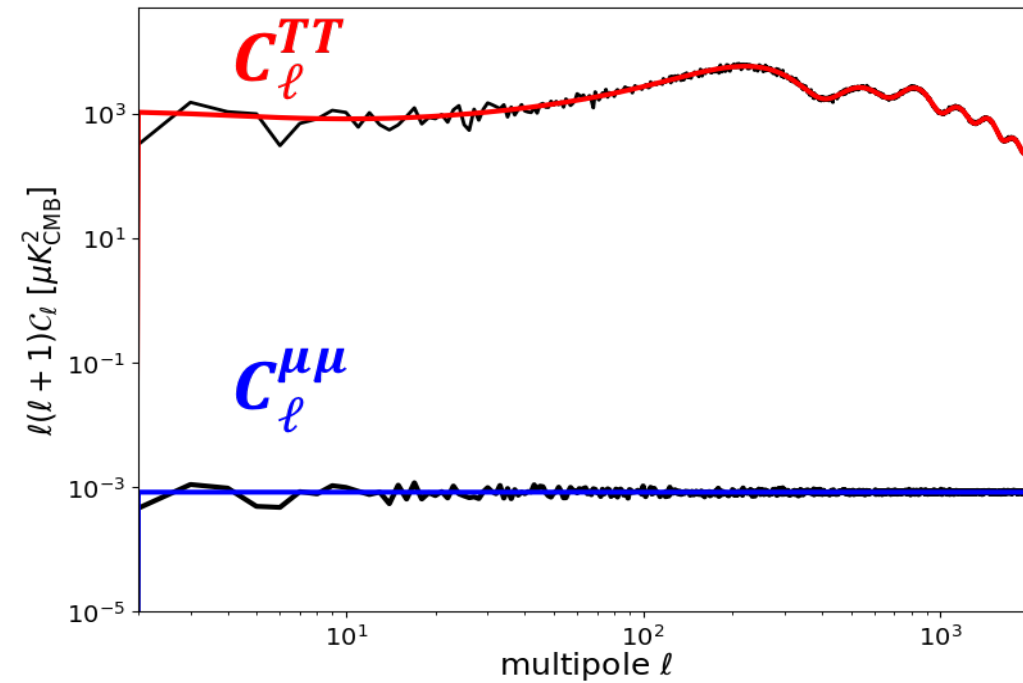
μ -type



$$\langle \mu \rangle = 2 \times 10^{-8}$$

$$f_{NL}(k \simeq 740 \text{ Mpc}^{-1}) = 4500$$

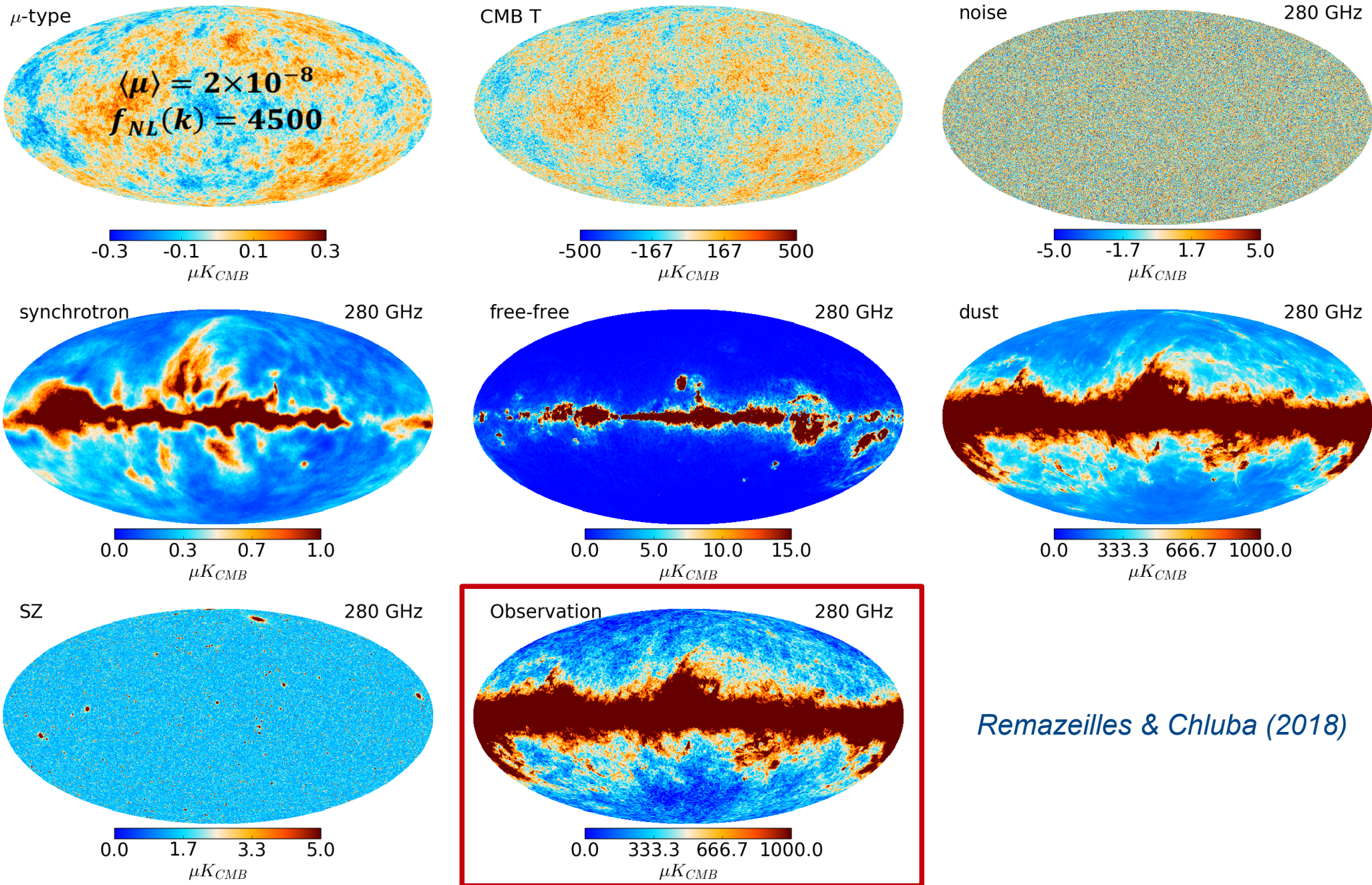
Simulation of correlated μ and T fields



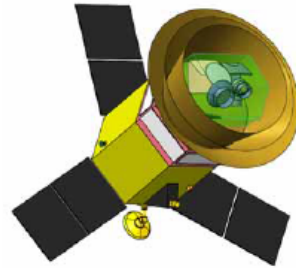
$$\langle \mu \rangle = 2 \times 10^{-8}$$

$$f_{NL}(k \simeq 740 \text{ Mpc}^{-1}) = 4500$$

Our full sky simulations



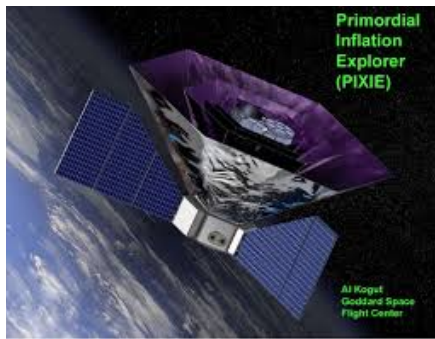
CMB satellite concepts



LiteBIRD (JAXA) – Phase A

Suzuki et al, 2018

40 – 402 GHz ; 2.5 μ K.arcmin



PIXIE (NASA)

Kogut et al., 2011

30 – 6000 GHz ;

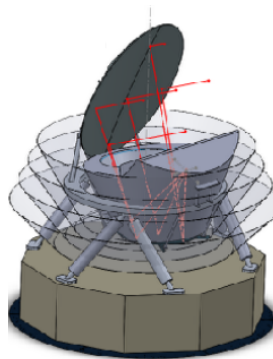
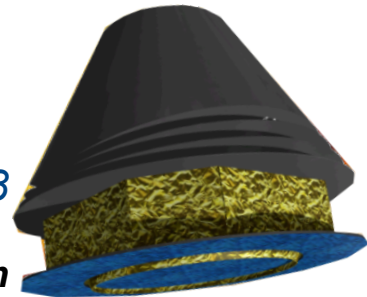
6.6 μ K.arcmin

($\Delta\nu=30$ GHz)

**CORE (ESA)
CMB-Bharat (ISRO)**

Delabrouille et al, 2018

60 – 600 GHz ; 1.7 μ K.arcmin



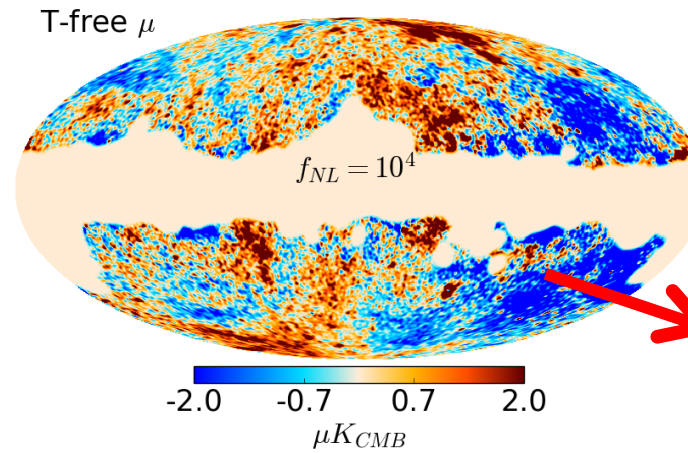
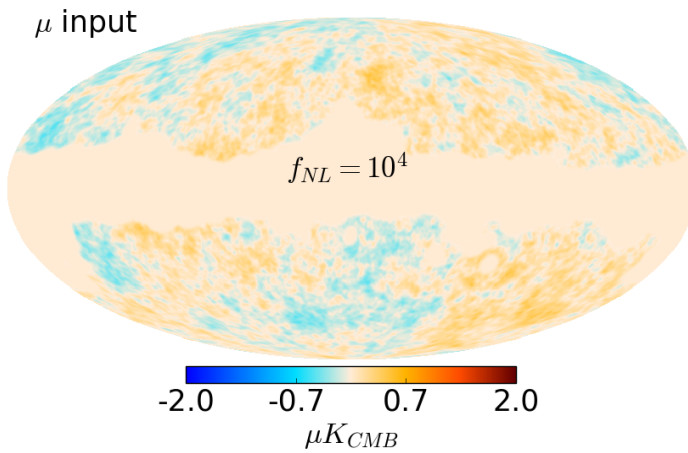
PICO (NASA)

Probe Mission Concept Study 2018

Shaul Hanany, priv. comm.

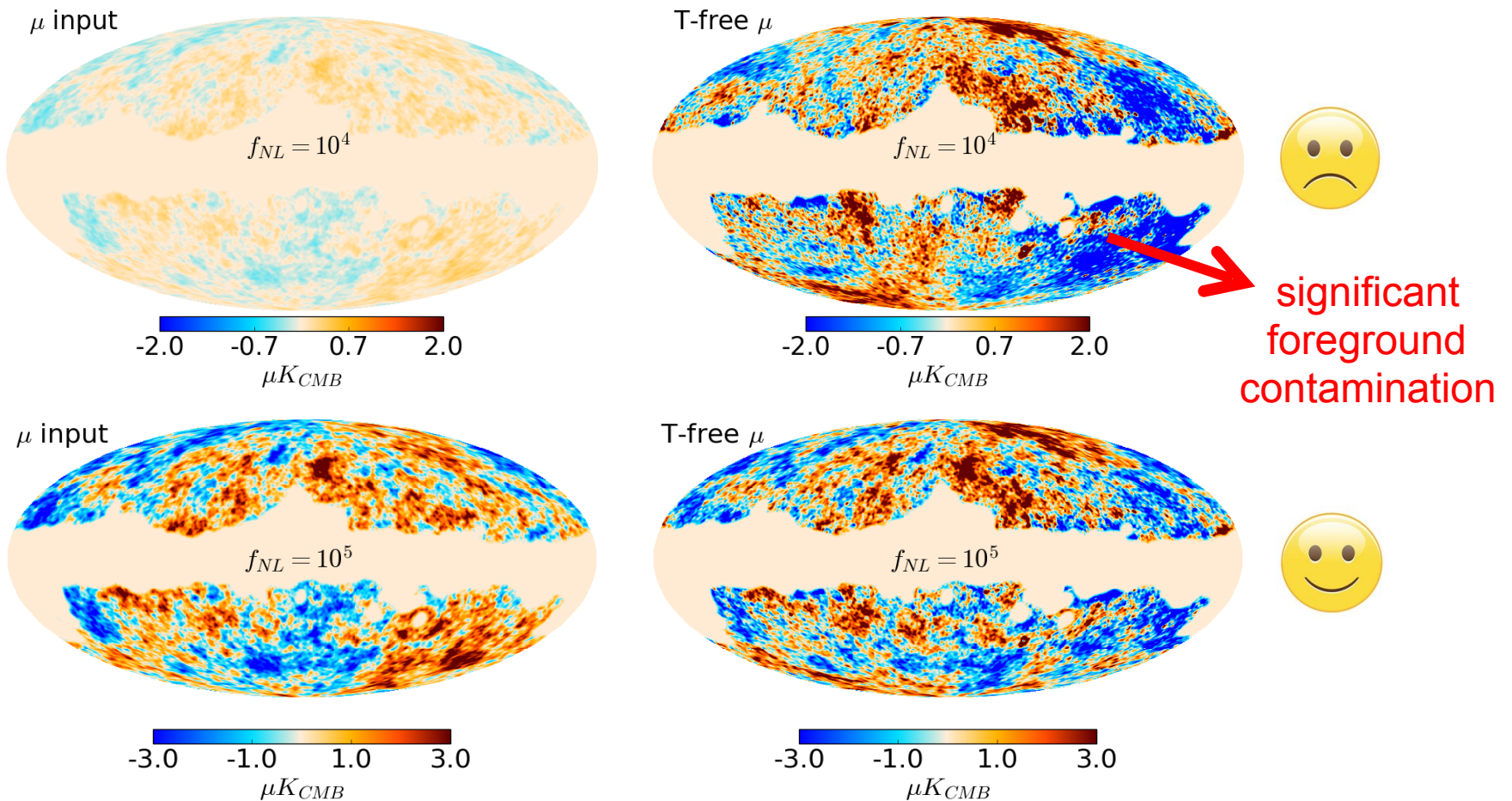
21 – 800 GHz ; 1.1 μ K.arcmin

Constrained-ILC μ -map reconstruction

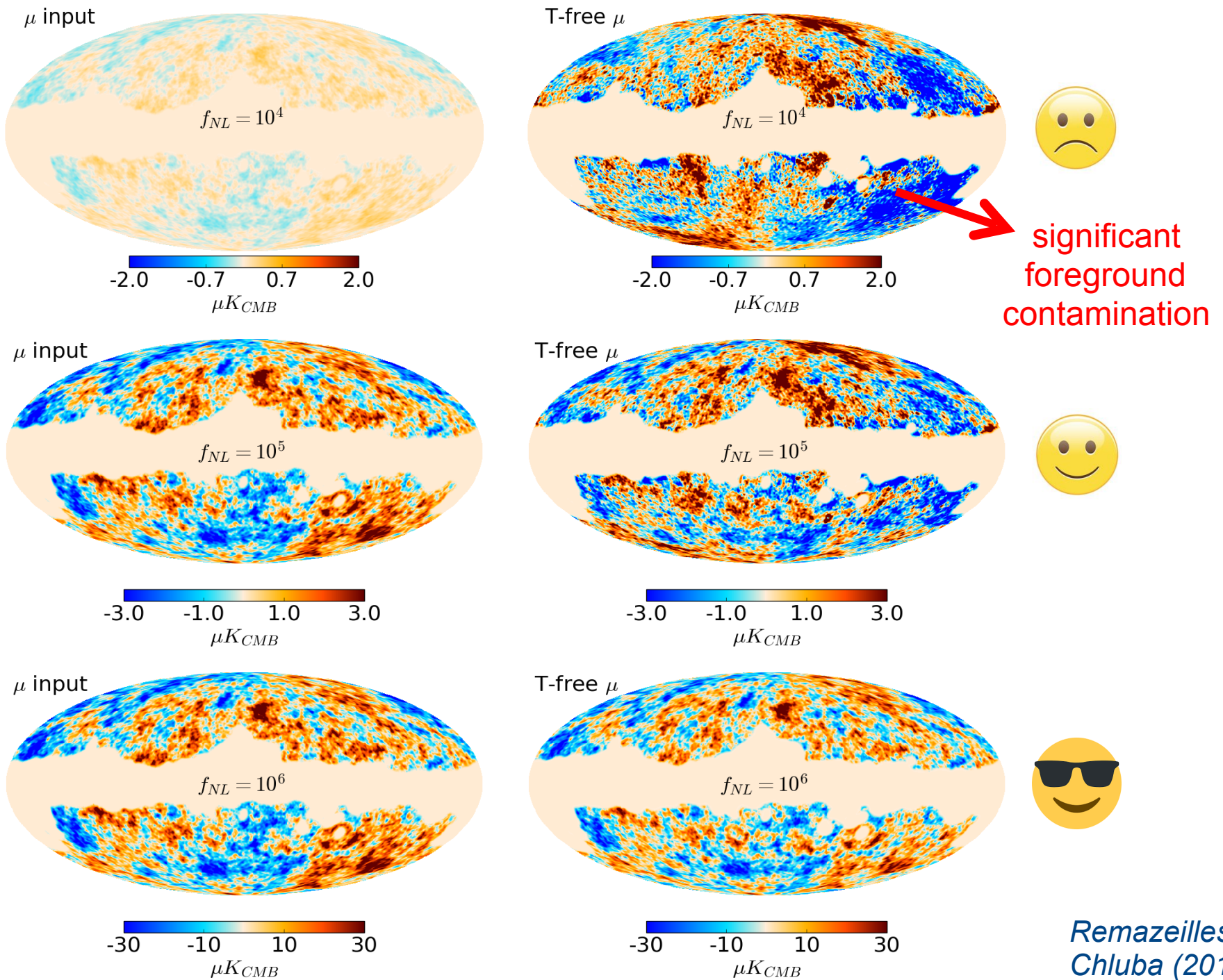


significant
foreground
contamination

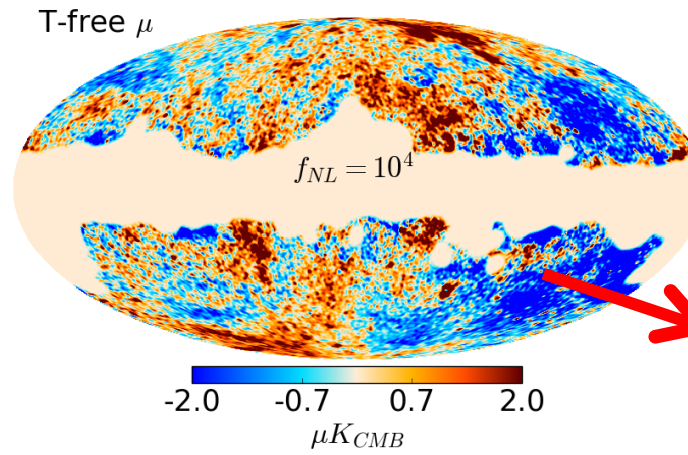
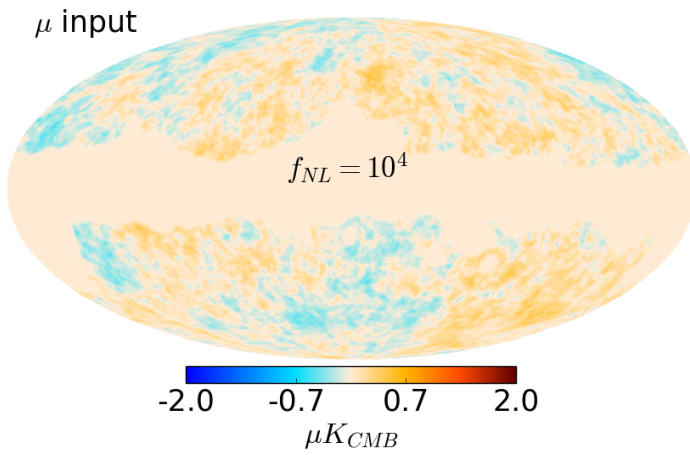
Constrained-ILC μ -map reconstruction



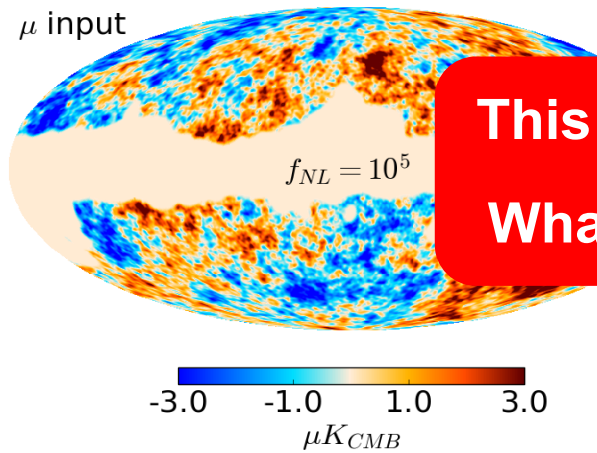
Constrained-ILC μ -map reconstruction



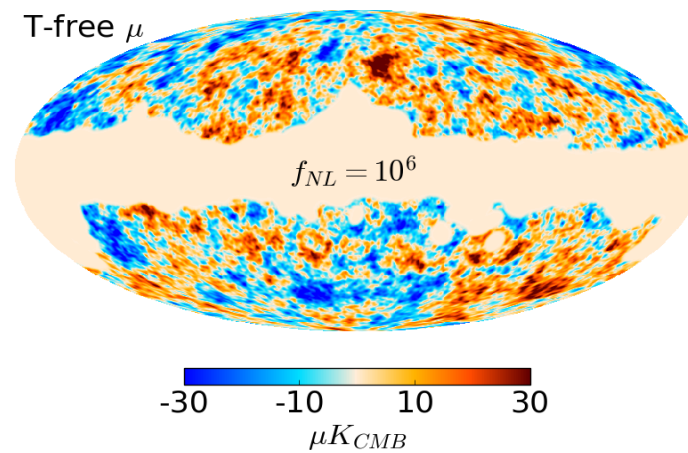
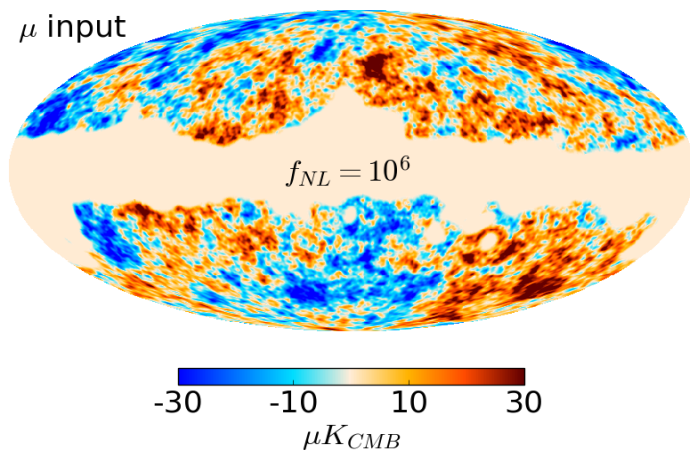
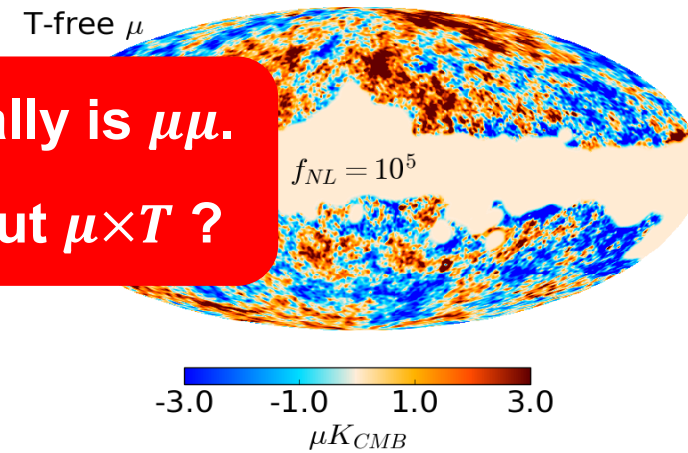
Constrained-ILC μ -map reconstruction



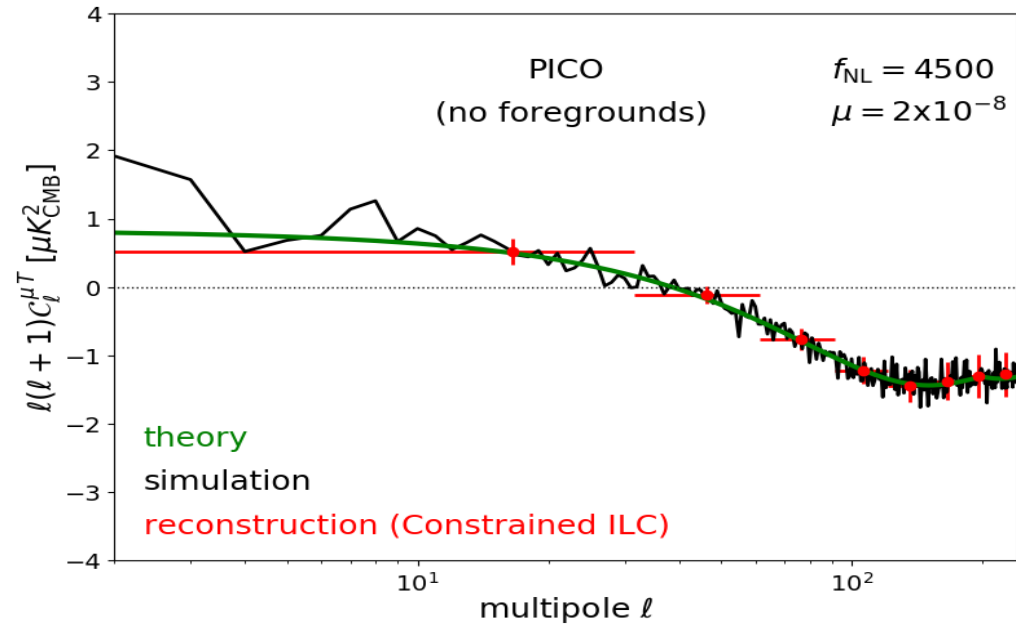
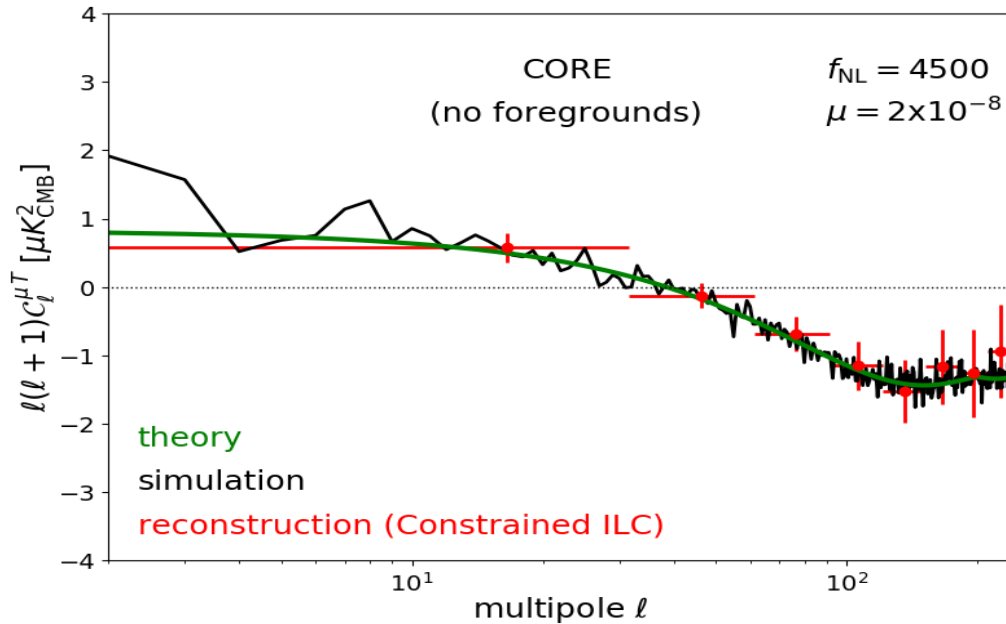
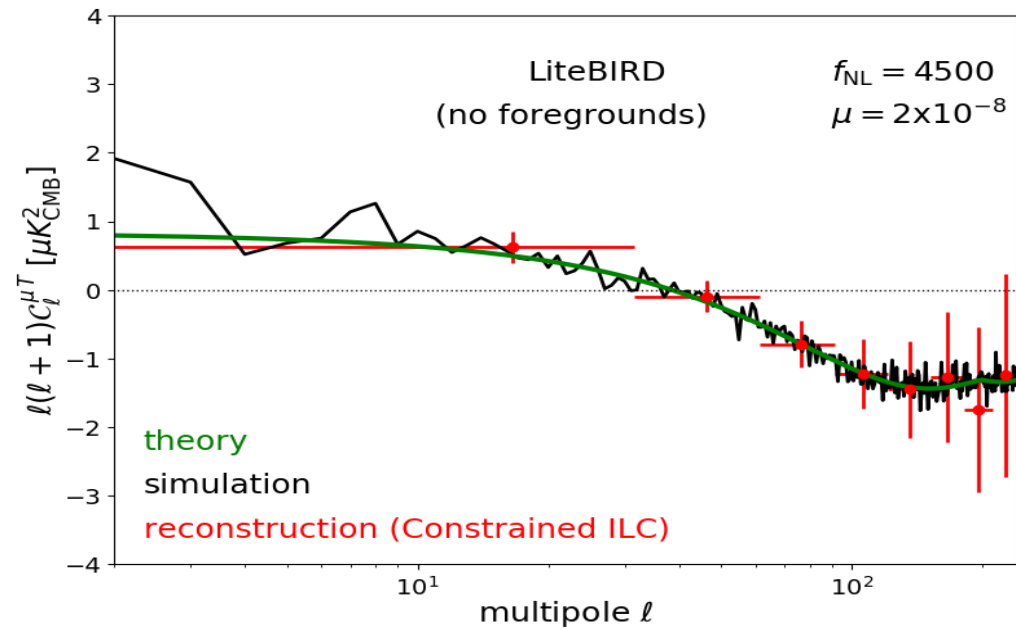
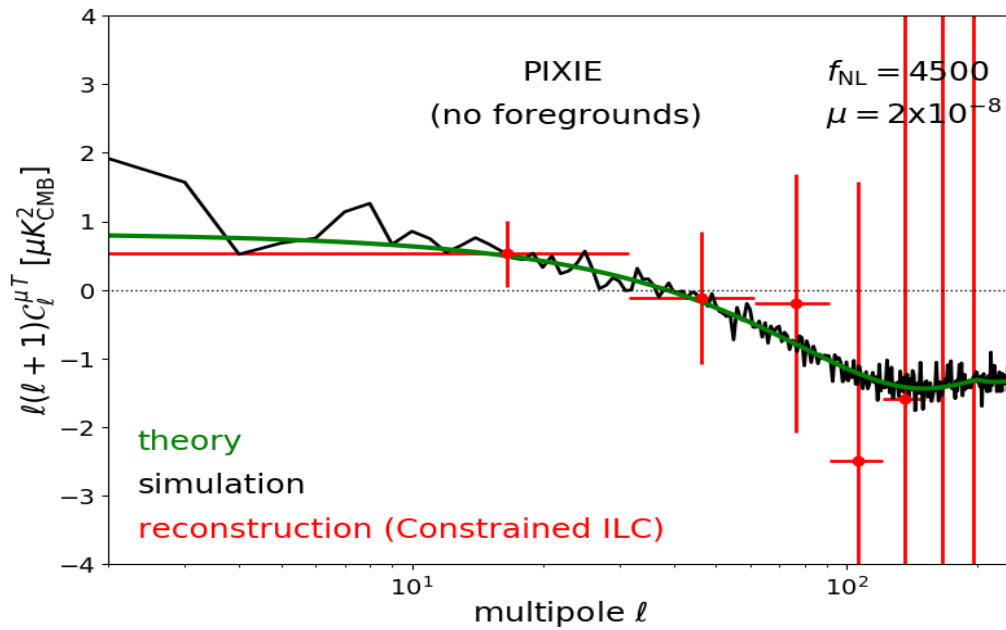
- significant foreground contamination



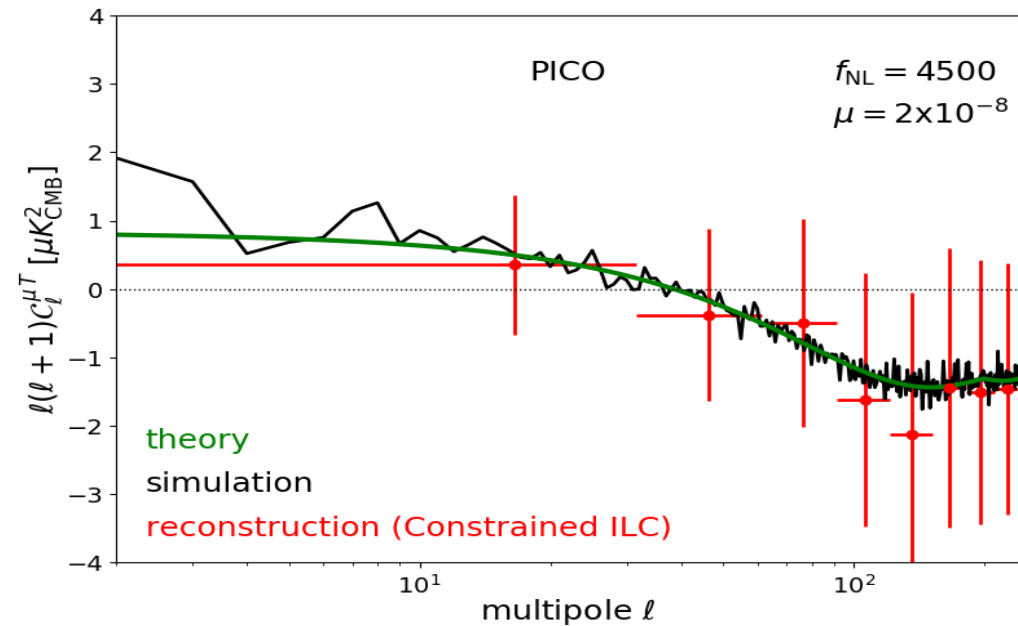
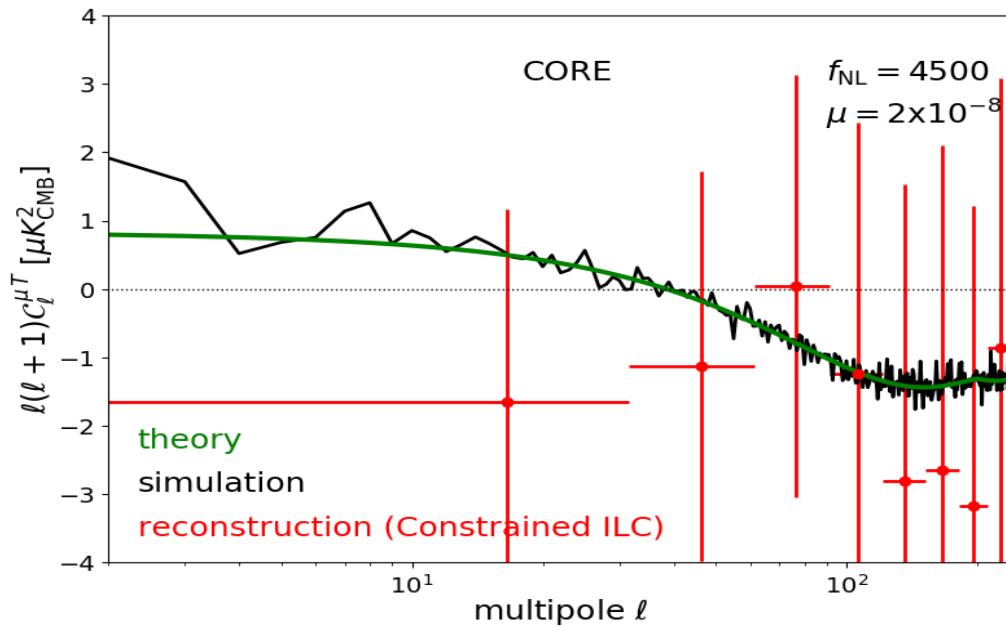
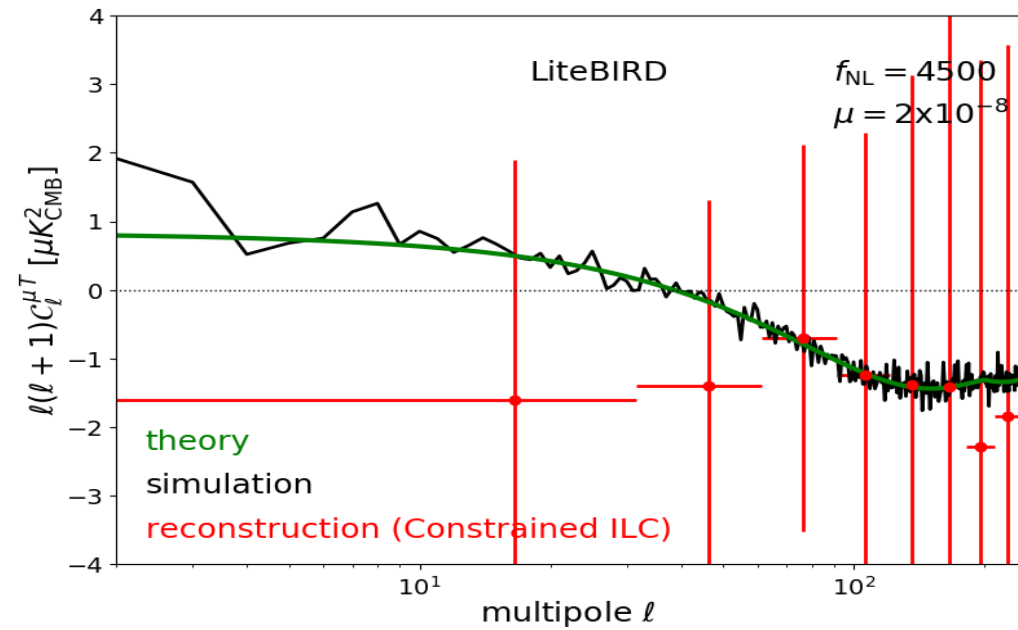
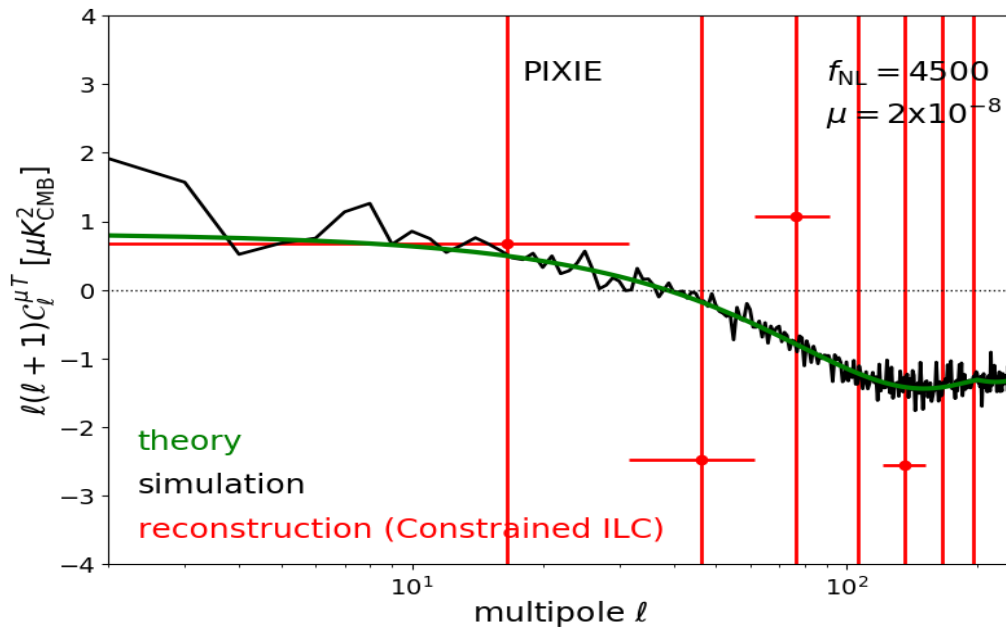
This actually is $\mu\mu$.
What about $\mu \times T$?



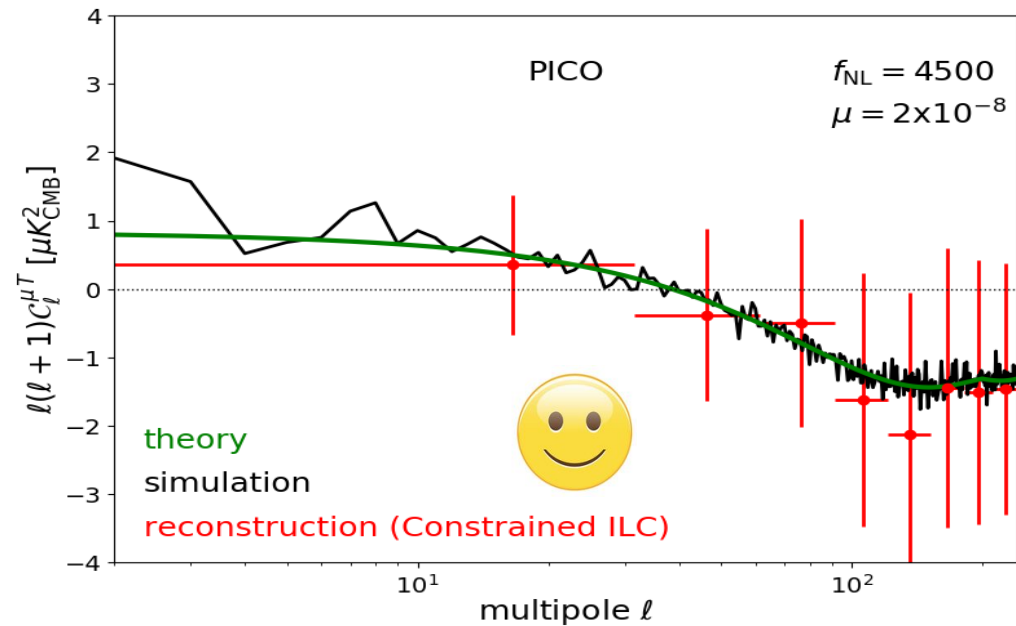
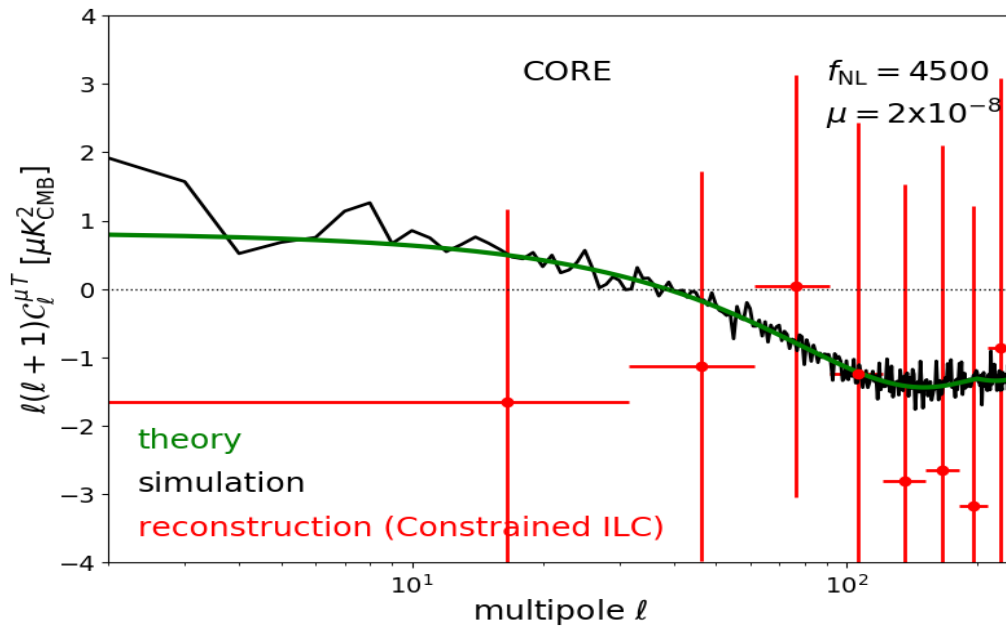
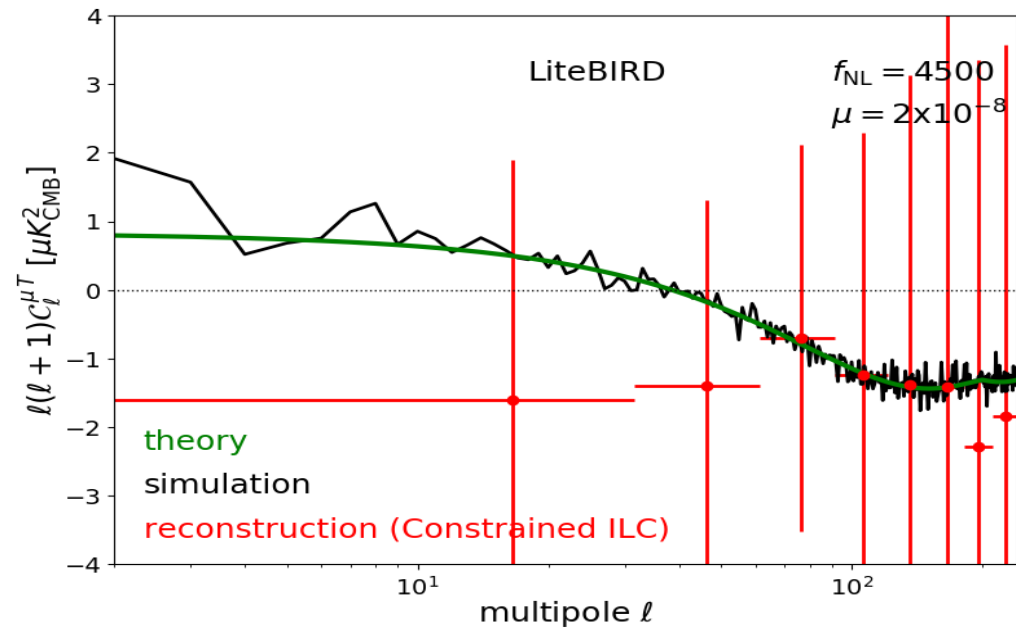
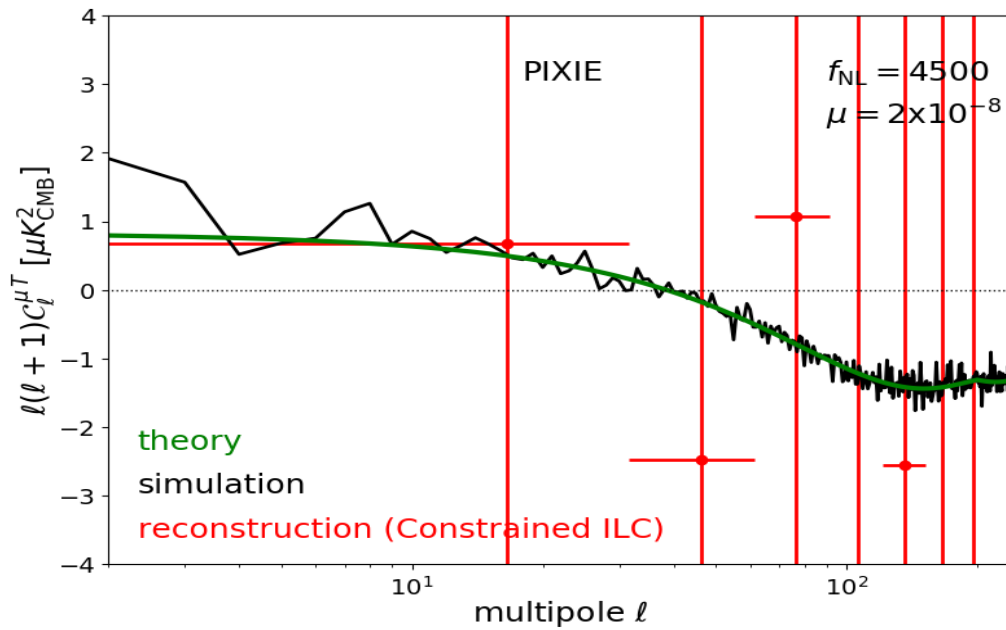
$C_\ell^{\mu T}$ reconstruction: $f_{NL} = 4500$ (w/o foregrounds)



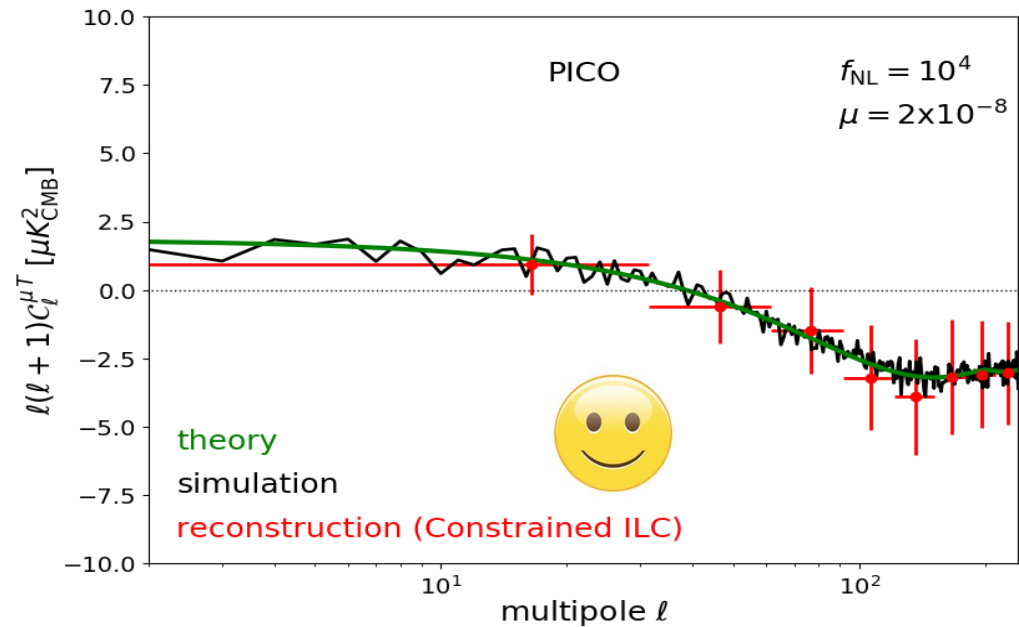
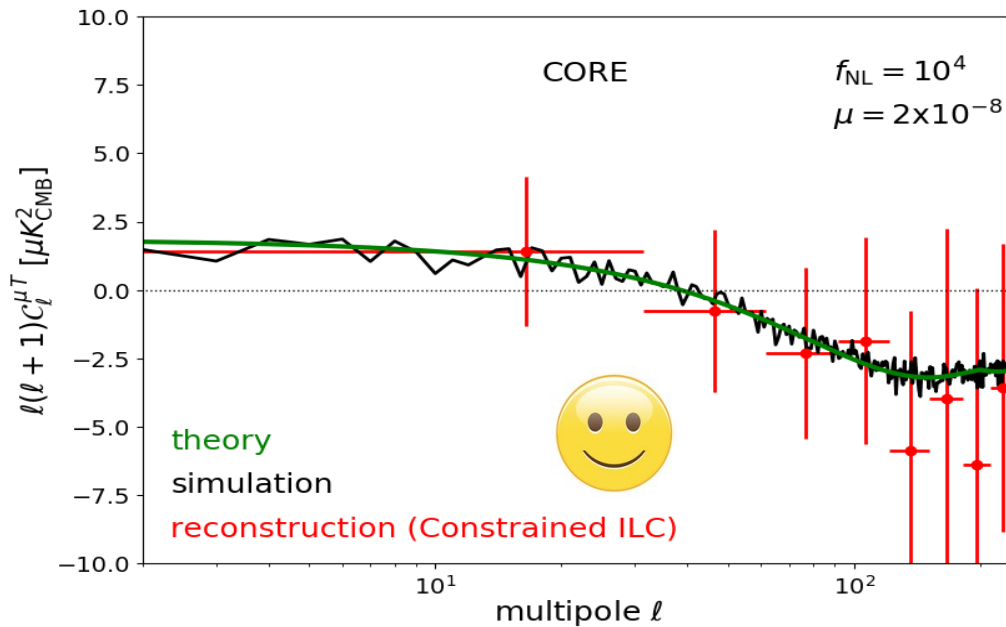
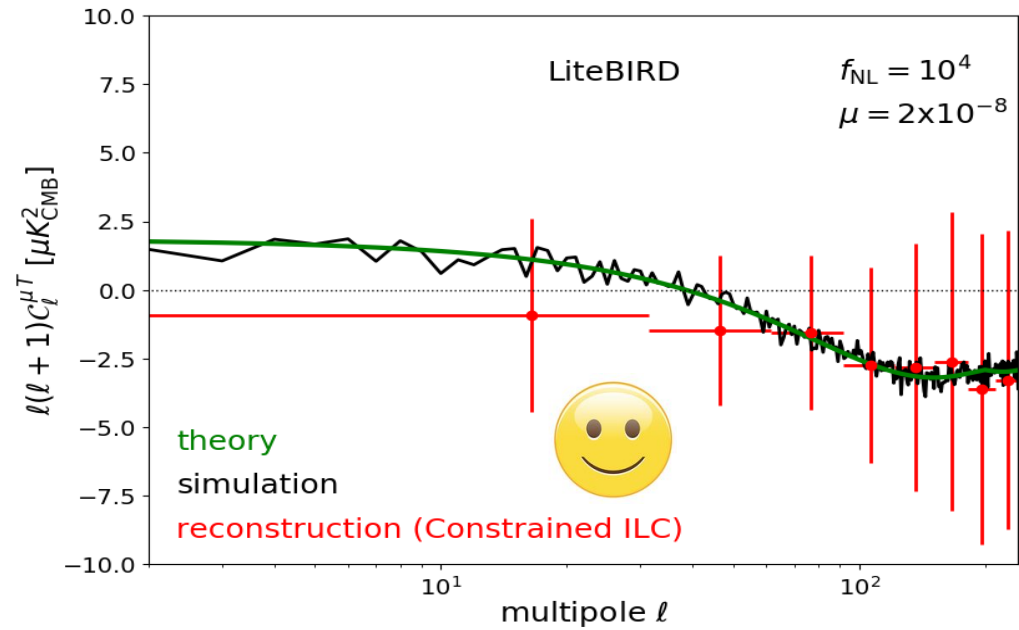
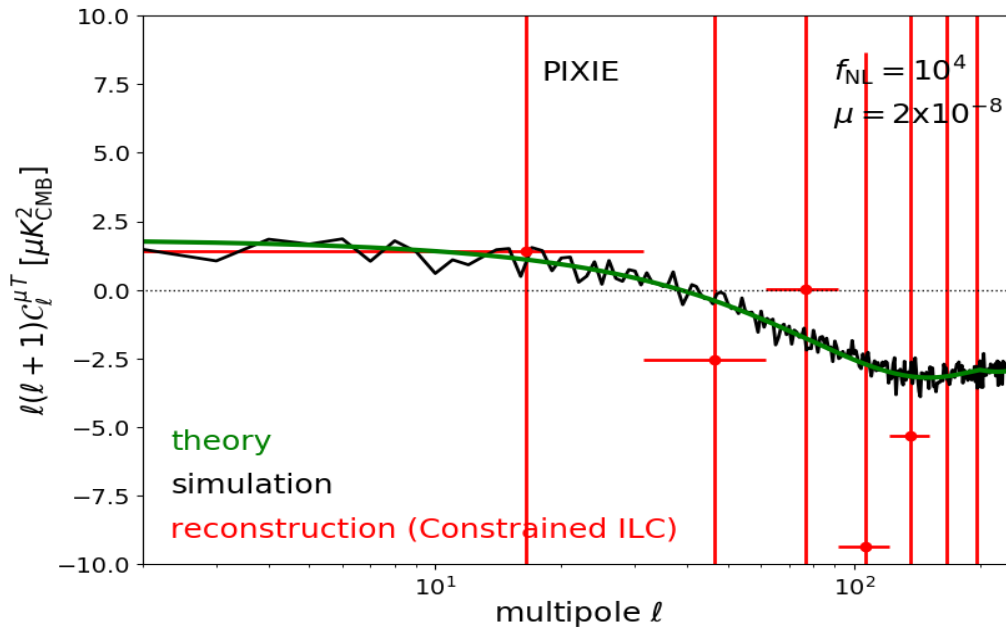
$C_\ell^{\mu T}$ reconstruction: $f_{NL} = 4500$ (with foregrounds)



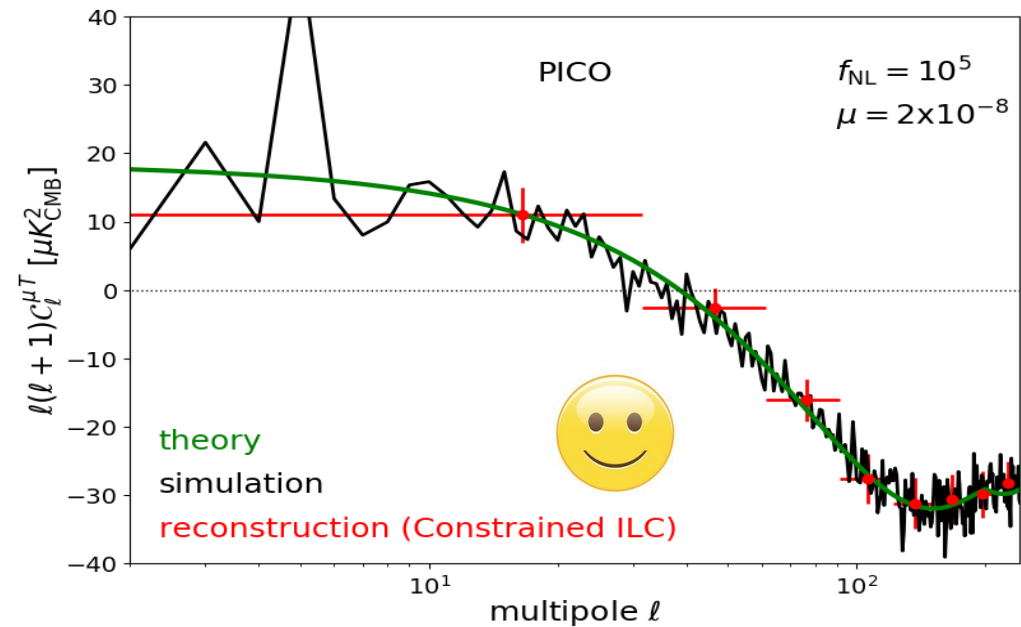
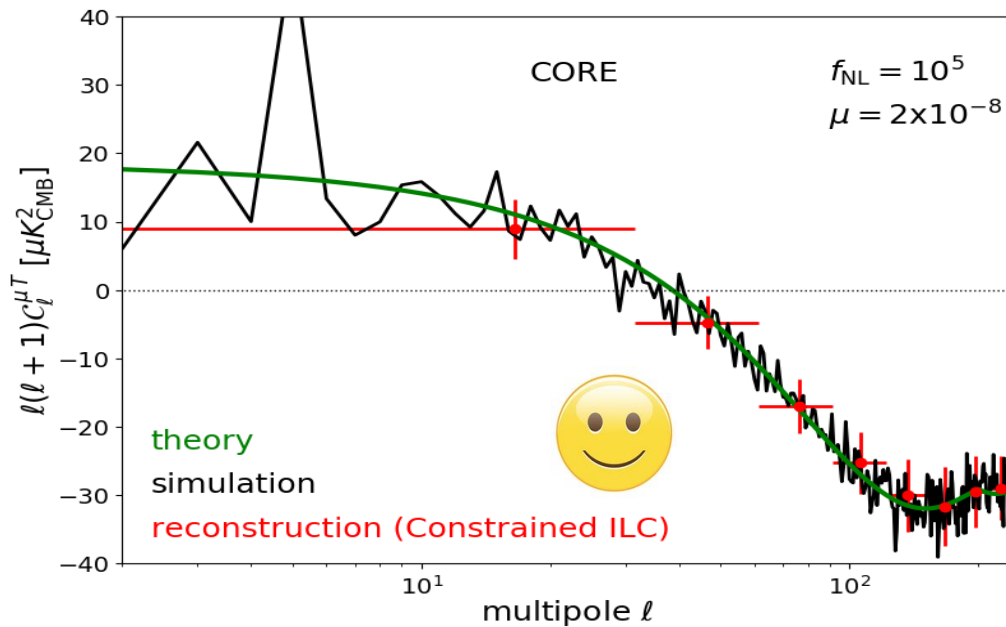
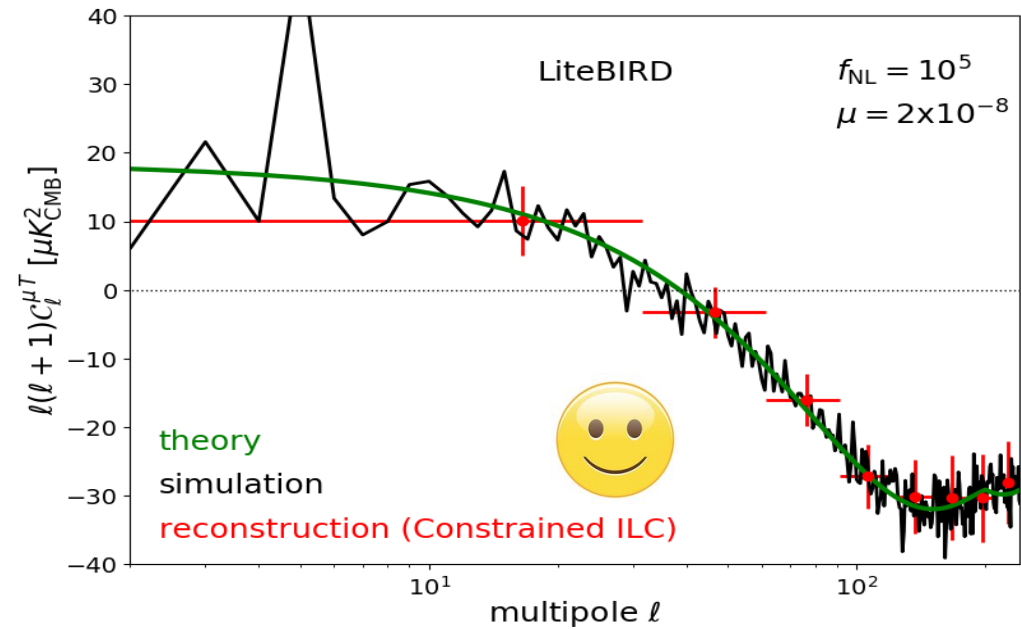
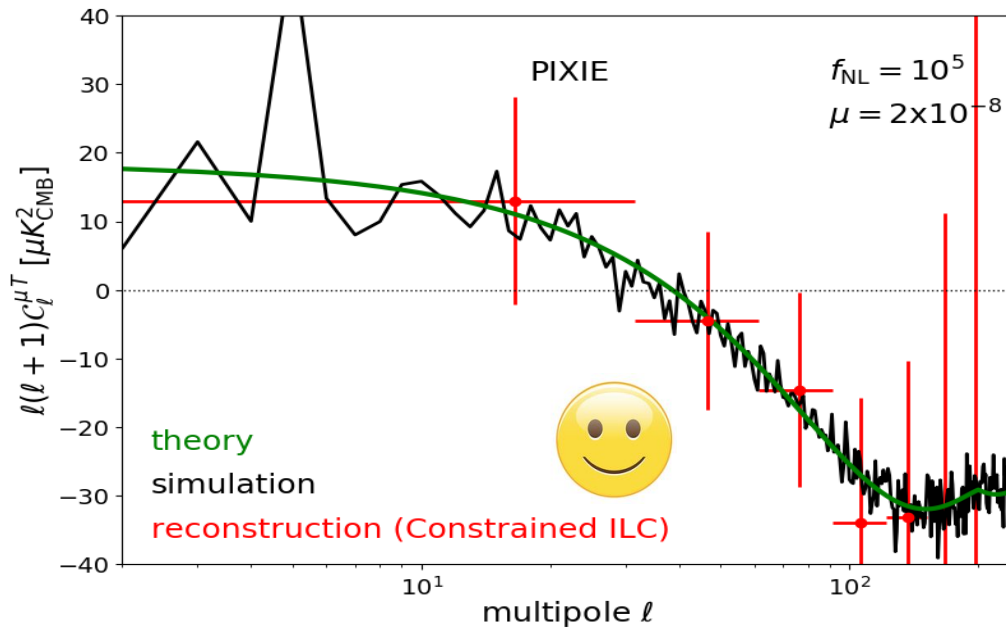
$C_\ell^{\mu T}$ reconstruction: $f_{NL} = 4500$ (with foregrounds)



$C_\ell^{\mu T}$ reconstruction: $f_{NL} = 10^4$ (with foregrounds)



$C_\ell^{\mu T}$ reconstruction: $f_{NL} = 10^5$ (with foregrounds)



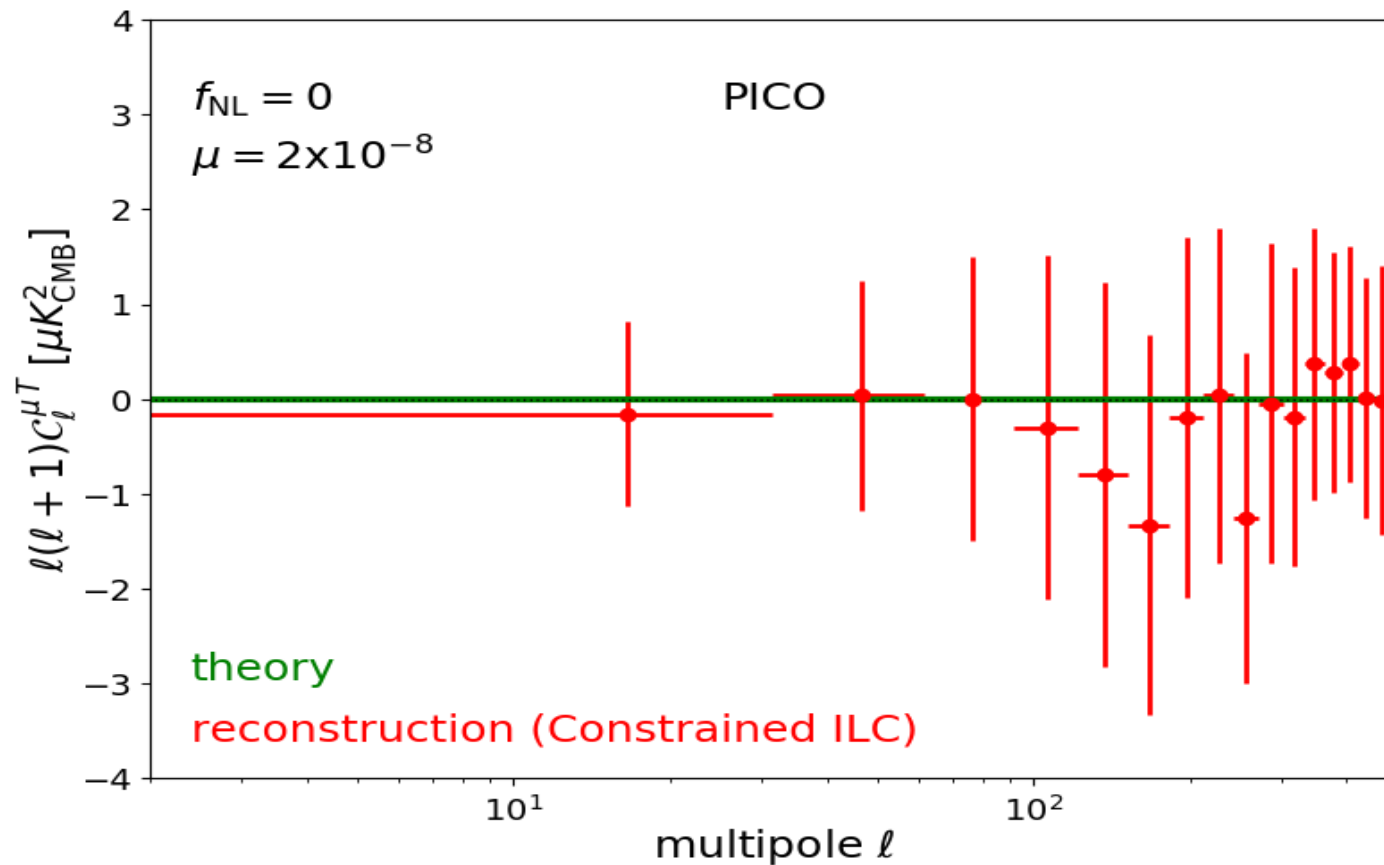
Forecasts on small-scale non-Gaussianity

Table 5. Detection forecasts on $f_{\text{NL}}(k \simeq 740 \text{ Mpc}^{-1})$ after component separation, based on multipoles $2 \leq \ell \leq 200$.

f_{NL} (fiducial)	10^5	10^4	4500	4500 w/o foregrounds
<i>PIXIE</i>	$(1.11 \pm 0.40) \times 10^5$ 2.5σ	$(2.17 \pm 3.90) \times 10^4$ –	$(1.5 \pm 3.9) \times 10^4$ –	4778 ± 3868 1.2σ
<i>LiteBIRD</i>	$(0.98 \pm 0.08) \times 10^5$ 12.5σ	$(0.91 \pm 0.68) \times 10^4$ 1.5σ	4272 ± 6788 –	4753 ± 930 4.8σ
<i>CORE</i>	$(0.97 \pm 0.08) \times 10^5$ 12.5σ	$(1.35 \pm 0.74) \times 10^4$ 1.4σ	5692 ± 6397 –	4336 ± 653 6.9σ
<i>PICO</i>	$(0.99 \pm 0.06) \times 10^5$ 17.8σ	$(1.07 \pm 0.30) \times 10^4$ 3.3σ	5094 ± 2929 1.5σ	4480 ± 371 12.1σ

PICO is in the best position to detect the μ -T correlation signal at $f_{\text{NL}}(k \simeq 740 \text{ Mpc}^{-1}) \lesssim 4500$ in the presence of foregrounds

Null-test: μ -T signal reconstruction for $f_{NL} = 0$



In the absence of μ -distortion anisotropies, the reconstruction by Constrained ILC is consistent with $f_{NL} = 0$

Minimum detection limit

Table 6. Detection limits for *PICO* on $f_{\text{NL}}(k \simeq 740 \text{ Mpc}^{-1})$ after component separation, based on the multipole range $2 \leq \ell \leq 500$ using the model of [Ravenni et al. \(2017\)](#) to describe the $\mu - T$ cross-correlation. Foregrounds are included in all cases and the fiducial f_{NL} parameter was varied.

f_{NL} (fiducial)	−4500	0	4500
<i>PICO</i>	-2996 ± 2112	1325 ± 2114	5698 ± 2121
	2σ	–	2σ

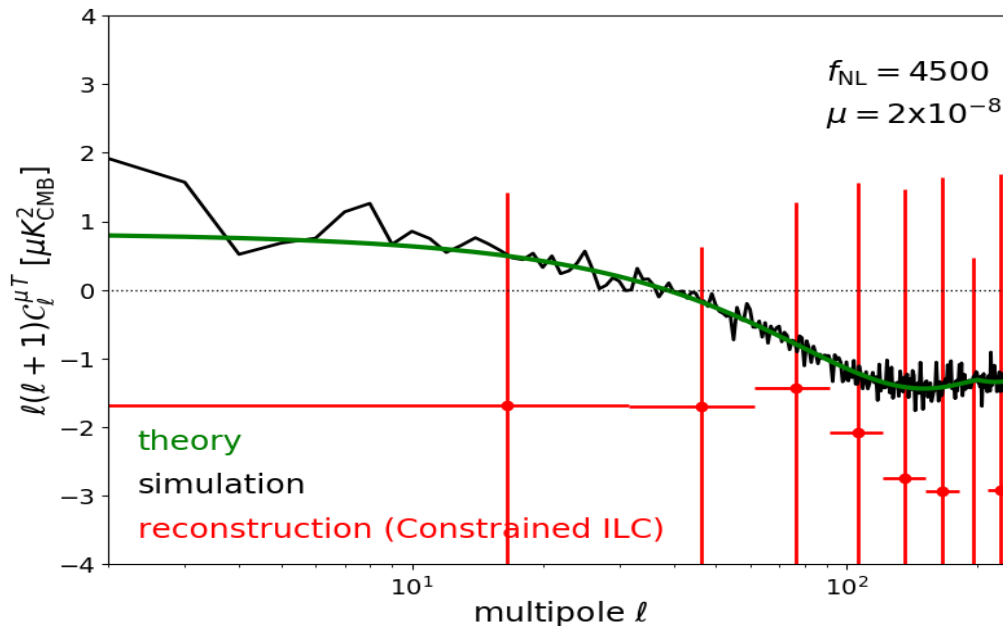
Minimum detection limit by PICO in the presence of foregrounds:

$$|f_{\text{NL}}(k \simeq 740 \text{ Mpc}^{-1})| \gtrsim 2100$$

More detectors or more frequencies?

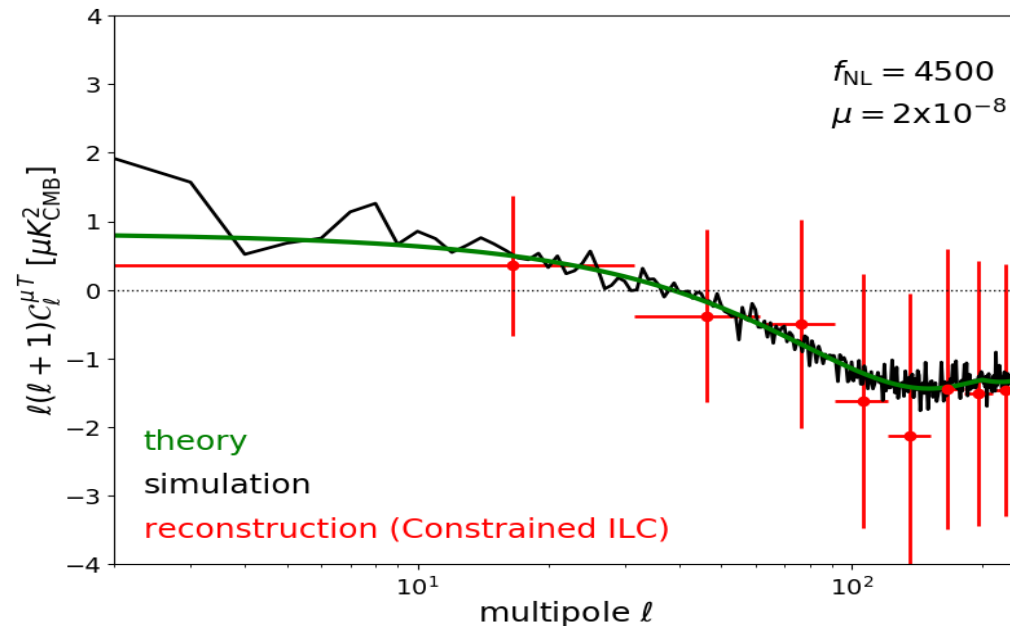
“Enhanced LiteBIRD” with $100 \times$ more detectors

0.2 $\mu\text{K. arcmin}$, 40 - 400 GHz



PICO

0.8 $\mu\text{K. arcmin}$, 20 - 800 GHz



Extended frequency coverage at frequencies $\nu \lesssim 40$ GHz and $\nu \gtrsim 400$ GHz provides more leverage than increased channel sensitivity

What part of the frequency range matters?

- Discarding PICO frequencies above $\nu > 400$ GHz degrades component separation results by $\simeq 7\%$
- Discarding PICO frequencies below $\nu < 40$ GHz degrades component separation results by $\simeq 30\%$

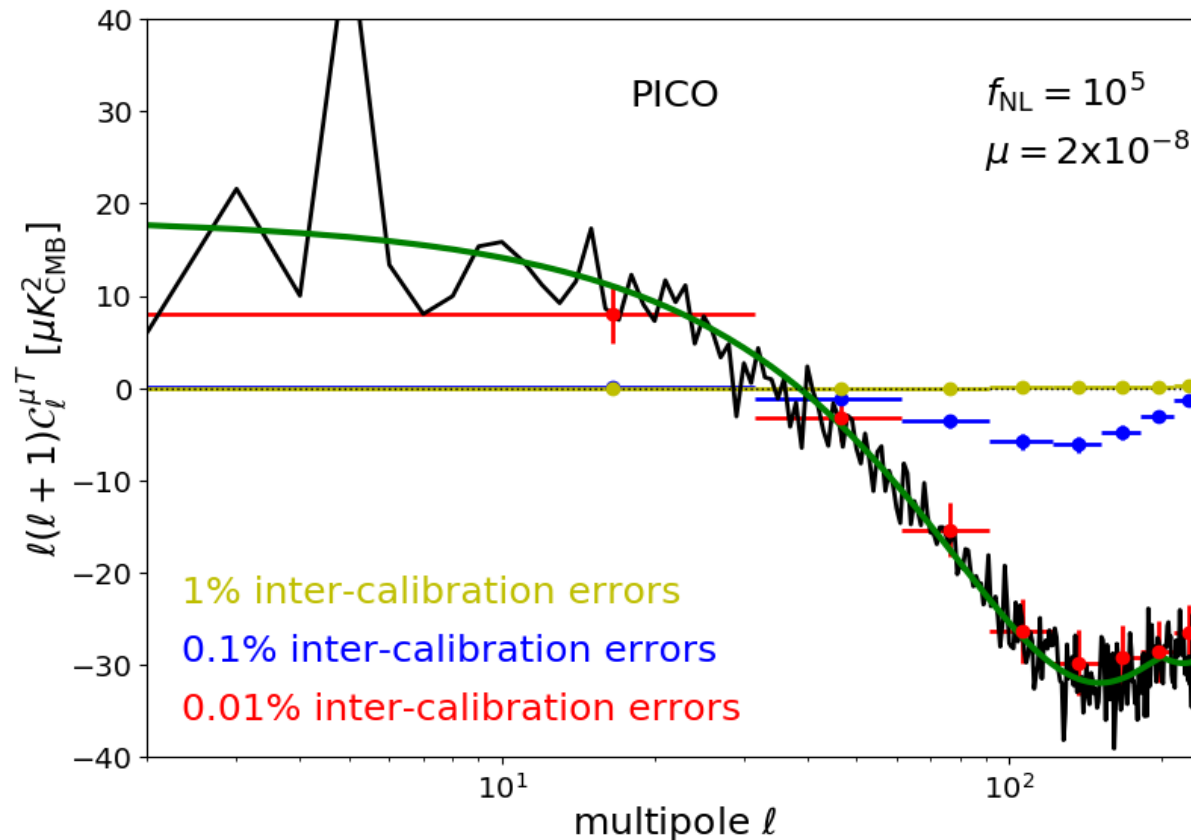
Low-frequencies $\nu < 40$ GHz have more constraining power on μ -distortion anisotropies than high-frequencies $\nu > 400$ GHz

Remazeilles & Chluba (2018)

→ consistent with the conclusions of *Abitbol et al (2017)* for monopole distortions

Inter-calibration errors kill CMB temperature (so may bias μ - T measurements)

Calibration errors can screw up an ILC in the high signal-to-noise regimes,
through partial cancellation of the variance of the CMB temperature anisotropies
Dick, Remazeilles, Delabrouille, MNRAS (2010)



The allowed inter-channel calibration uncertainty for PICO is 0.01 %
(The promise of CORE was to achieve such calibration accuracy)

Conclusions

Remazeilles & Chluba (2018)

- We have computed the first forecasts on the detection of the μ - T correlation signal and $f_{NL}(k \simeq 740 \text{ Mpc}^{-1})$ in the presence of foregrounds with future CMB satellites
- We have proposed a tricky component separation approach, the “Constrained-ILC”, to null out residual CMB TT correlations in the μ - T correlation signal
- Among CMB satellite concepts, PICO is in the best position to detect anisotropic μ -type distortions in the presence of foregrounds, with $f_{NL}(k \simeq 740 \text{ Mpc}^{-1}) \lesssim 2100$
- Optimization: more detectors or more frequencies?

Extended frequency coverage at frequencies $\nu \lesssim 40 \text{ GHz}$ and $\nu \gtrsim 400 \text{ GHz}$ provides more leverage than increased channel sensitivity

Low-frequencies $\nu \lesssim 40 \text{ GHz}$ have more constraining power on μ -distortions than high-frequencies $\nu \gtrsim 400 \text{ GHz}$

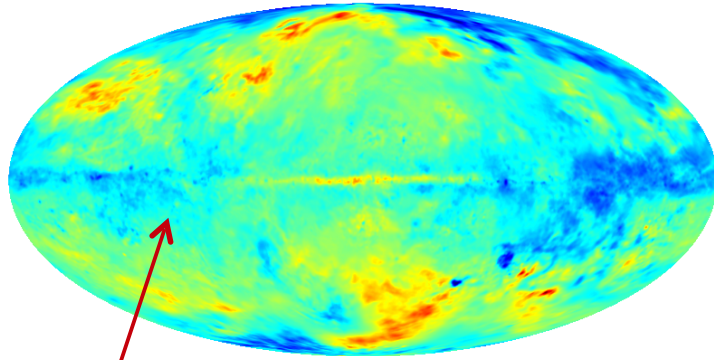
- Spectrometers like PIXIE or PRISTINE still needed for μ -distortion anisotropies to break the $f_{NL}\langle\mu\rangle$ degeneracy

Thank you for your attention!

Backup slides

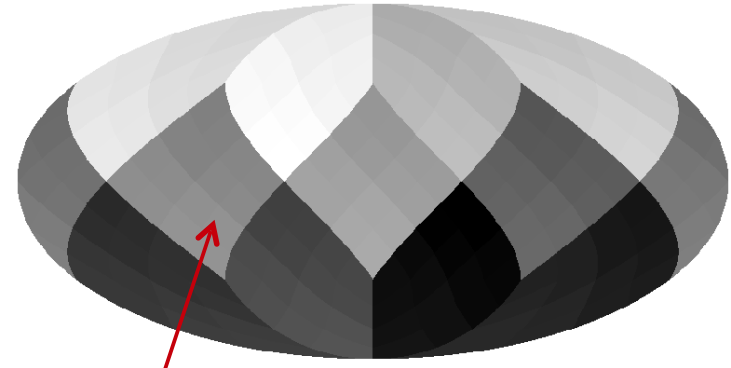
Averaging effects

dust spectral indices in the sky



one value β_{dust} per line-of-sight

mapping / pixelization



many β_{dust} values per pixel

→ effective SED: $\sum_i v^{\beta_i} \neq v^{\beta}$

- Because of averaging different line-of-sight SEDs within a pixel/beam, the actual SED of foregrounds *in the maps* differs from the physical SED *in the sky*
– Chluba et al 2017
- Spurious SED curvatures created by pixel averaging effects, if ignored in the parametric fit, may bias primordial B-modes at the level of $\Delta r \sim 10^{-3}$
– Remazeilles et al 2017, for the CORE collaboration
- The Constrained-ILC method is blind (no parametrization / assumption on foregrounds), therefore fairly insensitive to averaging effects
– Remazeilles & Chluba 2018

Importance of spatial resolution

Despite a very broad frequency coverage, PIXIE constraints on anisotropic μ -distortions are of poorer quality than those from LiteBIRD, CORE, PICO

→ *because of lower sensitivity and lower spatial resolution*

- We find that increasing PIXIE resolution from 96' to 40', while keeping its baseline sensitivity, would improve $\sigma(f_{NL})$ by 50%

→ *high-resolution channels enable using more correlated information to improve foreground cleaning*

- Suppose the foreground complexity can be captured by 10 degrees of freedom, then 15-20 frequency bands are enough to subtract foregrounds

→ *In this case, the most sensitive experiments will make a difference in the ILC trade-off of minimizing the balance between foreground and noise contaminations*

- However, if foreground complexity relies on more than 20 degrees of freedom, then the broad frequency range of PIXIE will make a difference with respect to imagers