

1. For *each* of the following expressions, say whether it is (or evaluates to) a ket, a bra, a complex number, or an operator:

(i)  $\langle\psi|\hat{L}_z$ , (ii)  $|1\rangle\langle 1|$ , (iii)  $|f\rangle \otimes |h\rangle$ , (iv)  $\langle\uparrow|\leftarrow\rangle$ .

2. Given that

$$|\leftarrow\rangle = \frac{|\uparrow\rangle + |\downarrow\rangle}{\sqrt{2}},$$

where  $\{|\uparrow\rangle, |\downarrow\rangle\}$  are the eigenstates of  $\hat{S}_z$ , evaluate the  $2 \times 2$  matrix representing the operator  $|\leftarrow\rangle\langle\leftarrow|$  in the  $S_z$  basis. What kind of operator is this?

3. (i) The following equations are written rather ‘loosely’: re-write in a form that is fully consistent:

$$\begin{aligned}\hat{M} &= \hat{S}_z^2 + 2\hbar^2 \\ \hat{J} &= \hat{L} + \hat{S}\end{aligned}$$

- (ii) In the equation

$$\hat{L}_z|m = 0\rangle = 0,$$

what does the symbol ‘0’ on the right-hand side stand for?

4. **Part of 2010 Exam Q2:** A certain ion in a crystal has a Hamiltonian  $\hat{H} = A\hat{L}_x^2$ , where  $\hat{L}_x$  is the  $x$ -component of orbital angular momentum and  $A$  is a positive constant. The ion has angular momentum quantum number  $l = 1$ . In the  $L_z$  representation, this gives

$$\hat{H} \xrightarrow{L_z} \frac{A\hbar^2}{2} \begin{pmatrix} 1 & 0 & 1 \\ 0 & 2 & 0 \\ 1 & 0 & 1 \end{pmatrix},$$

- (a) Find all the eigenvalues of  $\hat{H}$ , and the normalised eigenvector of the ground state (only).
- (b) The ion’s state at time  $t = 0$  is  $|m = -1\rangle$ . Analyse this into a component in the ground state and a component orthogonal to the ground state, and hence find the state at a later time  $t$ .
5. **2009 Exam Q1(c)** A hydrogen atom is in the state

$$|\psi\rangle = N(|1, 0, 0\rangle + 3|2, 1, -1\rangle + 2|3, 2, 2\rangle).$$

where  $|n, l, m\rangle$  is a hydrogenic state with the usual quantum number notation. If the  $z$  component of orbital angular momentum is measured, give the possible results with their probabilities of occurrence. [6 marks]

1. (i) Bra (operator acts to left on a bra to make a new bra).
  - (ii) Operator (outer products are operators; this is the projector to state  $|1\rangle$ ).
  - (iii) Ket (direct product of two kets is a ket in the direct product space).
  - (iv) Complex numbers (inner products are complex numbers).
2.  $|\leftarrow\rangle\langle\leftarrow|$  is the projection operator onto state  $|\leftarrow\rangle$ ; call it  $P$ . First basis ket is  $|\uparrow\rangle$ , second is  $|\downarrow\rangle$ , so

$$\begin{aligned} P_{11} &= \langle\uparrow|P|\uparrow\rangle = \langle\uparrow|\leftarrow\rangle\langle\leftarrow|\uparrow\rangle = \frac{1}{2}\langle\uparrow|(|\uparrow\rangle + |\downarrow\rangle)\rangle(\langle\uparrow| + \langle\downarrow|)|\uparrow\rangle \\ &= \frac{1}{2}(\langle\uparrow|\uparrow\rangle + \langle\uparrow|\downarrow\rangle)(\langle\uparrow|\uparrow\rangle + \langle\downarrow|\uparrow\rangle) = \frac{1}{2}(1+0)(1+0) = \frac{1}{2} \end{aligned}$$

using the orthonormality of  $\{|\uparrow\rangle, |\downarrow\rangle\}$ . Similarly:

$$\begin{aligned} P_{12} &= \langle\uparrow|P|\downarrow\rangle = \langle\uparrow|\leftarrow\rangle\langle\leftarrow|\downarrow\rangle = \frac{1}{2}(1+0)(0+1) = \frac{1}{2} \\ P_{21} &= \langle\downarrow|P|\uparrow\rangle = \langle\downarrow|\leftarrow\rangle\langle\leftarrow|\uparrow\rangle = \frac{1}{2}(0+1)(1+0) = \frac{1}{2} \\ P_{22} &= \langle\downarrow|P|\downarrow\rangle = \langle\downarrow|\leftarrow\rangle\langle\leftarrow|\downarrow\rangle = \frac{1}{2}(0+1)(0+1) = \frac{1}{2} \end{aligned}$$

So the matrix

$$[P] = \begin{pmatrix} P_{11} & P_{12} \\ P_{21} & P_{22} \end{pmatrix} = \begin{pmatrix} \frac{1}{2} & \frac{1}{2} \\ \frac{1}{2} & \frac{1}{2} \end{pmatrix}.$$

This is the one we met in the last example class.

3. (i)

$$\hat{M} = \hat{S}_z^2 + 2\hbar^2$$

This is an operator equation so all terms should be operators. However, since multiplying by the identity has no effect we often don't bother writing  $\hat{I}$ ; but to be accurate we should say

$$\hat{M} = \hat{S}_z^2 + 2\hbar^2\hat{I}.$$

$$\hat{J} = \hat{L} + \hat{S}$$

$J$  is an operator that acts on the combined (direct product) space of orbital angular momentum and spin, but  $\hat{L}$  and  $\hat{S}$  act only on one (each) of these factor spaces. To be precise we should write

$$\hat{J} = \hat{L} \otimes \hat{I} + \hat{I} \otimes \hat{S}.$$

(ii) In the equation

$$\hat{L}_z|m = 0\rangle = 0,$$

the '0' on the right-hand side is the zero vector (zero ket), since the LHS (operator acting on ket) evaluates to a ket, not to a scalar.

4. (a) Given that we know  $H = AL_x^2$ , which is a function of  $L_x$ , its eigenvalues will be the same functions of the  $L_x$  eigenvalues.  $L_x$  must have the same eigenvalues as  $L_z$ , since  $x$  and  $z$  are just conventional labels. For  $l = 1$ , the eigenvalues of  $L_z$  and  $L_x$  are  $\hbar$ ,  $0$ , and  $-\hbar$ . Hence we expect eigenvalues of  $H$  to be

$$A\hbar^2 \times (1, 0, 1)$$

Check by explicitly solving the matrix. For speed, solve the matrix without the pre-factor (call it  $H'$ ) and then multiply the eigenvalues by  $A\hbar^2/2$ :

$$\begin{aligned} \det(H' - mI) &= \begin{vmatrix} 1 - m & 0 & 1 \\ 0 & 2 - m & 0 \\ 1 & 0 & 1 - m \end{vmatrix} \\ &= (1 - m)[(2 - m) \times (1 - m) - 0] - 0 + 1 \times [0 - (2 - m) \times 1] \\ &= (2 - m)[1 - 2m + m^2 - 1] = (2 - m)m(m - 2), \end{aligned}$$

i.e.  $m = (2, 0, 2)$  and hence energy eigenvalues are  $E = A\hbar^2(1, 0, 1)$  as expected.

The ground state is the lowest energy value, in this case  $E = 0$ . Hence the eigenvector satisfies:

$$\begin{pmatrix} 1 & 0 & 1 \\ 0 & 2 & 0 \\ 1 & 0 & 1 \end{pmatrix} \begin{pmatrix} a \\ b \\ c \end{pmatrix} = 0,$$

i.e.  $a + c = 0$  and  $2b = 0$ . Hence the normalised ground state ket is

$$|E = 0\rangle \xrightarrow{L_z} \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 0 \\ -1 \end{pmatrix}.$$

- (b) The state with  $m = -1$  is an eigenstate of  $L_z$ , hence one of our basis kets; in fact it is the third basis ket as we always list the eigenvalues in the order  $\hbar$ ,  $0$ ,  $-\hbar$ . So

$$|m = -1\rangle \xrightarrow{L_z} \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

Per the question, write this as  $|m = -1\rangle = a|E = 0\rangle + |\phi\rangle$  where  $a$  and  $|\phi\rangle$  are unknown, but  $\langle\phi|E = 0\rangle = 0$ , i.e.  $|\phi\rangle$  is a vector in the subspace with excited energy  $E = A\hbar^2$ .

$$\begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} = \frac{a}{\sqrt{2}} \begin{pmatrix} 1 \\ 0 \\ -1 \end{pmatrix} + \begin{pmatrix} x \\ y \\ z \end{pmatrix}$$

where

$$(1, 0, -1) \begin{pmatrix} x \\ y \\ z \end{pmatrix} = (x - z) = 0,$$

so the excited subspace contains any vector with  $x = z$ . The solution may be obvious; formally we have to solve:

$$\begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} = \begin{pmatrix} a/\sqrt{2} + x \\ y \\ -a/\sqrt{2} + x \end{pmatrix},$$

hence  $y = 0$ ,  $x = -a/\sqrt{2}$ , and  $-\sqrt{2}a = 1$ :

$$\begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 1 \\ 0 \\ 1 \end{pmatrix} - \frac{1}{2} \begin{pmatrix} 1 \\ 0 \\ -1 \end{pmatrix}.$$

The two terms have energies  $A\hbar^2$  and 0, so we can apply the time evolution operator to get the state at time  $t$ :

$$|\Psi(t)\rangle = e^{-iHt/\hbar}|m = -1\rangle = e^{-iA\hbar t}|\phi\rangle - e^0 \frac{1}{\sqrt{2}}|E = 0\rangle \xrightarrow{L_z} \frac{1}{2} \begin{pmatrix} e^{-iA\hbar t} - 1 \\ 0 \\ e^{-iA\hbar t} + 1 \end{pmatrix}.$$

This is a bit more illuminating if we pull out an overall phase factor:

$$|\Psi(t)\rangle \xrightarrow{L_z} \frac{e^{-iA\hbar t/2}}{2} \begin{pmatrix} e^{-iA\hbar t/2} - e^{iA\hbar t/2} \\ 0 \\ e^{-iA\hbar t/2} + e^{iA\hbar t/2} \end{pmatrix} = e^{-iA\hbar t/2} \begin{pmatrix} -i \sin(A\hbar t/2) \\ 0 \\ \cos(A\hbar t/2) \end{pmatrix}.$$

The ion oscillates between  $L_z = -\hbar$  and  $L_z = \hbar$ , but without passing through  $L_z = 0$ .

5. First find the normalization constant  $N$ :

$$\langle\psi|\psi\rangle = 1 = |N|^2(1 + 3^2 + 2^2) = 14N^2$$

using the orthogonality of states with different eigenvalues. Hence  $N^2 = 1/14$ . The  $z$  component of angular momentum is given by the  $m$  quantum number, namely  $L_z = m\hbar$ , so the possibilities are:  $L_z = 0$ , probability  $1/14$ ;  $L_z = -\hbar$ , probability  $9/14$ ;  $L_z = 2\hbar$ , probability  $4/14$ .