CMB Cosmology, Particle Physics and What all this has to do with Recombination



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KEK Theory Meeting on Particle Physics Phenomenology

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Main Goals for this Lecture

- Give an overview of some of the recent CMB results
- Explain why the CMB anisotropies link early-universe, particle and recombination physics
- Motivate how CMB spectral distortion (in particular from z~1000) could help disentangling effects in the future
- Convince you that the CMB holds many additional treasures for us, promising an exciting future for CMB cosmology



Cosmic Microwave Background Anisotropies



Planck all sky map

CMB has a blackbody spectrum in every direction
tiny variations of the CMB temperature Δ*T*/*T* ~ 10⁻⁵

WMAP CMB Sky









Other cosmological Dataset:

small-scale CMB, Supernovae, large-scale structure/ BAO, Lyman- α forest, lensing, ...



BAO, Lyman- α forest, lensing, ...

Dependence of the Power Spectrum on the Main Cosmological Parameters



- Total density (curvature)
 → positions of peaks
- dark energy
 → ISW at large scales
- Baryon density
 → damping tail / ratio of peaks
- dark matter
 - → gravitational driving / enhancement of third peak over second
- spectral index n_S
 → tilt of the overall power spectrum
- Thomson optical depth au
- \rightarrow large scale E-mode polarization
- \rightarrow damping tail

Cosmic Microwave Background Anisotropies



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Cosmic Microwave Background Anisotropies



Planck all sky map

CMB has a blackbody spectrum in every direction

• tiny variations of the CMB temperature $\Delta T/T \sim 10^{-5}$

Dramatic improvements in angular resolution and sensitivity over the past decades!



Cosmic Microwave Background Anisotropies with ACT



ACT - collaboration, 148 GHz Map, Hajian et al. 2010

~ 0.02 degree beam!

CMB anisotropies clearly taught us a lot about the Universe we live in!



Precisio	Tiny error bars				
Class	Parameter	$W\!M\!AP$ 7-year $\rm ML^a$	$WMAP+BAO+H_0 ML$	WMAP 7-year Mean ^b	$WMAP+BAO+H_0$ Mean
Primary	$100\Omega_b h^2$ $\Omega_c h^2$	2.270 0.1107	2.246 0.1120	$2.258^{+0.057}_{-0.056}\\0.1109\pm0.0056$	2.260 ± 0.053 0.1123 ± 0.0035
	$\hat{\Omega}_{\Lambda}$ n_s $ au$ $ au$ $\Delta_{\mathcal{P}}^2(k_0)^{\mathrm{c}}$	0.738 0.969 0.086 2.38×10^{-9}	0.728 0.961 0.087 2.45×10^{-9}	$\begin{array}{c} 0.734 \pm 0.029 \\ 0.963 \pm 0.014 \\ 0.088 \pm 0.015 \\ (2.43 \pm 0.11) \times 10^{-9} \end{array}$	$\begin{array}{c} 0.728^{+0.015}_{-0.016}\\ 0.963\pm 0.012\\ 0.087\pm 0.014\\ (2.441^{+0.088}_{-0.092})\times 10^{-9}\end{array}$
Derived	$ \frac{\sigma_8}{H_0} $ $ \frac{\Omega_b}{\Omega_c} $ $ \Omega_m h^2 $ $ \frac{\sigma_8}{\Omega_c} $	0.803 71.4 km/s/Mpc 0.0445 0.217 0.1334	0.807 70.2 km/s/Mpc 0.0455 0.227 0.1344	$\begin{array}{c} 0.801 \pm 0.030 \\ 71.0 \pm 2.5 \text{ km/s/Mpc} \\ 0.0449 \pm 0.0028 \\ 0.222 \pm 0.026 \\ 0.1334^{+0.0056}_{-0.0055} \\ 10.5 \pm 1.2 \end{array}$	$\begin{array}{c} 0.809 \pm 0.024 \\ \hline 0.809 \pm 0.024 \\ \hline 70.4^{+1.3}_{-1.4} \text{ km/s/Mpc} \\ 0.0456 \pm 0.0016 \\ 0.227 \pm 0.014 \\ 0.1349 \pm 0.0036 \\ \hline 10.4 \pm 1.2 \end{array}$
	${z_{reion}}^{a}$	10.3 13.71 Gyr	10.5 13.78 Gyr	10.5 ± 1.2 $13.75 \pm 0.13 \text{ Gyr}$	10.4 ± 1.2 $13.75 \pm 0.11 \text{ Gyr}$

^aLarson et al. (2010). "ML" refers to the Maximum Likelihood parameters.

^bLarson et al. (2010). "Mean" refers to the mean of the posterior distribution of each parameter. The quoted errors show the 68% confidence levels (CL).

 $^{c}\Delta_{\mathcal{R}}^{2}(k) = k^{3}P_{\mathcal{R}}(k)/(2\pi^{2})$ and $k_{0} = 0.002 \text{ Mpc}^{-1}$.

^d "Redshift of reionization," if the universe was reionized instantaneously from the neutral state to the fully ionized state at z_{reion} . Note that these values are somewhat different from those in Table 1 of Komatsu et al. (2009b), largely because of the changes in the treatment of reionization history in the Boltzmann code CAMB (Lewis 2008).

^eThe present-day age of the universe.





e.g. Komatsu et al., 2011, ApJ, arXiv:1001.4538 Dunkley et al., 2011, ApJ, arXiv:1009.0866

CMB anisotropies clearly taught us a lot about the Universe we live in!







TABLE I. Standard Λ CDM parameters from the combination of WMAP9, ACT and SPT.

Parameter	WMAP9	WMAP9	WMAP9		
	+ACT	+SPT	+ACT+SPT		
$100\Omega_b h^2$	2.260 ± 0.041	2.231 ± 0.034	2.245 ± 0.032		
$100\Omega_c h^2$	11.46 ± 0.43	11.16 ± 0.36	11.23 ± 0.36		
$100\theta_A$	1.0396 ± 0.0019	1.0422 ± 0.0010	1.0420 ± 0.0010		
au	0.090 ± 0.014	0.082 ± 0.013	0.085 ± 0.013		
n_s	0.973 ± 0.011	0.9650 ± 0.0093	0.9678 ± 0.0088		
$10^9 \Delta_R^2$	2.22 ± 0.10	2.15 ± 0.10	2.17 ± 0.10		
$\Omega_{\Lambda}{}^{a}$	0.716 ± 0.024	0.737 ± 0.019	0.734 ± 0.019		
σ_8	0.830 ± 0.021	0.808 ± 0.018	0.814 ± 0.017		
t_0	13.752 ± 0.096	13.686 ± 0.065	13.682 ± 0.063		
H_0	69.7 ± 2.0	71.5 ± 1.7	71.2 ± 1.6		
$100r_s/D_{V0.57}$	7.50 ± 0.17	7.65 ± 0.14	7.65 ± 0.14		
$100r_s/D_{V0.35}$	11.29 ± 0.31	11.56 ± 0.26	11.55 ± 0.26		
best fit χ^2	7596.0	7617.1	7660.0		

Precision Cosmology with Planck

	Planck+WP		Planck+WP+highL		Planck+lensing+WP+highL		Planck+WP+highL+BAO	
Parameter	Best fit	68% limits	Best fit	68% limits	Best fit	68% limits	Best fit	68% limits
$\Omega_{\rm b}h^2$	0.022032	0.02205 ± 0.00028	0.022069	0.02207 ± 0.00027	0.022199	0.02218 ± 0.00026	0.022161	0.02214 ± 0.00024
$\Omega_c h^2 \dots \dots$	0.12038	0.1199 ± 0.0027	0.12025	0.1198 ± 0.0026	0.11847	0.1186 ± 0.0022	0.11889	0.1187 ± 0.0017
100θ _{MC}	1.04119	1.04131 ± 0.00063	1.04130	1.04132 ± 0.00063	1.04146	1.04144 ± 0.00061	1.04148	1.04147 ± 0.00056
τ	0.0925	$0.089^{+0.012}_{-0.014}$	0.0927	$0.091^{+0.013}_{-0.014}$	0.0943	$0.090^{+0.013}_{-0.014}$	0.0952	0.092 ± 0.013
<i>n</i> _s	0.9619	0.9603 ± 0.0073	0.9582	0.9585 ± 0.0070	0.9624	0.9614 ± 0.0063	0.9611	0.9608 ± 0.0054
$\ln(10^{10}A_s)$	3.0980	$3.089^{+0.024}_{-0.027}$	3.0959	3.090 ± 0.025	3.0947	3.087 ± 0.024	3.0973	3.091 ± 0.025
$A_{100}^{\rm PS}$	152	171 ± 60	209	212 ± 50	204	213 ± 50	204	212 ± 50
A_{143}^{PS}	63.3	54 ± 10	72.6	73 ± 8	72.2	72 ± 8	71.8	72.4 ± 8.0
$A_{217}^{\rm PS}$	117.0	107^{+20}_{-10}	59.5	59 ± 10	60.2	58 ± 10	59.4	59 ± 10
A ^{CIB} ₁₄₃	0.0	< 10.7	3.57	3.24 ± 0.83	3.25	3.24 ± 0.83	3.30	3.25 ± 0.83
A ^{CIB} ₂₁₇	27.2	29 ₋₉ ⁺⁶	53.9	49.6 ± 5.0	52.3	50.0 ± 4.9	53.0	49.7 ± 5.0
A ^{tSZ} ₁₄₃	6.80		5.17	$2.54^{+1.1}_{-1.9}$	4.64	$2.51^{+1.2}_{-1.8}$	4.86	$2.54^{+1.2}_{-1.8}$
$r_{143\times217}^{\rm PS}$	0.916	> 0.850	0.825	$0.823^{+0.069}_{-0.077}$	0.814	0.825 ± 0.071	0.824	0.823 ± 0.070
r ^{CIB} _{143×217}	0.406	0.42 ± 0.22	1.0000	> 0.930	1.0000	> 0.928	1.0000	> 0.930
γ^{CIB}	0.601	$0.53^{+0.13}_{-0.12}$	0.674	0.638 ± 0.081	0.656	0.643 ± 0.080	0.667	0.639 ± 0.081
$\xi^{tSZ \times CIB}$	0.03		0.000	< 0.409	0.000	< 0.389	0.000	< 0.410
A^{kSZ}	0.9		0.89	5.34 ^{+2.8} -1.9	1.14	$4.74^{+2.6}_{-2.1}$	1.58	$5.34^{+2.8}_{-2.0}$
Ω_{Λ}	0.6817	$0.685^{+0.018}_{-0.016}$	0.6830	0.685+0.017	0.6939	0.693 ± 0.013	0.6914	0.692 ± 0.010
<i>σ</i> ₈	0.8347	0.829 ± 0.012	0.8322	0.828 ± 0.012	0.8271	0.8233 ± 0.0097	0.8288	0.826 ± 0.012
Z _{re}	11.37	11.1 ± 1.1	11.38	11.1 ± 1.1	11.42	11.1 ± 1.1	11.52	11.3 ± 1.1
H_0	67.04	67.3 ± 1.2	67.15	67.3 ± 1.2	67.94	67.9 ± 1.0	67.77	67.80 ± 0.77
Age/Gyr	13.8242	13.817 ± 0.048	13.8170	13.813 ± 0.047	13.7914	13.794 ± 0.044	13.7965	13.798 ± 0.037
100 <i>0</i> ,	1.04136	1.04147 ± 0.00062	1.04146	1.04148 ± 0.00062	1.04161	1.04159 ± 0.00060	1.04163	1.04162 ± 0.00056
<i>r</i> _{drag}	147.36	147.49 ± 0.59	147.35	147.47 ± 0.59	147.68	147.67 ± 0.50	147.611	147.68 ± 0.45

- Massive amount of information! (close to 30 Planck papers in March 2013)
- Impressive consistency between different experiments!
- Amazing confirmation of ACDM



Planck Satellite

CMB anisotropies directly probe early-universe physics / inflation

Another way to plot small-scale power spectrum



- 6σ deviation from scale-invariance (previously ~ 3σ)
- single-field inflation predicts departure from scale-invariance (e.g., Mukhanov 2007)
- Degeneracies with, e.g., effective number of relativistic degrees of freedom, N_{eff}, Helium abundance, Y_p, and recombination physics!
- The power spectrum at small scales thus directly links early-Universe, particle and recombination physics!

All kind of fun science with the CMB (no time for this though)





Planck Collaboration, 2013, paper XVII

SZ clusters on the sky







Planck Collaboration, 2013, paper XXVII

- Non-Gaussianity (test of inflation models)
- Topology
- CMB anomalies
- CIB and Galactic science

Extension of ACDM and why these link earlyuniverse, particle and recombination physics

Simplest one parameter extensions of ACDM

	Planck+WP		Planck+WP+BAO		Planck-	Planck+WP+highL		Planck+WP+highL+BAO	
Parameter	Best fit	95% limits	Best fit	95% limits	Best fit	95% limits	Best fit	95% limits	
Ω _K	-0.0105	$-0.037^{+0.043}_{-0.049}$	0.0000	$0.0000^{+0.0066}_{-0.0067}$	-0.0111	$-0.042^{+0.043}_{-0.048}$	0.0009	$-0.0005^{+0.0065}_{-0.0066}$	
$\Sigma m_{\nu} [eV] \ldots \ldots$	0.022	< 0.933	0.002	< 0.247	0.023	< 0.663	0.000	< 0.230	
$N_{\rm eff}$	3.08	$3.51^{+0.80}_{-0.74}$	3.08	$3.40^{+0.59}_{-0.57}$	3.23	$3.36^{+0.68}_{-0.64}$	3.22	$3.30^{+0.54}_{-0.51}$	
$Y_{\rm P}$	0.2583	$0.283_{-0.048}^{+0.045}$	0.2736	$0.283^{+0.043}_{-0.045}$	0.2612	$0.266^{+0.040}_{-0.042}$	0.2615	$0.267^{+0.038}_{-0.040}$	
$dn_{\rm s}/d\ln k$	-0.0090	$-0.013^{+0.018}_{-0.018}$	-0.0102	$-0.013^{+0.018}_{-0.018}$	-0.0106	$-0.015^{+0.017}_{-0.017}$	-0.0103	$-0.014^{+0.016}_{-0.017}$	
$r_{0.002}$	0.000	< 0.120	0.000	< 0.122	0.000	< 0.108	0.000	< 0.111	
<i>w</i>	-1.20	$-1.49^{+0.65}_{-0.57}$	-1.076	$-1.13^{+0.24}_{-0.25}$	-1.20	$-1.51^{+0.62}_{-0.53}$	-1.109	$-1.13^{+0.23}_{-0.25}$	

Planck Collaboration, 2013, paper XV

- All consistent with standard \Lambda CDM
- slight tensions between different experiments (e.g., Neff, Yp and running)

CMB anisotropy constraints on running and the tensor to scalar ratio



• Single-field inflation: $dn/d \ln k \simeq (n_{\rm S}-1)^2$ • Big future goal. detection of B-polarization • Plenty of processing processing the next few years: ground/palle • plate of processing procesing processing procesing processing processing proces

n

Other experiments:

 $\mathrm{d}n/\mathrm{d}\ln k = -0.022\pm0.012$ Dunkley et al 2011 & Keisler et al 2011

 $\mathrm{d}n/\mathrm{d}\ln k = -0.003 \pm 0.013$ Sievers et al 2013

 $\mathrm{d}n/\mathrm{d}\ln k = -0.024\pm0.011$ Hou et al 2012

CMB as a test for BBN

Standard Big Bang Nucleosynthesis (SBBN)



Beyond SBBN



Beyond SBBN



Change in timing

non-equilibrium BBN

catalyzed BBN

Courtesy: to Josef Pradler

Beyond SBBN



Abundances of light-elements provide a unique test of non-standard BBN!

Change in timing

non-equilibrium BBN

catalyzed BBN

Courtesy: to Josef Pradler

SBBN Predictions for Helium and Deuterium



Planck Collaboration, 2013, paper XV

Interplay of N_{eff} and Y_p and other parameters



Hinshaw et al, 2012 (WMAP-9yr)

Bottom line: changes in the damping tail can be mimics by combination of many parameters

CMB constraints on N_{eff} and Y_p



Planck Collaboration, 2013, paper XV

- Helium determination from CMB consistent with SBNN prediction
- CMB constraint on N_{eff} competitive
- Partial degeneracy with Y_p and running
- Some tension between different data sets



CMB constraints on N_{eff} and Y_p



Consistent with SBBN and standard value for N_{eff}

• Future CMB constraints (SPTPol & ACTPol) on Yp will reach 1% level

Cosmological Time in Years



Cosmological Time in Years





BAO, Lyman- α forest, lensing, ...



small-scale CMB, Supernovae, large-scale structure/ BAO, Lyman- α forest, lensing, ...

How does cosmological recombination work?
Physical Conditions during Recombination

- Temperature $T_{\gamma} \sim 2.725 (1+z) \text{ K} \sim 3000 \text{ K}$
- Baryon number density $N_{\rm b} \sim 2.5 \times 10^{-7} {\rm cm}^{-3} (1+z)^3 \sim 330 {\rm cm}^{-3}$
- Photon number density N_γ ~ 410 cm⁻³ (1+z)³ ~ 2×10⁹ N_b
 ⇒ photons in very distant Wien tail of blackbody spectrum can keep hydrogen ionized until hv_α ~ 40 kT_γ
- Collisional processes negligible (completely different from stars!!!)
- Rates dominated by radiative processes
 (e.g. stimulated emission & stimulated recombination)
- Compton interaction couples electrons very tightly to photons until $z \sim 200 \Rightarrow T_{\gamma} \sim T_e \sim T_m$





continuum: *e p* (He)



Routes to the ground state ?

Zeldovich, Kurt & Sunyaev, 1968, ZhETF, 55, 278 Peebles, 1968, ApJ, 153, 1

Hydrogen atom



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- direct recombination to 1s
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- recombination to 2p followed by Lyman-α emission
 - medium optically thick to Ly- α phot.
 - many resonant scatterings
 - escape very hard (*p* ~10⁻⁹ @ *z* ~1100)

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- recombination to 2s followed by 2s two-photon decay
 - 2s \rightarrow 1s ~10⁸ times slower than Ly- α
 - 2s two-photon decay profile \rightarrow maximum at $\nu \sim$ 1/2 ν_{α}
 - immediate escape

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Routes to the ground state ?

•	 direct recombination to 1s Emission of photon is followed by immediate re-absorption 	≻ No
•	recombination to 2p followed by Lyman- α emission	
	 medium optically thick to Ly-α phot. many resonant scatterings escape very hard (<i>p</i> ~10⁻⁹ @ <i>z</i> ~1100) 	∼ 43%
•	recombination to 2s followed by 2s two-photon decay	
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	- immediate escape	

 $\Delta N_{\rm e}$ / $N_{\rm e}$ ~ 10% - 20%

First recombination computations completed in 1968!



Yakov Zeldovich



Vladimir Kurt (UV astronomer)

Moscow

Princeton



Rashid Sunyaev



Jim Peebles

Multi-level Atom ⇒ The Recfast-Code



Seager, Sasselov & Scott, 1999, ApJL, 523, L1 Seager, Sasselov & Scott, 2000, ApJS, 128, 407 Output of $N_{\rm e}/N_{\rm H}$

Hydrogen:

- up to 300 levels
- only 2s & 2p separately
- $n>2 \rightarrow$ full SE for *l*-sub-states

Helium:

- Hel 200-levels (z ~ 1400-1500)
- Hell 100-levels (z ~ 6000-6500)
- Helll 1 equation

Low Redshifts:

- H chemistry (important at low z)
- cooling of matter (Bremsstrahlung, collisional cooling, line cooling)

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 $\Delta N_{\rm e}$ / $N_{\rm e}$ ~ 1% - 3%

Getting Ready for Planck

Hydrogen recombination

- Two-photon decays from higher levels (Dubrovich & Grachev, 2005, Astr. Lett., 31, 359; Wong & Scott, 2007; JC & Sunyaev, 2007; Hirata, 2008; JC & Sunyaev 2009)
- Induced 2s two-photon decay for hydrogen (JC & Sunyaev, 2006, A&A, 446, 39; Hirata 2008)
- Feedback of the Lyman- α distortion on the 1s-2s two-photon absorption rate (Kholupenko & Ivanchik, 2006, Astr. Lett.; Fendt et al. 2008; Hirata 2008)
- Non-equilibrium effects in the angular momentum sub-states (Rubiño-Martín, JC & Sunyaev, 2006, MNRAS; JC, Rubiño-Martín & Sunyaev, 2007, MNRAS; Grin & Hirata, 2009; JC, Vasil & Dursi, 2010)
- Feedback of Lyman-series photons (Ly[n] → Ly[n-1]) (JC & Sunyaev, 2007, A&A; Kholupenko et al. 2010; Haimoud, Grin & Hirata, 2010)
- Lyman-α escape problem (atomic recoil, time-dependence, partial redistribution) (Dubrovich & Grachev, 2008; JC & Sunyaev, 2008; Forbes & Hirata, 2009; JC & Sunyaev, 2009)
- Raman scattering (Hirata 2008; JC & Thomas , 2010; Haimoud & Hirata, 2010)

Helium recombination

- Similar list of processes as for hydrogen (Switzer & Hirata, 2007a&b; Hirata & Switzer, 2007)
- Spin forbidden 2p-1s triplet-singlet transitions (Dubrovich & Grachev, 2005, Astr. Lett.; Wong & Scott, 2007; Switzer & Hirata, 2007; Kholupenko, Ivanchik&Varshalovich, 2007)
- Hydrogen continuum opacity during He I recombination

(Switzer & Hirata, 2007; Kholupenko, Ivanchik & Varshalovich, 2007; Rubiño-Martín, JC & Sunyaev, 2007)

 Detailed feedback of helium photons (Switzer & Hirata, 2007a; JC & Sunyaev, 2009, MNRAS)







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Cumulative Changes to the Ionization History





JC & Thomas, MNRAS, 2010; Shaw & JC, 2011

Cumulative Change in the CMB Power Spectra





Importance of recombination for inflation



CITA General Addaption Technology endored Addaption Shaw & JC, 2011, and references therein

Importance of recombination for measuring helium



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What if something unexpected happened?

- E.g., something *standard* was missed, or something *non-standard* happened !?
- A non-parametric estimation of possible corrections to the recombination history would be very useful → Principle component analysis (PCA)



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Measured mode amplitudes for ACT & SPT

parameters	6s	$\begin{array}{c} {\rm SPT+WMAP7} \\ + {\rm mode} \ 1 \end{array}$	$+ \mod 2$	6s	$\begin{array}{r} \mathrm{ACT}+\mathrm{WMAP7} \\ + \bmod 1 \end{array}$	$+ \mod 2$
$100\Omega_{ m b}h^2$	2.221 ± 0.042	2.253 ± 0.046	2.249 ± 0.047	2.219 ± 0.051	2.240 ± 0.050	2.236 ± 0.053
$\Omega_{ m c}h^2$	0.1110 ± 0.0048	0.1123 ± 0.0049	0.1118 ± 0.0052	0.1121 ± 0.0052	0.1155 ± 0.0056	0.1121 ± 0.0061
$100\theta_{s}$	1.041 ± 0.002	1.041 ± 0.002	1.040 ± 0.003	1.039 ± 0.002	1.039 ± 0.002	1.035 ± 0.004
au	0.086 ± 0.015	0.089 ± 0.015	0.089 ± 0.015	0.086 ± 0.015	0.089 ± 0.015	0.0875 ± 0.015
$n_{ m s}$	0.964 ± 0.011	0.977 ± 0.013	0.975 ± 0.016	0.963 ± 0.013	0.976 ± 0.015	0.960 ± 0.019
$10^9 \Delta_R^2$	2.43 ± 0.10	2.40 ± 0.10	2.40 ± 0.10	2.45 ± 0.11	2.43 ± 0.11	2.45 ± 0.11
μ_1	(0)	-0.77 ± 0.46	-0.76 ± 0.47	(0)	-1.27 ± 0.74	-1.67 ± 0.86
μ_2	(0)	(0)	-0.39 ± 1.09	(0)	(0)	-3.5 ± 2.7
σ_8 (derived)	0.807 ± 0.024	0.825 ± 0.027	0.818 ± 0.032	0.814 ± 0.028	0.841 ± 0.031	0.802 ± 0.040
$\delta z_{ m dec}/z_{ m dec}$ ^a	-	-0.6%	-0.7%	-	-1.0%	-1.7%
$\delta \sigma_{z, m dec} / \sigma_{z, m dec}$ ^b	-	1.5%	-0.5%	-	2.6%	-14.0%
$(\delta x_{\rm e} /x_{\rm e})_{\rm max}{}^{\rm c}$	-	5% (z $\sim 1196)$	5% $(z\sim1039)$	-	8% (z $\sim 1006)$	$31\%~(z\sim1076)$
$\Delta\chi^2$	-	2.5	2.5	-	2.1	2.5

^arelative change in the redshift of maximum visibility where $z_{dec} = 1088$ is the fiducial maximum visibility point. ^brelative change in the width of the visibility function.

^cmaximum relative change in the ionization fraction. The redshift corresponding to this maximum change is also included.

- First mode detected at ~ 2σ
- Similar for current Planck data
- Effect very similar to the one of helium
- In the future 2-3 modes detectable





Is there another more direct way to constrain the cosmological recombination history?

COBE / FIRAS (Far InfraRed Absolute Spectrophotometer)



 $T_0 = 2.725 \pm 0.001 \,\mathrm{K}$ $|y| \le 1.5 \times 10^{-5}$ $|\mu| \le 9 \times 10^{-5}$

Mather et al., 1994, ApJ, 420, 439 Fixsen et al., 1996, ApJ, 473, 576 Fixsen et al., 2003, ApJ, 594, 67



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Only very small distortions of CMB spectrum are still allowed!

Simple estimates for hydrogen recombination

Hydrogen recombination:

 per recombined hydrogen atom an energy of ~ 13.6 eV in form of photons is released

- at $z \sim 1100 \rightarrow \Delta \epsilon/\epsilon \sim 13.6 \text{ eV } N_b / (N_\gamma 2.7 \text{k} T_r) \sim 10^{-9} \text{--} 10^{-8}$
- \rightarrow recombination occurs at redshifts $z < 10^4$
- At that time the *thermalization* process doesn't work anymore!
- There should be some small spectral distortion due to additional Ly-α and 2s-1s photons! (Zeldovich, Kurt & Sunyaev, 1968, ZhETF, 55, 278; Peebles, 1968, ApJ, 153, 1)
- → In 1975 *Viktor Dubrovich* emphasized the possibility to observe the recombinational lines from n > 3 and $\Delta n << n!$

100-shell hydrogen atom and continuum CMB spectral distortions



JC & Sunyaev, 2006, A&A, 458, L29 (astro-ph/0608120)

100-shell hydrogen atom and continuum CMB spectral distortions



JC & Sunyaev, 2006, A&A, 458, L29 (astro-ph/0608120)

100-shell hydrogen atom and continuum CMB spectral distortions



JC & Sunyaev, 2006, A&A, 458, L29 (astro-ph/0608120)

100-shell hydrogen atom and continuum Relative distortions



JC & Sunyaev, 2006, A&A, 458, L29 (astro-ph/0608120)

Wien-region:

- L_{α} and 2s distortions
 - are very strong
- but CIB more dominant

@ CMB maximum:

- relative distortions extremely small
- strong v-dependence

RJ-region:

- relative distortion exceeds
 level of ~10⁻⁷ below v ~
 1-2 GHz
- oscillatory frequency dependence with ~1-10 percent-level amplitude:
- hard to mimic by known
 foregrounds or systematics

Cosmological Time in Years



What about the contributions from helium recombination?

• Nuclear reactions: $Y_p \sim 0.24 \leftrightarrow N_{Hel} / N_H \sim 8 \%$

 \rightarrow expected photon number rather small

• *BUT:*

(i) two epochs of He recombination HeIII→HeII at z~6000 and HeII→HeI at z~2500
(ii) Helium recombinations faster → more narrow features with larger amplitude
(iii) non-trivial superposition

 \rightarrow local amplification possible

(iv) reprocessing of HeII & HeI photons by HeI and HI

 \rightarrow increases the number of helium-related photons

Any opens a way to *directly* measure the primordial (pre-stellar!!!) helium abundance!

Grotrian diagram for neutral helium



Helium contributions to the cosmological recombination spectrum







Cosmological Time in Years


What would we actually learn by doing such hard job?

Cosmological Recombination Spectrum opens a way to measure:

- \rightarrow the specific *entropy* of our universe (related to $\Omega_{b}h^{2}$)
- \rightarrow the CMB *monopole* temperature T_0
- \rightarrow the pre-stellar abundance of helium Y_p

→ If recombination occurs as we think it does, then the lines can be predicted with very high accuracy!

→ In principle allows us to directly check our understanding of the standard recombination physics



Figure 7.3: The 1 and 2 dimensional marginalized parameter posterior using the CMB spectral distortions. All three cases constrain the CMB power spectrum using a Gaussian likelihood based on Planck noise levels. The black line adds constraints due to a 10% measurement of the spectral distortions, while the blue line assumes a 1% measurement. The red line does not include the data from the spectral distortions. CMB based cosmology alone

 Spectrum helps to break some of the parameter degeneracies

 Planning to provide a module that computes the recombination spectrum in a fast way

detailed forecasts: which lines to measure; how important is the absolute amplitude; how accurately one should measure; best frequency resolution;

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But is the standard cosmological recombination spectrum really interesting enough?

Extra Sources of Ionizations or Excitations



- ,Hypothetical' source of extra photons parametrized by ε_α & ε_i
- Extra excitations \Rightarrow delay of Recombination
- Extra ionizations ⇒ affect 'freeze out' tail
- This affects the Thomson visibility function
- From WMAP $\Rightarrow \epsilon_{\alpha} < 0.39 \& \epsilon_i < 0.058$ at 95% confidence level (Galli et al. 2008)
- Extra ionizations & excitations should also lead to additional photons in the recombination radiation!!!
- This in principle should allow us to check for such sources at z~1000

Peebles, Seager & Hu, ApJ, 2000

Dark Matter Annihilation: Energy Branching Ratios





Efficiencies according to Chen & Kamionkowski, 2004 & Shull & van Steenberg 1985

- N^2 dependence $\Rightarrow dE/dt \propto (1+z)^6$ and $dE/dz \propto (1+z)^{3...3.5}$
- only part of the energy is really deposited ($f_d \sim 0.1$)
- Branching into heating (100% at high z), ionizations and excitations (mainly during recombination)
- Branching depends on considered DM model

Dark Matter Annihilation: Effect on CMB Anisotropies and the Recombination Spectrum





- 'Delay of recombination'
- Affects Thomson visibility function
- Possibility of Sommerfeld-enhancement
- Clumpiness of matter at z<100

- Additional photons at all frequencies
- Broadening of spectral features
- Shifts in the positions

CMB limits on annihilation cross section

95% c.l.



Galli et al, 2011 (similar to Huetsi et al, 2011)

What could the recombination spectrum add?

- WMAP constraints on possible dark matter annihilation efficiencies already very tight (e.g. see Galli et al. 2009; Slatyer et al. 2009, Huetsi et al., 2009, Huetsi et al., 2011, Galli et al., 2011)
 - absolute changes to CMB power spectra have to be small (~ 1%-5%)
 - changes to cosmological recombination spectrum are of similar order

What could the recombination spectrum add?

- WMAP constraints on possible dark matter annihilation efficiencies already very tight (e.g. see Galli et al. 2009; Slatyer et al. 2009, Huetsi et al., 2009, Huetsi et al., 2011, Galli et al., 2011)
 - absolute changes to CMB power spectra have to be small (~ 1%-5%)
 - changes to cosmological recombination spectrum are of similar order
- So why bother anymore? What could the cosmological recombination spectrum teach us in addition?
 (JC, 2009, arXiv:0910.3663)
 - spectrum is sensitive to cases for which the C_l's are not affected!
 - DM annihilation parameters are ,degenerate' with $n_{\rm S}$ & $\Omega_{\rm b}h^2$
 - spectrum could help breaking this degeneracy
 - very direct way to check for sources of extra ionizations and excitations during all three recombination epochs
 - broad y and µ distortions will give another handle! (see tomorrow)

Decaying particle during & after recombination



- Modify recombination history
- this changes Thomson visibility function and thus the CMB temperature and polarization power spectra
- ⇒ CMB anisotropies allow

probing particles with lifetimes $\gtrsim 10^{12}$ sec

 CMB spectral distortions provide complementary probe! (more tomorrow)

Chen & Kamionkowski, 2004

What would we actually learn by doing such hard job?

Cosmological Recombination Spectrum opens a way to measure:

- \rightarrow the specific *entropy* of our universe (related to $\Omega_{b}h^{2}$)
- \rightarrow the CMB *monopole* temperature T_0
- \rightarrow the pre-stellar abundance of helium Y_p
- → If recombination occurs as we think it does, then the lines can be predicted with very high accuracy!

→ In principle allows us to directly check our understanding of the standard recombination physics

If something unexpected or non-standard happened:

- → non-standard thermal histories should leave some measurable traces
- → direct way to measure/reconstruct the recombination history!
- → possibility to distinguish pre- and post-recombination y-type distortions
- → sensitive to energy release during recombination
- → variation of fundamental constants

Conclusions

- CMB anisotropies provide an outstanding confirmation of ACDM cosmology
- The data has become so precise that one can start testing SBBN and non-standard extensions of ACDM
- The recombination process is crucial for the interpretation of the data at this level of precision
 - Future observation of the cosmological recombination radiation will allow confirming the recombination model
- If something non-standard happened around z~1000, then this should show up in the recombination spectrum, allowing us to break degeneracies and providing independent confirmation

Cosmological Time in Years

